

The NIST 60-Millimeter Diameter Cylindrical Cavity Resonator: Performance Evaluation for Permittivity Measurements

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Abstract

Uncertainty estimates are developed for dielectric permittivity calculations made using the NIST 60-mm diameter cylindrical resonator. A mode-filtering helical waveguide makes up the cavity's cylindrical wall, which permits the generation of high-purity TE_{01p} resonant modes for high accuracy permittivity measurements. The cavity's length can be varied from 408 to 433 mm. Fixed-length and fixed-frequency measurements in the X-band frequency range are evaluated with particular emphasis on 10 GHz. Resonator theory and design, measurement tolerances, and software are included.

Keywords: cylindrical cavity, dielectric permittivity, measurement software, resonator theory and design, TE_{01} helical waveguide, uncertainty estimates, x-band 10 GHz microwave.

Chapter 1

Introduction

1.1 Definition of Permittivity

Permittivity is a physical quantity that describes the relation between electric field \bar{E} in volts per meter and electric displacement \bar{D} in coulombs per square meter. In a vacuum, this relation is given by

$$\bar{D} = \epsilon_0 \bar{E}, \quad (1.1)$$

where $\epsilon_0 = 8.85418782 \times 10^{-12}$ farads per meter is the permittivity of free space. In matter, electric displacement is affected by the material's electric susceptibility, χ_e^* , which is a measure of the material's polarizability. In an isotropic, linear media the polarization vector \bar{P} is parallel and proportional to the electric field,

$$\bar{P} = \epsilon_0 \chi_e^* \bar{E}, \quad (1.2)$$

and the electric displacement vector then becomes

$$\begin{aligned} \bar{D} &= \epsilon_0 \bar{E} + \bar{P} \\ &= \epsilon_0 (1 + \chi_e^*) \bar{E} \\ &= \epsilon^* \bar{E}, \end{aligned} \quad (1.3)$$

where

$$\begin{aligned} \epsilon^* &= \epsilon_0 (1 + \chi_e^*), \\ &= \epsilon_0 \epsilon_R^* \end{aligned} \quad (1.4)$$

is the permittivity of the material, and $\epsilon_R^* = (1 + \chi_e^*)$ is the material's relative permittivity or dielectric constant. In general, electric susceptibility has real

1.2. OPERATION OF MODE-FILTERED CYLINDRICAL RESONATORS 3

(i.e., dispersive) and imaginary (i.e., absorptive) parts which makes relative permittivity a complex quantity. We define relative permittivity as

$$\begin{aligned}\epsilon_R^* &= 1 + \chi_e^* \\ &= (1 + \Re\{\chi_e^*\}) - j\Im\{\chi_e^*\} \\ &= \epsilon'_R - j\epsilon''_R.\end{aligned}\tag{1.5}$$

For low-loss materials ϵ'_R is often loosely called the relative permittivity, with the distinction between ϵ_R^* and ϵ'_R implicitly understood. The term ϵ''_R is called the dielectric loss factor.

If the induced conduction current is directly proportional to the electric field, the curl of the magnetic field \overline{H} is

$$\begin{aligned}\nabla \times \overline{H} &= (j\omega\epsilon^* + \sigma)\overline{E}, \\ &= j\omega\epsilon_0\epsilon'_R \left(1 - j\frac{\epsilon''_R}{\epsilon'_R} - j\frac{\sigma}{\omega\epsilon_0\epsilon'_R}\right)\overline{E}, \\ &= j\omega\epsilon_0\epsilon'_R(1 - j\tan\delta)\overline{E}.\end{aligned}\tag{1.6}$$

This relation for loss tangent $\tan\delta$ now explicitly separates dielectric and conductive losses. In most discussions dielectric loss of nonconductive materials at microwave and millimeter wave frequencies is described by a frequency dependent conductivity of the material at the measurement frequency. In this case loss tangent $\epsilon''_R = \sigma/\omega\epsilon_0$, and loss tangent becomes

$$\tan\delta = \frac{\epsilon''_R}{\epsilon'_R} = \frac{\sigma}{\omega\epsilon_0\epsilon'_R}.\tag{1.7}$$

1.2 Operation of Mode-Filtered Cylindrical Resonators

The cylindrical cavity resonator described in this Technical Note was built by the NIST Electromagnetic Properties of Materials program to accurately measure the complex permittivity of dielectric materials. The application of cavity resonator methods to measure the complex permittivity of dielectric samples is not new [1]. However, resonator methods are inherently accurate and to this day prove to be the most accurate measurement method available. NIST chose to build a cavity resonator to accurately characterize dielectric materials that in turn can be used as check standards to test the validity of other measurement methods. Special test services are performed regularly for those outside NIST,

and, since its introduction, the NIST resonator has characterized a large number of materials both for internal and external use.

Cavity resonators can measure permittivity only at fixed-frequencies and over a narrow range of frequency. This is the case especially for the shorter, solid-walled cavities because higher-order modes occur very easily. Measurement of permittivity over a broader range of frequencies is partially solved by the use of helical waveguide cavities [2, 3]. The mode filtering capability of helical waveguides allows construction of cavity resonators that can support several highly pure modes in a limited frequency range.

The NIST resonator has a 60-mm diameter helical wire-wound cylindrical wall, with a variable length of approximately 408 to 433 mm. The cavity's helical winding allows only circumferential currents to flow, which permits propagation of only the TE_{0n} family of circular transmission line modes, where n describes the number of radial variations. Mode-filtering capability permits the use of a long cavity that would otherwise be highly overmoded. This allows permittivity measurements at several TE_{01} resonant frequencies throughout X-band (8.2 to 12.4 GHz) rather than at only one or two specific frequencies.

The well defined geometry of the NIST cylindrical resonator allows standing waves to occur at certain frequencies. Power enters through a coupling iris located in one of the cavity's endplates and excites the TE_{01} waveguide mode. The present coupling endplate connects to X-band waveguide feeds. The cylindrical TE_{01} waveguide mode has an easily calculable wavelength that depends on waveguide diameter and the permittivity of the medium inside the waveguide (7.1). Power propagates in this mode from the coupling endplate to the other endplate and reflects back. At certain frequencies the cavity length will be an integer number p of half-wavelengths, resulting in constructive interference. The integer p is called the axial mode number. At these frequencies a standing wave is set up, and a resonance condition occurs. For the NIST cavity, the first TE_{01} resonant mode occurs near 6.1 GHz, and the first TE_{02} resonant mode occurs near 11.2 GHz.

The NIST cavity resonator is ideally suited for accurate permittivity measurements because the resonator's quality factor Q is very high (60 000 to 80 000). Q is sometimes defined as the number of radians a wave propagates before the wave's power decays to e^{-1} , or approximately 37%, of its initial power. For a high- Q resonator this means that the excitation wave reflects off the cavity endplates and travels the length of the cavity several times before the wave's power is significantly reduced. For example, if $p = 25$ and the Q of the $TE_{01(25)}$ mode is 80 000, then an excitation wave travels the length of the cavity approximately $Q/(25\pi) \approx 1000$ times before decaying to e^{-1} of its initial power. If this high- Q cavity contains a sample, the excitation wave also passes through the material

more than 1000 times. In effect this increases the effective thickness of the sample and makes cavity resonator measurements sensitive to small changes in ϵ'_R and $\tan\delta$.

1.3 Methods of Measurement

To determine the permittivity of a material, two measurement methods can be used, as shown graphically in Fig. 1.1. One method involves reducing the sample-loaded cavity's length so that it resonates at the same frequency as the empty cavity. The change in length is a measure of permittivity, and the change in Q is a measure of the sample's dielectric loss (Sec. 2.2.3). The second method holds the cavity length fixed. The change in resonant frequency determines ϵ'_R , and the change in Q determines dielectric loss.

The theory from which ϵ'_R is calculated assumes that the cavity and sample have zero loss (infinite Q), so field orthogonality can be assumed. In practice, this assumption causes little problem because of the cavity's low surface impedance and because dielectric samples are usually low loss. The cavity resonator is so sensitive, however, that it is possible to determine and make corrections to frequency measurements that improve ϵ'_R measurement accuracy. These corrections are discussed in Chapter 4. In a sense, the concept of having two methods to measure the real part of permittivity is artificial. All that is needed to calculate ϵ'_R are cavity length and diameter, sample thickness and axial mode number. The empty cavity needs to be measured only because we have to determine the cavity's dimensions, where cavity length and diameter are calculated from the frequencies and axial mode numbers p of the TE_{01p} resonance spectrum.

The fixed-frequency method is more accurate than the fixed-length method for calculating dielectric loss. Cavity Q is frequency dependent. Skin depth (resistive losses) of the cavity walls decrease and coupling port losses increase as frequency increases. Presently, the measured empty-cavity Q is used as the reference that determines the cavity's conductor and port losses. When we measure the Q of the empty cavity, we determine the cavity loss at that resonant frequency. When a sample is inserted the resonant frequency is lowered. If we use the Q of the empty cavity resonating at a different frequency as our reference, our loss tangent calculation will be in error. This is especially true when the sample is very low loss ($\tan\delta < 0.0003$), because small differences in cavity Q become important. When sample loss is greater than $\tan\delta > 0.0008$ the differences between empty-cavity losses and the cavity losses of the sample-loaded cavity are small enough that the fixed-length method can be used with nearly the same accuracy as the fixed-frequency method.

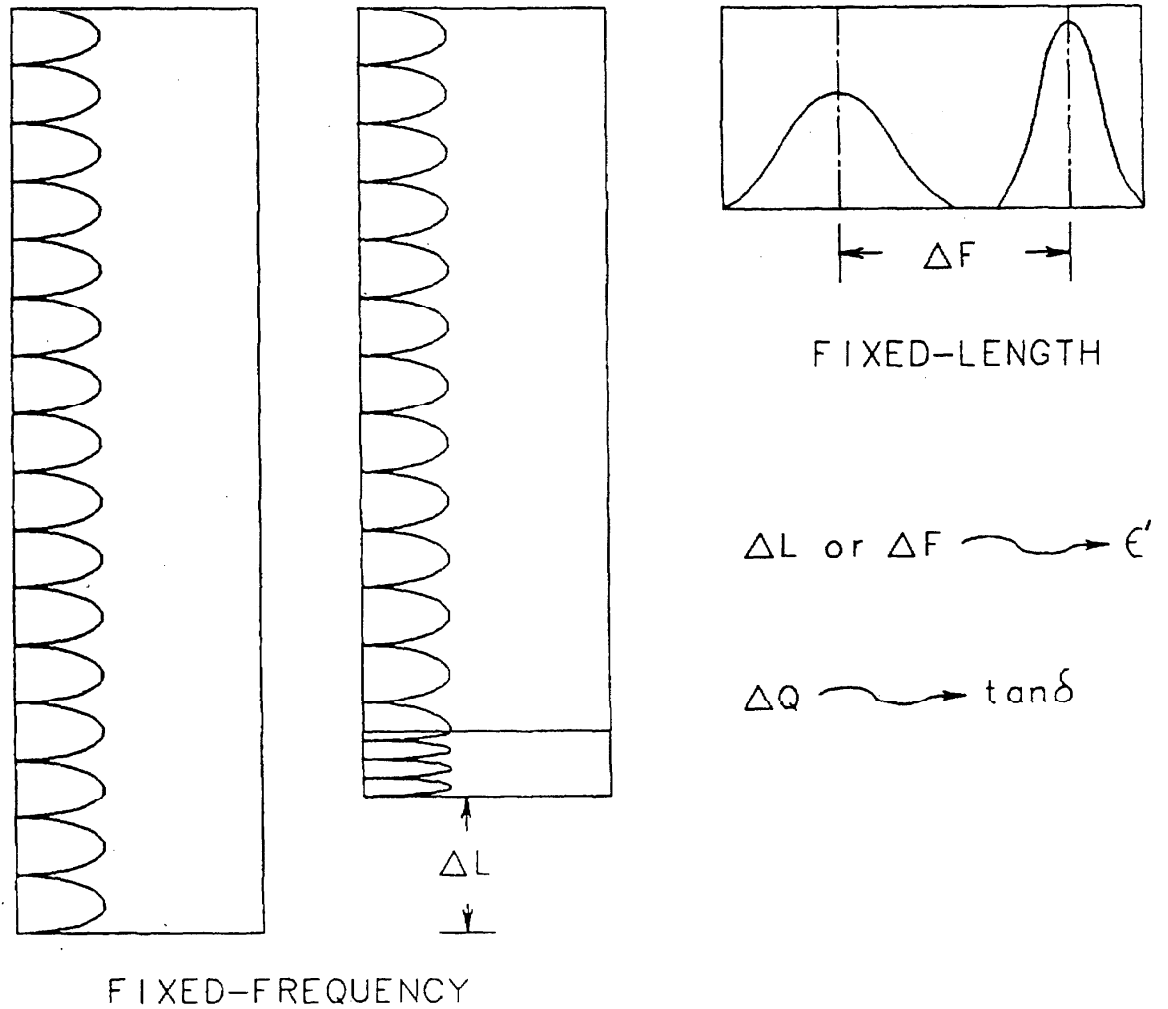


Figure 1.1: The fixed-frequency/variable-length and fixed-length/variable-frequency permittivity measurement methods rely on the empty-cavity dimensions and Q as a reference.

Chapter 2

Theory of Microwave Resonators

2.1 Introduction

Resonators create, filter, and select frequencies in oscillators, amplifiers, and tuners. Recent advances in the miniaturization of microwave circuits have spawned the development of low-loss, temperature-stable dielectric resonators. These resonators replace waveguide filters in many communication systems where microstrip and stripline resonators cannot be used due to their intrinsic high loss. Dielectric resonators are also very important fixtures for measuring the electrical properties of low-loss solids in the microwave region. An important resonator circuit at microwave frequencies is the closed cylindrical cavity operated in a resonant TE_{01p} mode. The magnetic field for this mode is sketched in Fig. 2.1. For a distant observer this mode appears as a magnetic dipole, and for this reason some authors call it a “magnetic dipole mode.” The electric field lines are simply circles concentric with the axis of the cylinder (Fig. 2.2).

For a TE mode in a waveguide, the electric field is everywhere *transverse* to the propagation direction or to the axial z -direction of the guide. Coordinate systems for circular and rectangular waveguides are sketched in Figs. 2.3 and 2.4, respectively. For a TE_{mn} mode in a rectangular guide, m simply represents the number of antinodes occurring in the electric field E in the x -direction and n represents the number of antinodes in E in the y -direction. For a TE_{mnp} mode in a finite circular guide, m represents the number of antinodes occurring in E in the ϕ -direction, n the number of antinodes in the radial r -direction, and p the number in the z -direction. The same notation is used for transverse magnetic or TM_{mn} modes, with the m and n indices representing the antinodes occurring

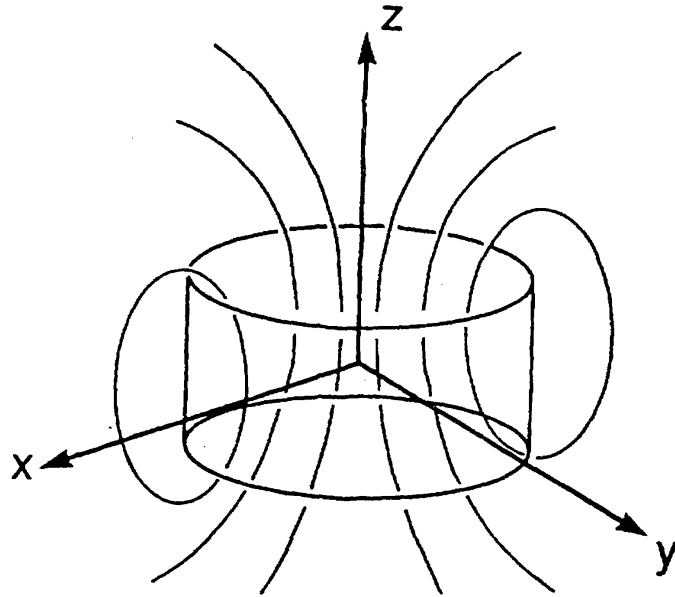


Figure 2.1: Magnetic field lines of the resonant mode TE_{01p} in an isolated dielectric resonator.

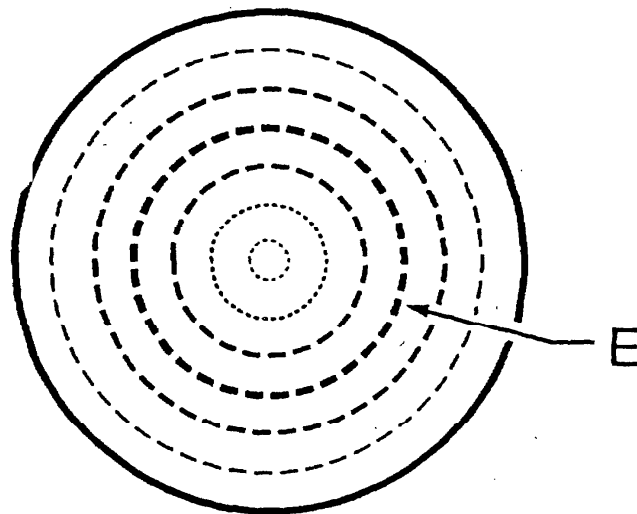


Figure 2.2: Electric field distribution in equatorial plane for TE_{01p} mode. Breadth of dashed lines is proportional to the electric field.

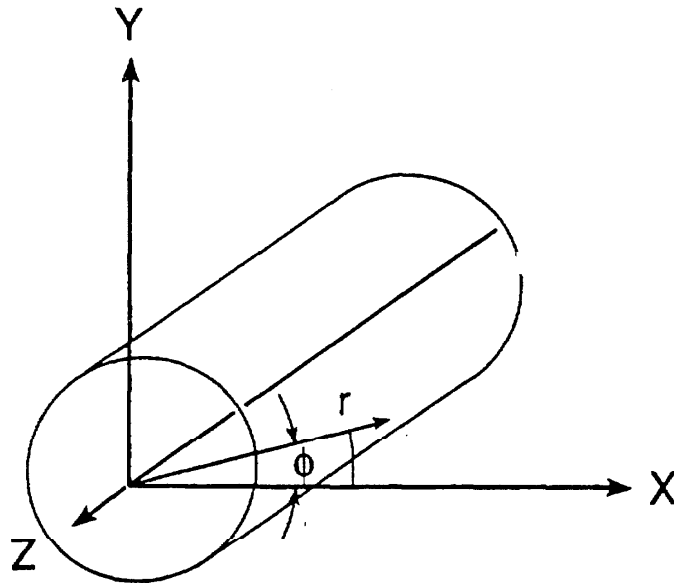


Figure 2.3: Cylindrical coordinate system for waveguides.

in the magnetic field H for a given coordinate direction. Generally, the choice of cavity to be used in dielectric metrology is influenced by the shape and size of the sample to be measured. Insertion and removal of the sample from the cavity should be convenient and not significantly alter the Q [4] of the air-filled cavity or increase the radiation from the cavity. Hollow circular cylindrical waveguide resonators, which are terminated by two short-circuit endplates, are the most commonly used for dielectric measurements of low-loss solids. The reasons for this are the relative ease of fabrication and the very high Q -factors and accompanying narrow bandwidths that can be obtained with this fixture in the microwave region [1]. If the cavity is made from helically wound waveguide, it also acts as a mode filter in that all other waveguide modes except the TE_{01p} mode [3] are greatly attenuated.

In general, microwave energy is coupled to the cavity through transmission line probes, as illustrated in Fig. 2.5. As in resonant transmission lines where resonance occurs at many frequencies, the hollow cylindrical waveguide has many resonance frequencies and associated field distributions or modes. The field of the mode with the lowest or dominant resonant frequency is termed the dominant mode of that fixture. Only enough energy is provided by the input excitation to match cavity losses. Once the internal fields at a resonant frequency are determined, an equivalent lumped parameter circuit of the fixture may be ascertained. In addition, the internal power dissipation, stored energies, and energy flow out of the cavity can be determined.

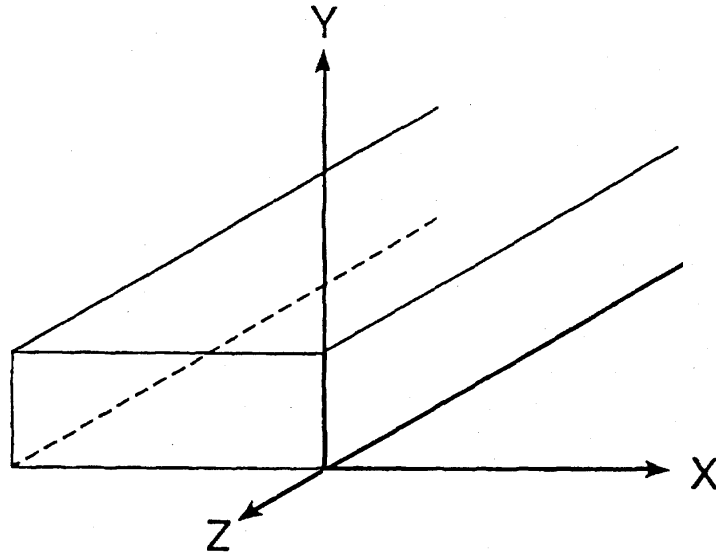


Figure 2.4: Rectangular coordinate system for waveguides.

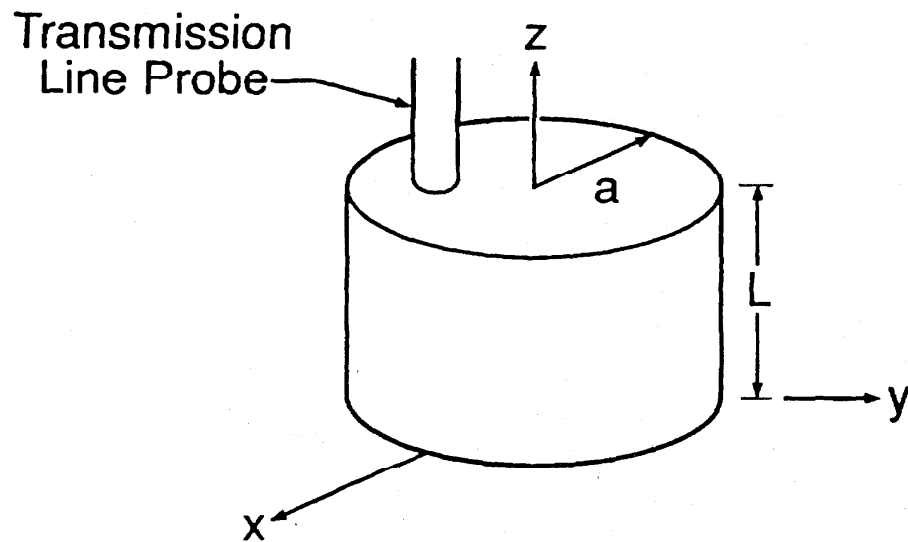


Figure 2.5: Circular cylindrical hollow cavity of length L and radius a .

As expected, the different microwave modes that can be set up in a hollow right circular cylinder depend on both the dimension of the cylinder and the microwave frequency. A hollow right circular cylindrical cavity is supportive of both TE - and TM -mode field structures. For steady-state $\exp(j\omega t)$ field time dependence, the equations for the axial components of the electric and magnetic field interior to a waveguide of arbitrary cross-section and filled with material of complex permittivity ϵ^* and complex permeability μ^* satisfy the homogeneous Helmholtz equation

$$\nabla^2 \begin{pmatrix} E_z \\ H_z \end{pmatrix} + k^2 \begin{pmatrix} E_z \\ H_z \end{pmatrix} = 0, \quad (2.1)$$

where z is the direction of propagation and E_z, H_z are taken for TM or TE modes, respectively. The differential operator ∇^2 is the Laplacian, which can be expressed in any curvilinear coordinate system by evaluation of the metric. All other field components can be expressed in terms of E_z and H_z . In general, any field in a hollow waveguide, however complicated, may be represented by a combination of TE and TM modes. The wavenumbers

$$k^2 = k_c^2 - \gamma^2 = \omega^2 \epsilon^* \mu^*, \quad (2.2)$$

are the characteristic eigenvalues of (2.1). In other words, to each value of k^2 there will correspond a function E_z or H_z (for TM or TE modes) which is a *characteristic* function from which may be derived the other components of the fields. In general ϵ^* and μ^* in (2.2) can be complex dyadics, but for this discussion the material is assumed isotropic and each dyadic becomes a single complex value multiplied by the unit dyadic. As defined in (1.7), the conductivity of the medium is contained in the imaginary part of the dielectric constant, $\epsilon^* = \epsilon' - j\epsilon''$. For a plane-wave field $H_z(z) = H_z e^{-\gamma z}$ in which the electric field is entirely transverse (TE waves),

$$\gamma^2 = \left(\frac{2\pi}{\lambda}\right)^2 = k_c^2 - \omega^2 \epsilon^* \mu^*, \quad (2.3)$$

where the value of λ is the wavelength of a plane TE wave in the medium that fills the hollow pipe and k_c is the cutoff wavenumber of the TE mode. In order to have propagation down the pipe, γ^2 must be negative. The quantity k_c^2 is always real because, for a circular waveguide of radius a ,

$$(k_c)_{mn} = \frac{t'_{mn}}{a}, \quad (2.4)$$

where t'_{mn} is the n th root of the first derivative of the Bessel function $J'_m(k_c a) = 0$ for a TE_{mn} mode. When γ^2 is positive or $\Re(\omega^2 \epsilon^* \mu^*) \leq k_c^2$ there is no propagation

of energy through the waveguide. This condition is termed cutoff. The cutoff or critical frequency is given by

$$\omega_c^2 = \Re \left\{ \frac{1}{\epsilon^* \mu^*} \right\} k_c^2. \quad (2.5)$$

Substituting (2.4) into (2.5) gives the cutoff frequency for TE_{mn} modes in a circular waveguide:

$$(f_c)_{TE_{mn}}^2 = \left(\frac{1}{2\pi} \right)^2 \Re \left\{ \frac{1}{\epsilon^* \mu^*} \right\} \left(\frac{t'_{mn}}{a} \right)^2, \quad (2.6)$$

or

$$(\lambda_c)_{TE_{mn}} = \frac{c}{(f_c)_{TE_{mn}}} = \frac{c \cdot 2\pi a \cdot \sqrt{\Re(\epsilon^* \mu^*)}}{t'_{mn}}, \quad (2.7)$$

where c is the velocity of light in the medium.

Similarly, for a TM_{mn} mode in circular waveguide,

$$(k_c)_{mn} = \frac{t_{mn}}{a}, \quad (2.8)$$

where t_{mn} is the n th root of $J_m(k_c a) = 0$.

As before, the cutoff wavelength for TM_{mn} modes is given by:

$$(\lambda_c)_{TM_{mn}} = \frac{c \cdot 2\pi a \cdot \sqrt{\Re(\epsilon^* \mu^*)}}{t_{mn}} \quad (2.9)$$

If λ_g is the wavelength in the guide and k_g the guide wavenumber, then

$$\gamma = j k_g = j \frac{2\pi}{\lambda_g}, \quad (2.10)$$

or, from (2.3),

$$\left(\frac{2\pi}{\lambda_g} \right)^2 = \left(\frac{2\pi}{\lambda} \right)^2 - \left(\frac{2\pi}{\lambda_c} \right)^2, \quad (2.11)$$

from which we derive the axial wavelength in the waveguide,

$$\lambda_g = \frac{\lambda}{\sqrt{1 - \left(\frac{\lambda}{\lambda_c} \right)^2}}, \quad (2.12)$$

where λ is the wavelength of a uniform plane wave in the medium:

$$\begin{aligned}\lambda &= \frac{\lambda_0}{\sqrt{\frac{\Re\{\epsilon^*\mu^*\}}{\epsilon_0\mu_0}}} \\ &\approx \frac{\lambda_0}{\sqrt{\frac{\epsilon'\mu'}{\epsilon_0\mu_0}}}\end{aligned}\quad (2.13)$$

where λ_0 is the free-space wavelength and we assume that $\Re\{\epsilon^*\mu^*\} \approx \epsilon'\mu'$. Equation (2.12) then becomes

$$\lambda_g = \frac{\lambda_0}{\sqrt{\frac{\epsilon'\mu'}{\epsilon_0\mu_0} - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} = \frac{\lambda_0}{\sqrt{\epsilon'_R\mu'_R - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}. \quad (2.14)$$

Equation (2.14) shows that when the free-space wavelength is much less than the cutoff wavelength, that waves propagate in the guide with a guide wavelength normalized by the square root of the product of relative permittivity and relative permeability of the material filling the guide. Of course, in the empty cavity situation, this means that the guide wavelength is essentially that of air. Similarly, when wavelengths are much greater than the cutoff wavelength for a given mode, the guide wavelength becomes pure imaginary or, stated physically, *nonpropagating* evanescent modes are set up.

2.1.1 Fields in a Right Circular Cylindrical Cavity

For the problem at hand, we are dealing with a cylindrical structure having two metallic end plates. In other words, it is bounded in axial extent and the associated internal fields may have variations in *both* the transverse and axial directions. In the case of *TE* waves, the wave equation (2.1) for H_z expressed in cylindrical coordinates becomes

$$\frac{1}{r} \frac{\partial}{\partial r} \left[r \frac{\partial}{\partial r} H_z \right] + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} H_z + \frac{\partial^2}{\partial z^2} H_z + k^2 H_z = 0, \quad (2.15)$$

where

$$k_c^2 = \gamma^2 + k^2, \quad (2.16)$$

and from (1.7)

$$\begin{aligned}k^2 &= \omega^2 \mu^* \epsilon^* = \omega^2 \mu^* (\epsilon' - j\epsilon'') \\ &= \omega^2 \mu^* \epsilon_0 \epsilon'_R (1 - j \tan \delta).\end{aligned}\quad (2.17)$$

The wavenumber is complex for lossy dielectric materials; that is,

$$k = k_{\text{real}} - jk_{\text{imaginary}} , \quad (2.18)$$

where it is easily shown that

$$k_{\text{real}}^2 = \frac{\omega^2 \mu \epsilon}{2} \left[\sqrt{\left(\frac{\sigma}{\omega \epsilon}\right)^2 + 1} + 1 \right] , \quad (2.19)$$

and

$$k_{\text{imaginary}}^2 = \frac{\omega^2 \mu \epsilon}{2} \left[\sqrt{\left(\frac{\sigma}{\omega \epsilon}\right)^2 + 1} - 1 \right] . \quad (2.20)$$

In (2.19) and (2.20) $\mu = \mu'_R \mu_0$ and $\epsilon = \epsilon'_R \epsilon_0$, where $\epsilon_0 \approx 8.854 \times 10^{-12}$ F/m and $\mu_0 = 4\pi \times 10^{-7}$ H/m. For lossless dielectrics, $\sigma = 0$ and $\epsilon''_R = 0$ with $k_{\text{real}} = \omega \sqrt{\mu \epsilon}$ and $k_{\text{imaginary}} = 0$.

The method of separation of variables [5] yields a solution to (2.15) of the form

$$H_z(r, \phi, z) = R(r)\Phi(\phi)Z(z) . \quad (2.21)$$

Substitution of (2.21) into (2.15) and then division by $H_z(r, \phi, z)$ yields

$$\frac{1}{rR} \frac{d}{dr} \left[r \frac{d}{dr} R \right] + \frac{1}{r^2 \Phi} \frac{d^2 \Phi}{d\phi^2} + \frac{1}{Z} \frac{d^2 Z}{dz^2} + k^2 = 0 . \quad (2.22)$$

The third term in (2.22) is explicitly independent of r and ϕ . It is also necessarily independent of z if (2.22) is to sum to zero for all (r, ϕ, z) . Therefore,

$$\frac{1}{Z} \frac{d^2 Z}{dz^2} = -\beta^2 , \quad (2.23)$$

where β is a constant. Substitution of (2.23) into (2.22) and multiplication by r^2 results in

$$r \frac{1}{R} \frac{d}{dr} \left[r \frac{d}{dr} R \right] + \frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} + (k^2 - \beta^2) r^2 = 0 . \quad (2.24)$$

The second term in (2.24) is a function of ϕ only, whereas the rest of the equation is a function of r only. By the same argument the azimuthal component of H_z obeys the relation,

$$\frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = -m^2 , \quad (2.25)$$

where m is a constant. Substitution of (2.25) into (2.24) and multiplication throughout by $R(r)$ yields

$$r \frac{d}{dr} \left[r \frac{d}{dr} R \right] + [(k^2 - \beta^2) r^2 - m^2] R = 0 , \quad (2.26)$$

where β is the axial waveguide propagation constant and k is the medium wavenumber. Now $k^2 - \beta^2$ represents the square of the transverse radial wavenumber k_c , where

$$k_c = \sqrt{k^2 - \beta^2}, \quad (2.27)$$

so (2.26) can be written

$$r \frac{d}{dr} \left[r \frac{dR}{dr} \right] + [(k_c r)^2 - m^2] R = 0. \quad (2.28)$$

The original Helmholtz equation (2.1) is now separated into three equations, each of which determines only one of the functions $R(r)$, $\Phi(\phi)$ or $Z(z)$. Equations (2.23) and (2.25), are harmonic equations, whose solutions are harmonic functions, or linear combinations of sines and cosines. Equation (2.28) is a Bessel equation of the m th order with independent solutions $J_m(k_c r)$ and $N_m(k_c r)$, where J_m and N_m represent Bessel and Neumann functions, respectively, of order m . Because $N_m(k_c r)$ is not finite at $r = 0$, the solution for $R(r)$ is

$$R(r) = J_m(k_c r). \quad (2.29)$$

The choice of the constants β and m , as well as the solutions for (2.23) and (2.25), depends on the physical geometry of the fixture, conditions at the boundaries, and the type of field to be supported by the fixture. For nonvanishing modes in a right circular cylindrical cavity, k_c can take on only *characteristic* or certain discrete values that correspond to different modes of propagation. In the case of the cylindrical cavity, the constant m must be an integer if the solution for E_z is to be single-valued in ϕ (periodic). The radical that defines the transverse radial wavenumber k_c calls for some comment. The branch of the square root is usually chosen such that $k_c \rightarrow k$ as $|\beta^2| \rightarrow 0$ and $k_c \rightarrow |\beta|$ as $|k^2| \rightarrow 0$.

2.1.2 TE Modes

For *TE* mode structure ($E_z = 0$) we need solve for only H_z . All other field components interior to the cavity are derived from H_z by Maxwell's equations. The complete solution for $H_z(r, \phi, z)$ is

$$H_z(r, \phi, z) = J_m(k_c r) [A \cos m\phi + B \sin m\phi] \sin \beta z, \quad (2.30)$$

which satisfies the boundary conditions at the cavity's cylindrical walls $r = a$ and at the cavity endplate $z = 0$. This includes the case where the cavity is filled with dielectric materials such that the transverse wavenumber k_c is complex. A and B are constants in (2.30) which determine the phase of the azimuthal field

Table 2.1: Zeroes t'_{mn} of the first derivative of the Bessel function of first kind and order m .

m	n		
	1	2	3
0	3.8317	7.0156	10.1735
1	1.8412	5.3314	8.5363
2	3.0542	6.7061	9.9695
3	4.2012	8.0152	11.3459

orientation relative to the coupling port(s). We are currently considering only an empty cylindrical cavity with $k = \omega\sqrt{\mu_0\epsilon_0}$ and arbitrary azimuthal phase. In order to enforce the boundary condition on E_ϕ , derivable from H_z from Maxwell's equations, we must have

$$E_\phi(a, \phi, z) = 0, \quad (2.31)$$

or

$$J'_m(k_c a) = 0, \quad (2.32)$$

where J'_m represents the first derivative of the m th-order Bessel function. If the zeroes of J'_m are denoted by t'_{mn} , where $n = 1, 2, 3, \dots$ represents the zero-crossing number, then k_c must be chosen to take only certain discrete values; that is,

$$k_c = \frac{t'_{mn}}{a}. \quad (2.33)$$

Table 2.1 gives some representative values for t'_{mn} .

The boundary condition on the azimuthal electric field at $z = L$ results in certain allowable values for the longitudinal propagation wavenumber:

$$\sin \beta L = 0, \quad (2.34)$$

or

$$\beta = p \frac{\pi}{L}, \quad (2.35)$$

where $p = 1, 2, 3, \dots$. The final result for the axial magnetic field of the TE mode within the cavity is:

$$H_z(TE_{mnp}) = J_m\left(\frac{t'_{mn} r}{a}\right) [A \cos m\phi + B \sin m\phi] \sin\left[\frac{p\pi z}{L}\right], \quad (2.36)$$

where m, n, p are integers describing the TE mode. An identical analysis can be performed to derive the axial electric field for the TM case

The characteristic equation used to find the resonant frequencies for TE modes in the empty cavity is given by Harrington [5] pp. 213–216 as,

$$f_0(TE_{mnp}) = \frac{c_{air}}{2\pi} \left[\left(\frac{t'_{mn}}{a} \right)^2 + \left(\frac{p\pi}{L} \right)^2 \right]^{\frac{1}{2}}, \quad (2.37)$$

where c_{air} is the speed of light under testing conditions in air. When $f = f_0(TE_{mnp})$, we have a TE -mode solution to Maxwell's equations. From (2.2), the resonant frequency for the cylindrical cavity filled with material of complex permittivity ϵ^* and real permeability μ' is

$$f_0(TE_{mnp}) = \frac{1}{2\pi\sqrt{\mu'\epsilon'(\tan^2\delta + 1)^{\frac{1}{2}}}} \left[\left(\frac{t'_{mn}}{a} \right)^2 + \left(\frac{p\pi}{L} \right)^2 \right]^{\frac{1}{2}}. \quad (2.38)$$

As expected, when the length or diameter of a cylindrical cavity increases, the resonant frequency decreases for any given mode.

The figure of merit for assessing the performance or quality of a cavity resonator is the quality factor Q which is a measure of energy stored in the fields inside the resonator compared to the energy loss or dissipation per cycle. The Q -factor is defined by

$$\begin{aligned} Q &= 2\pi \frac{\text{maximum energy stored during a cycle}}{\text{average energy dissipated per cycle}}, \\ &= \frac{2\pi W}{PT} = \frac{\omega_0 W}{P}, \end{aligned} \quad (2.39)$$

where W is stored energy, P is power dissipation, ω_0 is resonant radian frequency, and T is period = $\frac{2\pi}{\omega_0}$. The higher the axial TE -mode number for any given azimuthal and radial mode number, the greater the cavity quality factor, for any given cavity diameter or length. The quality factor Q will be discussed in more detail in connection with the evaluation of dielectric loss measurements of materials. Some examples of normalized cavity Q values are shown in Fig. 2.6 for some representative TM_{mnp} modes and in Fig. 2.7 for some TE_{0np} modes as a function of cavity diameter-to-length ratio. In Fig. 2.7, the optimal (highest) Q value is obtained for TE_{mnp} modes when the diameter of the cavity equals its length. Our 60-mm diameter resonator's length is roughly seven times longer than its diameter. Resonance Q -value is compromised somewhat by this length, but the cavity's helical windings yield better than 30-dB mode purity. The length of the resonator allows for the measurement of several TE_{01p} modes in the X-band frequency range.

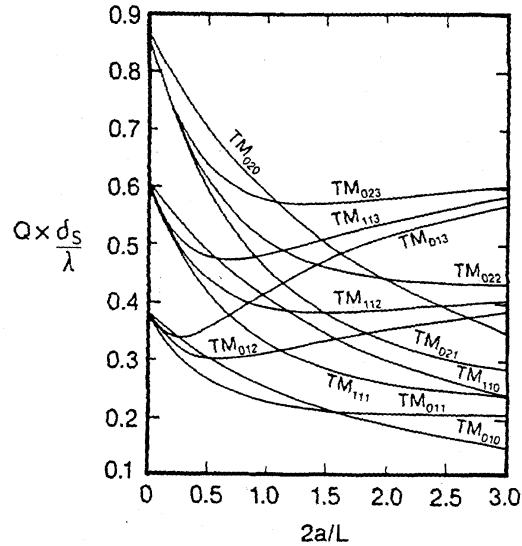


Figure 2.6: Normalized cavity Q versus diameter-to-length ratio for a right circular cylinder for TM_{mnp} modes. δ_s is the skin depth of the cavity wall material and λ is the free-space wavelength.

2.2 TE_{01p} Mode-Filter Cylindrical Cavity

One very useful cavity resonator for microwave dielectric property measurements on low-loss materials is constructed to filter all modes resulting from current other than that flowing circumferentially about the cavity wall. This yields a very high- Q cavity with a very pure TE -mode structure for precision electrical property measurements.

Consider, for example, the wall and endplate currents flowing in the cylindrical cavity for TE_{01p} and TM_{11p} modes as shown in Fig. 2.8. For the TM_{11p} mode there are both azimuthal and radial currents, whereas in the case of the TE_{01p} mode there are only azimuthal currents flowing in the wall and end plates. As Cook [2] notes, the presence of unwanted modes produces larger data scatter of loss tangent and permittivity with change in specimen length than when there is only one effective mode propagating, such as the TE_{01p} .

If the currents associated with the TM_{11p} mode can be interrupted while at the same time leaving those associated with the TE_{01p} mode unchanged, then signif-

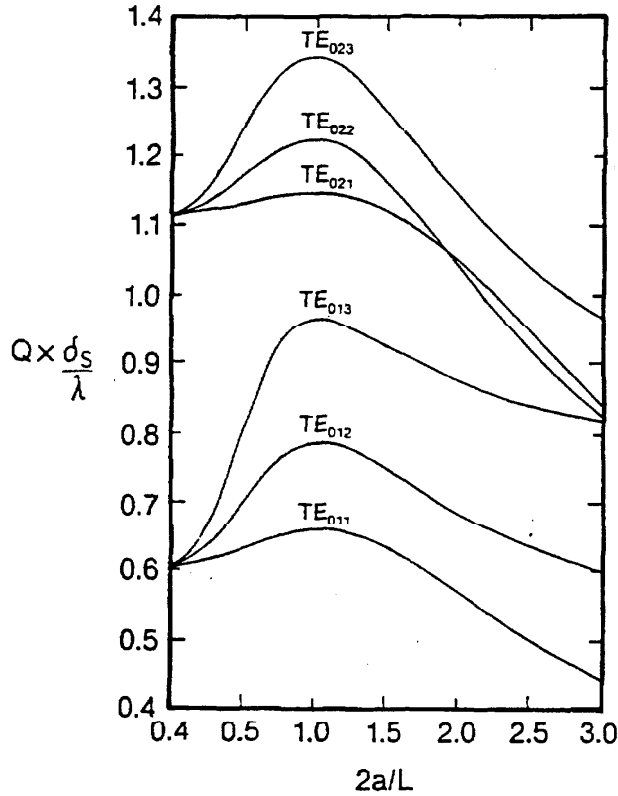


Figure 2.7: Normalized cavity Q versus diameter-to-length ratio for a right circular cylinder for TE_{0np} modes. δ_s is the skin depth of the cavity wall material and λ is the free-space wavelength.

icant attenuation of the TM_{11p} mode will take place. This can be accomplished by constructing a cavity wall in which the conductivity is discontinuous along the length of the cylinder but continuous around its circumference. This type of cavity becomes a mode filter. One approach to the construction of a mode-filter cavity is to construct a cylinder of annular copper rings electrically insulated from each other [2] to form a continuous cylinder. In this case, the currents flowing from ring to ring will be greatly impeded while those flowing around the rings are unimpeded.

Another approach is to make a helical waveguide by winding fine enameled wire into a precision cylindrical former (Fig. 2.9). This type of waveguide has been discussed in the literature by Morgan and Young [6]; Unger [7]; Young [8]; Cook

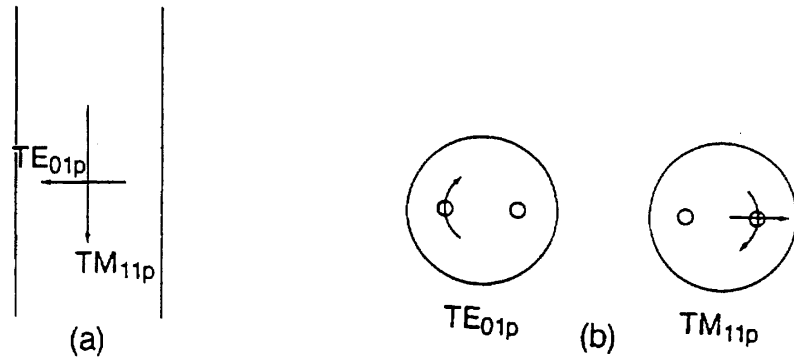


Figure 2.8: Currents flowing in (a) cylinder wall and (b) cavity end plates for TE_{01p} and TM_{11p} modes.

and Jones [9]; Cook [2]; Waldron [10] [11]; Waldron and Bowe [12]; Shimba [13]; Kazantsev, Kaznacheev and Meriakri [14]; Mikoshiba [15]; Piefke [16]; Waldron, Bowe, Wackrill and Wescott [17]; and Noda, Yamaguchi and Suzuki [18]. In a helical waveguide the modes whose wall currents follow the conducting helix possess attenuation constants which are essentially the same as for copper pipe. All other modes, however, have a high transmission loss. In other words, all modes other than the TE_{01p} have very large attenuation constants in a helical waveguide. The exact calculation of the attenuation constants for modes other than TE_{01p} depends on the helix pitch angle and the electrical properties of the thick lossy dielectric jacket surrounding the helix; the fields of those modes penetrate into the jacket and are attenuated, whereas that of the TE_{01p} mode does not and is therefore minimally attenuated. The helical waveguide, then, is the equivalent of an anisotropically conducting cylinder. Whether we use helically wound wire as opposed to annular copper rings is a question of convenience in manufacture.

Waldron [10], and Waldron and Bowe [12] have analyzed the azimuthal electric field dependence as a function of the radius of the helical winding, the complex permittivity of the jacket material, and the pitch angle of the helix. To a first-order approximation the TE_{01p} mode is independent of the pitch angle. Waldron and Bowe [12] also treat the effect of pitch for the TE_{01p} mode to the second order and find that the effect of a finite pitch does indeed cause a slight attenuation of the TE_{01p} mode. Cook [2] compares the attenuation of the TE_{01p} mode to other modes in a 50-mm diameter helix at 35 GHz; these results are given in Table 2.2. Note that the modes other than TE_{01p} have attenuation constants larger by several orders of magnitude than that for TE_{01p} . Piefke [16] demonstrates that

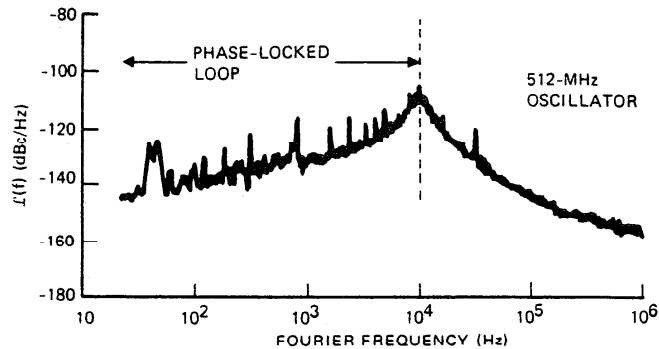


FIG. 11 Phase-locked-loop characteristics of the H.P. 8640B signal generator, showing the normalized phase-noise sideband power spectral density.

4. Measurements, Calculations, and Data Plots

The measurement sequence is automated except for the case where manual adjustments are required to maintain phase quadrature of the signals. After phase quadrature of the signals into the mixer is established, the IF attenuator is returned to the zero-decibel reference setting. This attenuator is set to a high value [assumed to be 50 dB in Eq. (65)] to prevent saturation of the spectrum analyzer during the calibration process.

The automated measurements are executed, and the direct measurement and data plot of $\mathcal{L}(f)$ is obtained in decibels (carrier) per hertz using the equation

$$\mathcal{L}(f) = - [\text{carrier power level} - (\text{noise power level} - 6 + 2.5 - 10 \log B - 3)]. \quad (69)$$

The noise power (dBm) is measured relative to the carrier-power level (dBm), and the remaining terms of the equation represent corrections that must be applied because of the type of measurement and the characteristics of the measurement equipment, as follows.

- (1) The measurement of noise sidebands with the signals in phase quadrature requires the -6 -dB correction that is noted in Eq. (69).
- (2) The nonlinearity of the spectrum analyzer's logarithmic IF amplifier results in compression of the noise peaks which, when average-detected, require the 2.5-dB correction.
- (3) The bandwidth correction is required because the spectrum analyzer measurements of random or white noise are a function of the particular bandwidth used in the measurement.

Figure 8.5 gives ϵ'_R results for these three samples. Samples 1 and 2 yield nearly identical results, while sample 3 has a noticeably lower permittivity. Figure 8.6 gives loss tangent results from these same three samples. All three samples have a noticeably different loss, and, surprisingly, the fused quartz sample has lower loss than the fused silica. We normally expect impurities to cause greater loss. The discontinuity in the results near 11.7 GHz are due to mode interference by TE_{02p} modes.

These results indicate that the permittivity and loss of high-purity, high homogeneity fused silica are consistent from manufacturer to manufacturer. However there are repeatable differences in permittivity and loss between the two fused silica samples. The fused-quartz sample has a noticeably lower permittivity and loss than the fused-silica samples. This could be because the fused-quartz sample is less dense and has more seed crystals and inclusions, or because of impurities. Also, the fixed-frequency and fixed-length results for sample 2 are nearly identical.

Uncertainty estimates for sample number 1 are given in Figs. 8.7 and 8.8.

8.2.3 Alumina

Eleven samples of 99.9% pure Al_2O_3 are under evaluation. Samples 8-11 were made separately from the first seven samples, and permittivity results for samples 1-7 are given in Figs. 8.9 and 8.10. A demonstrable difference between samples from the same batch can be seen. This is not due to measurement variations because results for Sample 6 are repeatable. Although further investigation is needed, the higher permittivity results at the lower frequencies are most likely due to our having used an effective cavity length that was too long for those frequencies (Sec. 4.1). Loss tangent for this high-purity alumina is the lowest of any material we have measured. Also, the fixed-frequency method gives stable and repeatable loss tangent results for most of the frequency range, while the fixed-length method varies much more. This is due to the limitations in the loss tangent model which assumes an equivalence between the empty cavity and a cavity containing a hypothetical lossless sample of the same permittivity as the real sample (Sec. 4.2). The negative loss tangent values near 11.5 GHz are due to the interfering TE_{021} mode. Uncertainty estimates for samples with $\epsilon'_R \approx 10$ are given in Chapter 9, so no uncertainty results are given here.

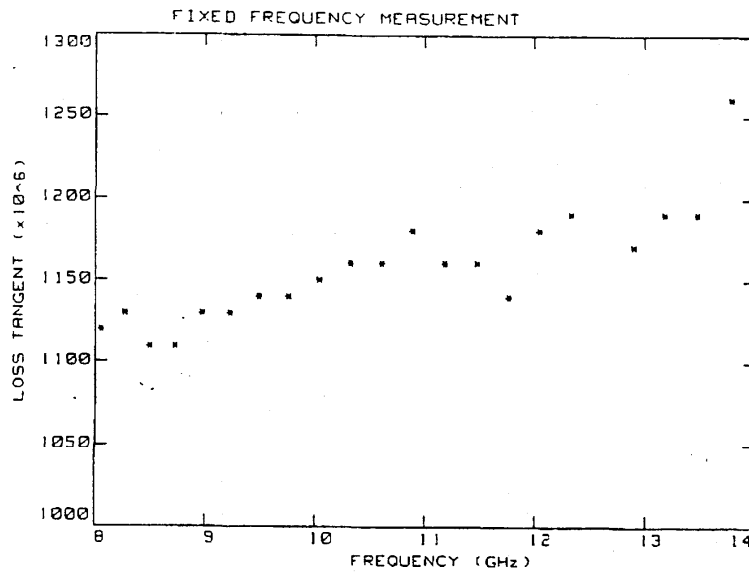


Figure 8.4: Loss tangent $\tan\delta$ of cross-linked polystyrene.

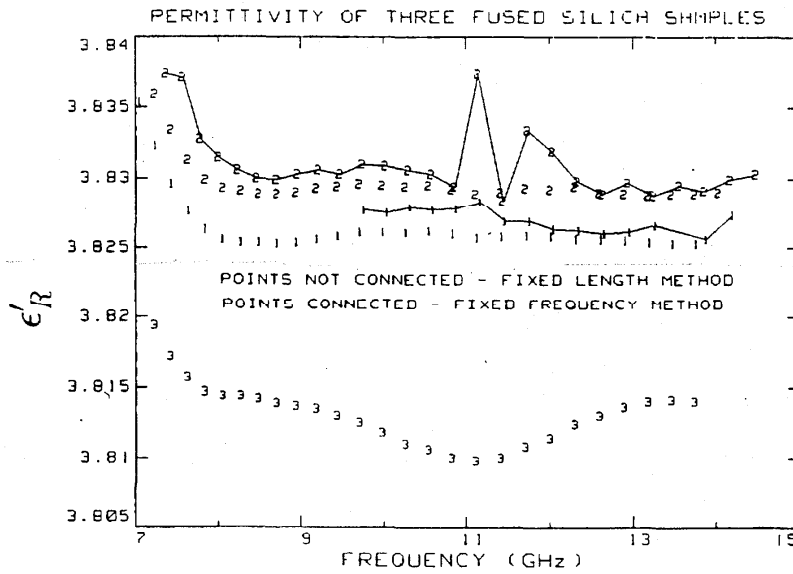


Figure 8.5: Permittivity ϵ'_R of two fused silica samples (1 and 2) and one fused quartz sample (3).

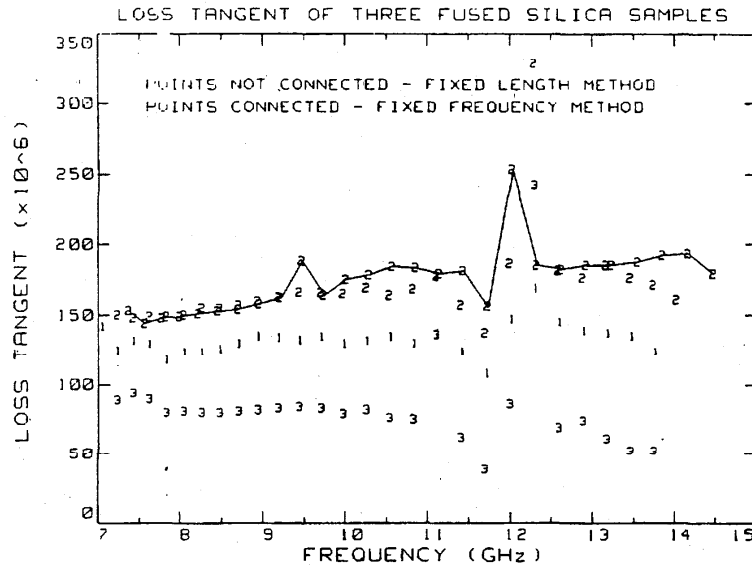


Figure 8.6: Loss tangent $\tan\delta$ of two fused silica samples (1 and 2) and one fused quartz sample (3). Note effect of interference by a TE_{02} mode between 11 and 13 GHz.

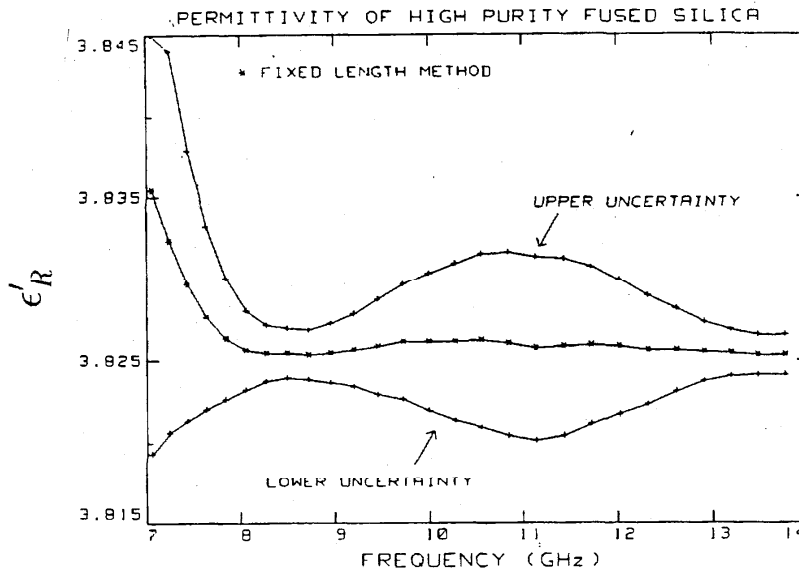


Figure 8.7: Permittivity ϵ'_R and estimated uncertainty of fused silica sample # 1.

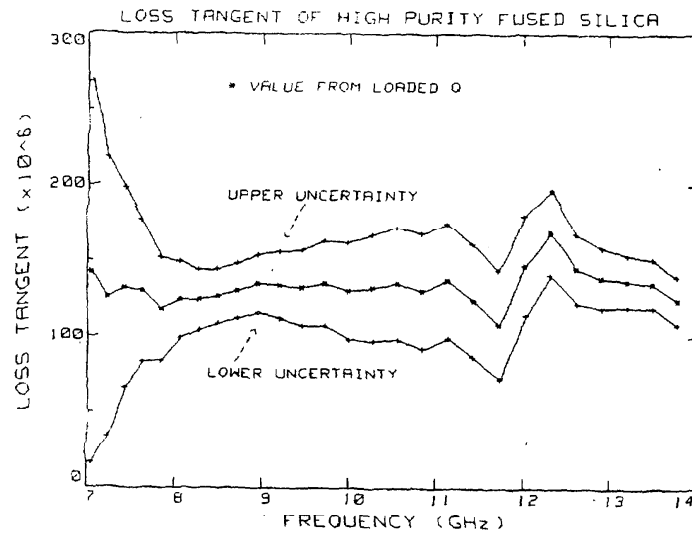


Figure 8.8: Loss tangent $\tan\delta$ and estimated uncertainty of fused silica sample # 1.

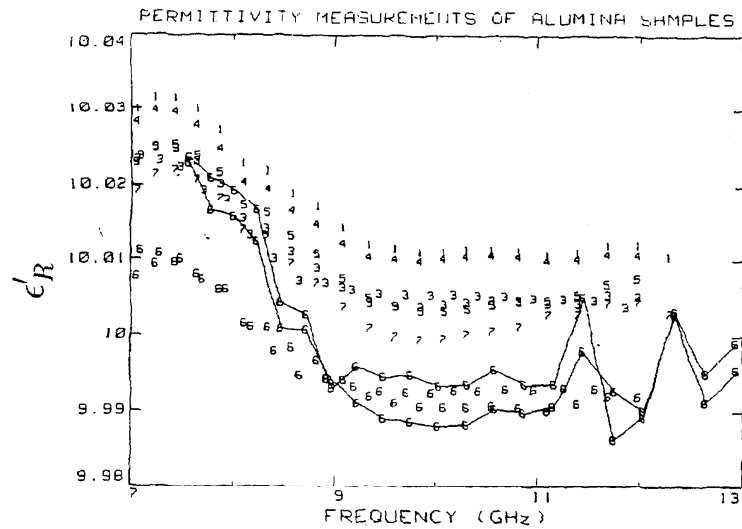


Figure 8.9: Measured permittivity ϵ'_R of several alumina samples from the same batch.

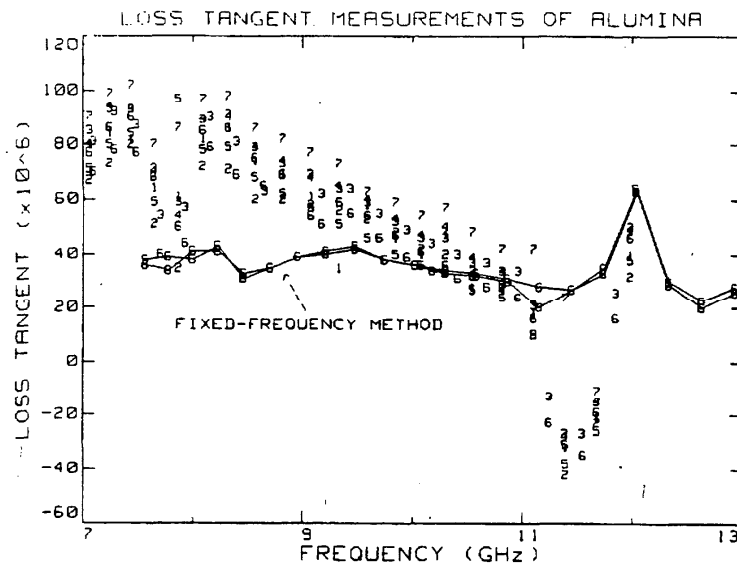


Figure 8.10: Measured loss tangent $\tan\delta$ of several alumina samples from the same batch.

Chapter 9

Uncertainty Analysis

This chapter describes how uncertainty estimates are obtained from the computer program CAVITYPROG. Typical uncertainty estimates are given for a 10-mm thick sample with a relative permittivity value of $10 - j0.001$. CAVITYPROG generates uncertainty tables for ϵ'_R and $\tan\delta$ due to individual parameter uncertainties. The uncertainties in cavity diameter and length, resonant frequency, micrometer reading, and resonance bandwidth were found in Chapter 4 and are summarized in Table 9.1.

The ϵ'_R and $\tan\delta$ results for the samples presented in Chapter 8 show estimated total uncertainty. This chapter presents estimated ϵ'_R and $\tan\delta$ uncertainties due to uncertainties in individual parameters. Because there can be any number of combinations of sample permittivity, loss, thickness, cavity length and resonant frequency, for the sake of brevity, we present results for a hypothetical sample with a relative permittivity of $10 - j0.001$. Uncertainty estimates are given for several frequencies in the X-band range. The uncertainty estimates are similar

Table 9.1: Estimated parameter uncertainties for the NIST resonator.

Cavity diameter	a	$\pm 2 \mu\text{m}$
Cavity length	L_r	$\pm 9 \mu\text{m}$
Micrometer reading	ΔL_r	$\pm 1 \mu\text{m}$
Sample thickness	b	$\pm 3 \mu\text{m}$
Resonance frequency	F_{sample}	$\pm 1000 \text{ Hz}$
Speed of light	c_{air}	$\pm 1000 \text{ m/s}$
Empty cavity Q	Q_0	$\pm 3\%$
With-sample Q	Q_s	$\pm 3\%$

for other permittivities, and must be evaluated on a case by case basis. For example, the measurement uncertainties for a thin sample with low dielectric constant can be much larger than one might expect (on the order of 20%), whereas a thicker sample or one with higher permittivity would have much smaller estimated measurement uncertainty.

9.1 Permittivity Uncertainty ($\Delta\epsilon'_R$)

To find uncertainties in ϵ'_R , CAVITYPROG calculates the change in ϵ'_R when one parameter is changed by its given uncertainty, while all other parameters are held fixed. This is repeated for each parameter to find the estimated uncertainty $\Delta\epsilon'_R$ due to individual parameters. CAVITYPROG then finds worst-case total uncertainty by summing the individual uncertainties. This variational method is equivalent to taking the total differential of ϵ'_R with respect to all parameters.

Figure 9.1 shows the estimated uncertainty $\Delta\epsilon'_R$ due to the individual measured parameter uncertainties given in Table 9.1.

Cavity diameter and length uncertainties are the greatest sources of uncertainty for ϵ'_R permittivity calculation. As shown previously in Fig. 4.3, the calculated cavity diameter changed less than 0.002 mm when the tuner endplate varied the cavity length by 24 mm. This result is corroborated at all frequencies in X-band as shown in Fig. 4.7 of Sec. 4.1.1 where the cavity dimensions were calculated from different subsets of the 7–14 GHz mode spectrum.

When we measure samples using the fixed-frequency method, the cavity length can be shortened by its maximum 24 mm range. From dimensional evaluation experiments described in Chapter 4, the uncertainty in diameter is less than ± 0.002 mm and the uncertainty in length is no greater than ± 0.009 mm. The overall uncertainty in cavity length is taken as the root-sum-square of the cavity length and the micrometer reading. These uncertainties are primarily due to endplate-travel accuracy (Sec. 4.1.2) and the agreement of calculated cavity dimensions from subsets of the 7–14 GHz resonance spectra (Sec. 4.1.1).

Sample thickness uncertainty is another significant error source that tends to have a maximum contribution at the frequencies at which the other error sources are minimized, as shown in Fig. 9.1. If the sample is thin < 2 mm, this estimated error can become very significant especially if the sample surfaces are rough.

The speed of light in Boulder, Colorado is approximately $c_{air} = 2.99709 \times 10^8$ m/s at 10 GHz, which corresponds to a relative velocity factor of 0.99972. At standard temperature and pressure (23°C, 1.013×10^5 Pa (1013 mbar)) and 50% relative humidity the speed of light is approximately 2.99695×10^8 m/s with

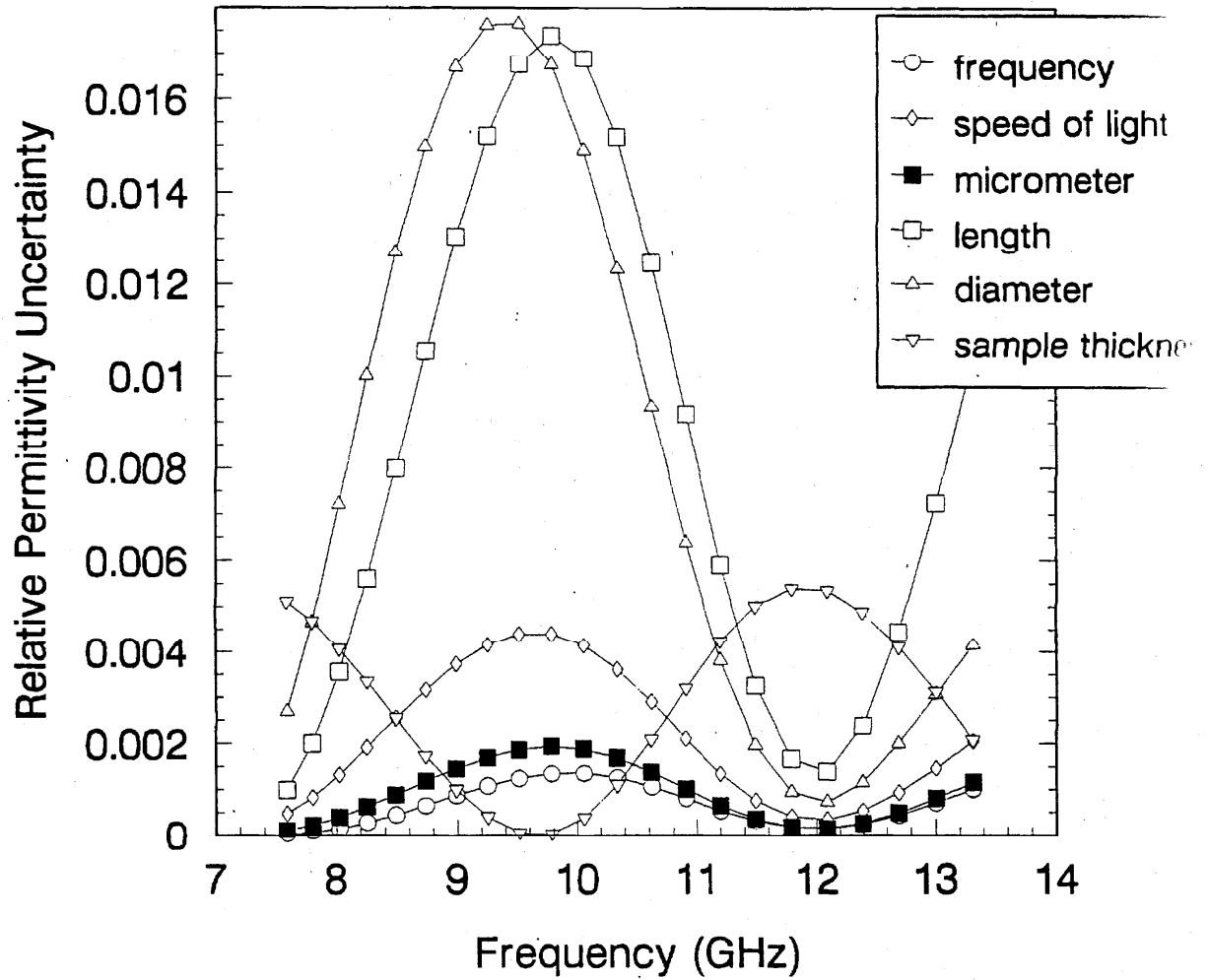


Figure 9.1: Permittivity uncertainty ($\Delta\epsilon'_R$) versus frequency due to the uncertainties in various measured parameters for a 10-mm thick sample with $\epsilon'_R = 10 - j0.001$.

relative velocity factor 0.99967. In the calculations for ϵ'_R , erroneously setting the speed of light in the air-filled portion of the cavity to the speed of light in a vacuum gives a significant upward bias to ϵ'_R results. If we use the proper value for the speed of light in calculations, the residual uncertainty in c_{air} is approximately ± 300 m/s. To accommodate for miscellaneous atmospheric disturbances we have specified the uncertainty in the speed of light to be approximately three times greater ($\approx \pm 1000$ m/s). Figure 9.1 shows that this uncertainty in the speed of light can indeed contribute an appreciable error to ϵ'_R .

Resonant frequency can be very accurately determined to within a few hundred hertz. However, if the resonant frequency is set and observed for one hour the resonant frequency drifts no more than 1 kHz due to changing cavity length caused by settling of the endplate drive micrometer or because of slight changes in cavity temperature. The reading of the sensing micrometer changes less than 0.001 mm to reflect this change in length. Because this problem occurs when the resonator is left alone for a period of time that is approximately the same as the time the computer takes to measure the X-band resonance spectrum, we use a 1 kHz uncertainty for the empty and sample-filled resonant frequencies. Fig. 9.1 shows the uncertainty in ϵ'_R due to a resonant frequency uncertainty of ± 1 kHz. There are cases in which the Q of a resonance can be low (≈ 2000) because of a sample's extreme thickness or lossiness. In this case, resonant frequency and Q can be difficult to define because of asymmetries that may occur in the resonance response. When this is the case, one must use special precautions in defining the uncertainties in resonance frequency and Q because systematic errors having to do with the resonator's frequency dependent characteristics may be occurring.

9.2 Loss Tangent Uncertainty

Loss tangent uncertainty is calculated from the total differential of the loss tangent equation given by Cook [2]. We have verified that this method is identical to the variational method used in calculating $\Delta\epsilon'_R$ in which the uncertainties are added to individual parameters, one at a time, to find the change in $\tan\delta$. The resulting individual estimates for $\Delta\tan\delta$ are added together to get a worst-case uncertainty estimate. The total differential method can be used to estimate $\Delta\tan\delta$ because we have an explicit equation for $\tan\delta$. The equation is given by

$$\tan\delta = \frac{p(2b-s) + (1/\epsilon'_R)[2(L_r-b) - q]}{p(2b-s)} \left(\frac{1}{Q_d} - \frac{1}{Q'} \right), \quad (9.1)$$

where

$$p = \frac{\sin^2 \beta_0 (L_r - b)}{\sin^2 \beta_1 b}, \quad (9.2)$$

$$q = \frac{\sin 2\beta_0 (L_r - b)}{\beta_0}, \quad (9.3)$$

$$s = \frac{\sin 2\beta_1 b}{\beta_1}. \quad (9.4)$$

The total differential is found from

$$\begin{aligned} \partial \tan \delta = & \partial L_r \left\{ \frac{\partial \tan \delta}{\partial L_r} + \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial L_r} + \frac{\partial \tan \delta}{\partial q} \frac{\partial q}{\partial L_r} \right\} \\ & + \partial b \left\{ \frac{\partial \tan \delta}{\partial b} + \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial b} + \frac{\partial \tan \delta}{\partial q} \frac{\partial q}{\partial b} + \frac{\partial \tan \delta}{\partial s} \frac{\partial s}{\partial b} \right\} \\ & + \partial \epsilon'_R \left\{ \frac{\partial \tan \delta}{\partial \epsilon'_R} + \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial \beta_\epsilon} \frac{\partial \beta_\epsilon}{\partial \epsilon'_R} + \frac{\partial \tan \delta}{\partial s} \frac{\partial s}{\partial \beta_\epsilon} \frac{\partial \beta_\epsilon}{\partial \epsilon'_R} \right\} \\ & + \partial Q_{\text{sample}} \left\{ \frac{\partial \tan \delta}{\partial Q_{\text{sample}}} \right\} \\ & + \partial Q_{\text{empty}} \left\{ \frac{\partial \tan \delta}{\partial Q_{\text{empty}}} \right\} \\ & + \partial c_{\text{air}} \left\{ \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial \beta_0} \frac{\partial \beta_0}{\partial c_{\text{air}}} + \frac{\partial \tan \delta}{\partial q} \frac{\partial q}{\partial \beta_0} \frac{\partial \beta_0}{\partial c_{\text{air}}} \right\} \\ & + \partial a \left\{ \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial \beta_0} \frac{\partial \beta_0}{\partial a} + \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial \beta_\epsilon} \frac{\partial \beta_\epsilon}{\partial a} + \frac{\partial \tan \delta}{\partial q} \frac{\partial q}{\partial \beta_0} \frac{\partial \beta_0}{\partial a} + \frac{\partial \tan \delta}{\partial s} \frac{\partial s}{\partial \beta_\epsilon} \frac{\partial \beta_\epsilon}{\partial a} \right\} \\ & + \partial f \left\{ \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial \beta_0} \frac{\partial \beta_0}{\partial f} + \frac{\partial \tan \delta}{\partial p} \frac{\partial p}{\partial \beta_\epsilon} \frac{\partial \beta_\epsilon}{\partial f} + \frac{\partial \tan \delta}{\partial q} \frac{\partial q}{\partial \beta_0} \frac{\partial \beta_0}{\partial f} + \frac{\partial \tan \delta}{\partial s} \frac{\partial s}{\partial \beta_\epsilon} \frac{\partial \beta_\epsilon}{\partial f} \right\} \end{aligned} \quad (9.5)$$

Uncertainty is found by replacing the partials ∂L_r , ∂b , $\partial \epsilon'_R$, $\partial Q_{\text{sample}}$, $\partial Q_{\text{empty}}$, ∂c_{air} , ∂a and ∂f with their respective estimated uncertainties ΔL_r , Δb , $\Delta \epsilon'_R$, ΔQ_{sample} , ΔQ_{empty} , Δc_{air} , Δa and Δf . The partial derivatives are given by

$$\frac{\partial \tan \delta}{\partial L_r} = \frac{2(Q_{\text{empty}} - Q_{\text{sample}})}{\epsilon'_R p Q_{\text{empty}} Q_{\text{sample}} (2b - s)}, \quad (9.6)$$

$$\frac{\partial \tan \delta}{\partial b} = \frac{-2(Q_{\text{empty}} - Q_{\text{sample}})(2L_r - q - s)}{\epsilon'_R p Q_{\text{empty}} Q_{\text{sample}} (2b - s)^2}, \quad (9.7)$$

$$\frac{\partial \tan \delta}{\partial \epsilon'_R} = \frac{(Q_{\text{empty}} - Q_{\text{sample}})(2b - 2L_r + q)}{(\epsilon'_R)^2 p Q_{\text{empty}} Q_{\text{sample}} (2b - s)}, \quad (9.8)$$

$$\frac{\partial \tan \delta}{\partial Q_{\text{sample}}} = -\frac{2b(\epsilon'_R p - 1) - \epsilon'_R p s + 2L_r - q}{\epsilon'_R p Q_{\text{sample}}^2 (2b - s)}, \quad (9.9)$$

$$\frac{\partial \tan \delta}{\partial Q_{\text{empty}}} = \frac{2b(\epsilon'_R p - 1) - \epsilon'_R p s + 2L_r - q}{\epsilon'_R p Q_{\text{empty}}^2 (2b - s)}. \quad (9.10)$$

The partial derivatives of p , q and s are given by

$$\frac{\partial p}{\partial \beta_0} = \frac{2(b - L_r) \sin(\beta_0(b - L_r)) \cos(\beta_0(b - L_r))}{\sin^2(b\beta_\epsilon)}, \quad (9.11)$$

$$\frac{\partial p}{\partial \beta_\epsilon} = -\frac{2b \cos(b\beta_\epsilon) \sin^2(\beta_0(b - L_r))}{\sin^3(b\beta_\epsilon)}, \quad (9.12)$$

$$\begin{aligned} \frac{\partial p}{\partial b} &= \frac{2\beta_0 \sin(\beta_0(b - L_r)) \cos(\beta_0(b - L_r))}{\sin^2(b\beta_\epsilon)} \\ &\quad - \frac{2\beta_\epsilon \cos(b\beta_\epsilon) \sin^2(\beta_0(b - L_r))}{\sin^3(b\beta_\epsilon)}, \end{aligned} \quad (9.13)$$

$$\frac{\partial p}{\partial L_r} = -\frac{2\beta_0 \sin(\beta_0(b - L_r)) \cos(\beta_0(b - L_r))}{\sin^2(b\beta_\epsilon)}, \quad (9.14)$$

$$\frac{\partial q}{\partial \beta_0} = \frac{\sin(2\beta_0(b - L_r))}{\beta_0^2} - \frac{2(b - L_r) \cos(2\beta_0(b - L_r))}{\beta_0}, \quad (9.15)$$

$$\frac{\partial q}{\partial b} = -2 \cos(2\beta_0(b - L_r)), \quad (9.16)$$

$$\frac{\partial q}{\partial L_r} = 2 \cos(2\beta_0(b - L_r)) \text{ and} \quad (9.17)$$

$$\frac{\partial s}{\partial \beta_\epsilon} = \frac{2b \cos(2b\beta_\epsilon)}{\beta_\epsilon} - \frac{2b\beta_\epsilon}{\beta_\epsilon^2}, \quad (9.18)$$

$$\frac{\partial s}{\partial b} = 2 \cos 2b\beta_\epsilon. \quad (9.19)$$

The the wavenumbers by the guide wavelength defined as

$$\beta = \sqrt{(2\pi f)^2 \mu \epsilon - \left(\frac{t'_{01}}{a}\right)^2}, \quad (9.20)$$

$$\beta_0 = \frac{2\pi}{\lambda_g} \sqrt{\left(\frac{2\pi f}{c_{\text{air}}}\right)^2 - \left(\frac{t'_{01}}{a}\right)^2}, \quad (9.21)$$

$$\beta_\epsilon = \frac{2\pi}{\lambda_{g\epsilon}} \sqrt{\left(\frac{2\pi f}{c_0}\right)^2 \epsilon'_R - \left(\frac{t'_{01}}{a}\right)^2}, \quad (9.22)$$

where $\epsilon'' \ll \epsilon'$ so that $\epsilon^* \approx \epsilon' - j\epsilon''$. The uncertainties in the wavenumbers are then found from the total differentials

$$\Delta\beta_0 = \frac{\partial\beta_0}{\partial c_{air}}\Delta c_{air} + \frac{\partial\beta_0}{\partial a}\Delta a + \frac{\partial\beta_0}{\partial f}\Delta f \text{ and} \quad (9.23)$$

$$\Delta\beta_\epsilon = \frac{\partial\beta_\epsilon}{\partial\epsilon'_R}\Delta\epsilon'_R + \frac{\partial\beta_0}{\partial a}\Delta a + \frac{\partial\beta_0}{\partial f}\Delta f, \quad (9.24)$$

where the partial derivatives are

$$\frac{\partial\beta_\epsilon}{\partial f} = \frac{4\pi^2 f \epsilon'_R}{c_0^2} \frac{1}{\sqrt{\left(\frac{2\pi f}{c_0}\right)^2 \epsilon'_R - \left(\frac{t'_{01}}{a}\right)^2}}, \quad (9.25)$$

$$\frac{\partial\beta_\epsilon}{\partial a} = \frac{(t'_{01})^2}{a^3} \frac{1}{\sqrt{\left(\frac{2\pi f}{c_0}\right)^2 \epsilon'_R - \left(\frac{t'_{01}}{a}\right)^2}}, \quad (9.26)$$

$$\frac{\partial\beta_\epsilon}{\partial\epsilon'_R} = \frac{(2\pi f)^2}{2c_0^2} \frac{1}{\sqrt{\left(\frac{2\pi f}{c_0}\right)^2 \epsilon'_R - \left(\frac{t'_{01}}{a}\right)^2}}, \quad (9.27)$$

and

$$\frac{\partial\beta_0}{\partial f} = \frac{4\pi^2 f}{c_{air}^2} \frac{1}{\sqrt{\left(\frac{2\pi f}{c_{air}}\right)^2 - \left(\frac{t'_{01}}{a}\right)^2}}, \quad (9.28)$$

$$\frac{\partial\beta_0}{\partial a} = \frac{(t'_{01})^2}{a^3} \frac{1}{\sqrt{\left(\frac{2\pi f}{c_{air}}\right)^2 - \left(\frac{t'_{01}}{a}\right)^2}}, \quad (9.29)$$

$$\frac{\partial\beta_0}{\partial c_{air}} = -\frac{(2\pi f)^2}{c_{air}^3} \frac{1}{\sqrt{\left(\frac{2\pi f}{c_{air}}\right)^2 - \left(\frac{t'_{01}}{a}\right)^2}} \quad (9.30)$$

Loss tangent uncertainty due to the uncertainties in the measured parameters is shown in Fig. 9.2. We can observe a highly frequency-dependent behavior in loss tangent uncertainty due to cavity dimensions uncertainties. Typically we measure samples with much lower loss than our hypothetical 10-mm thick sample with permittivity $10 - j0.001$, and loss tangent uncertainty is usually dominated by uncertainties in the measurement of resonance Q.

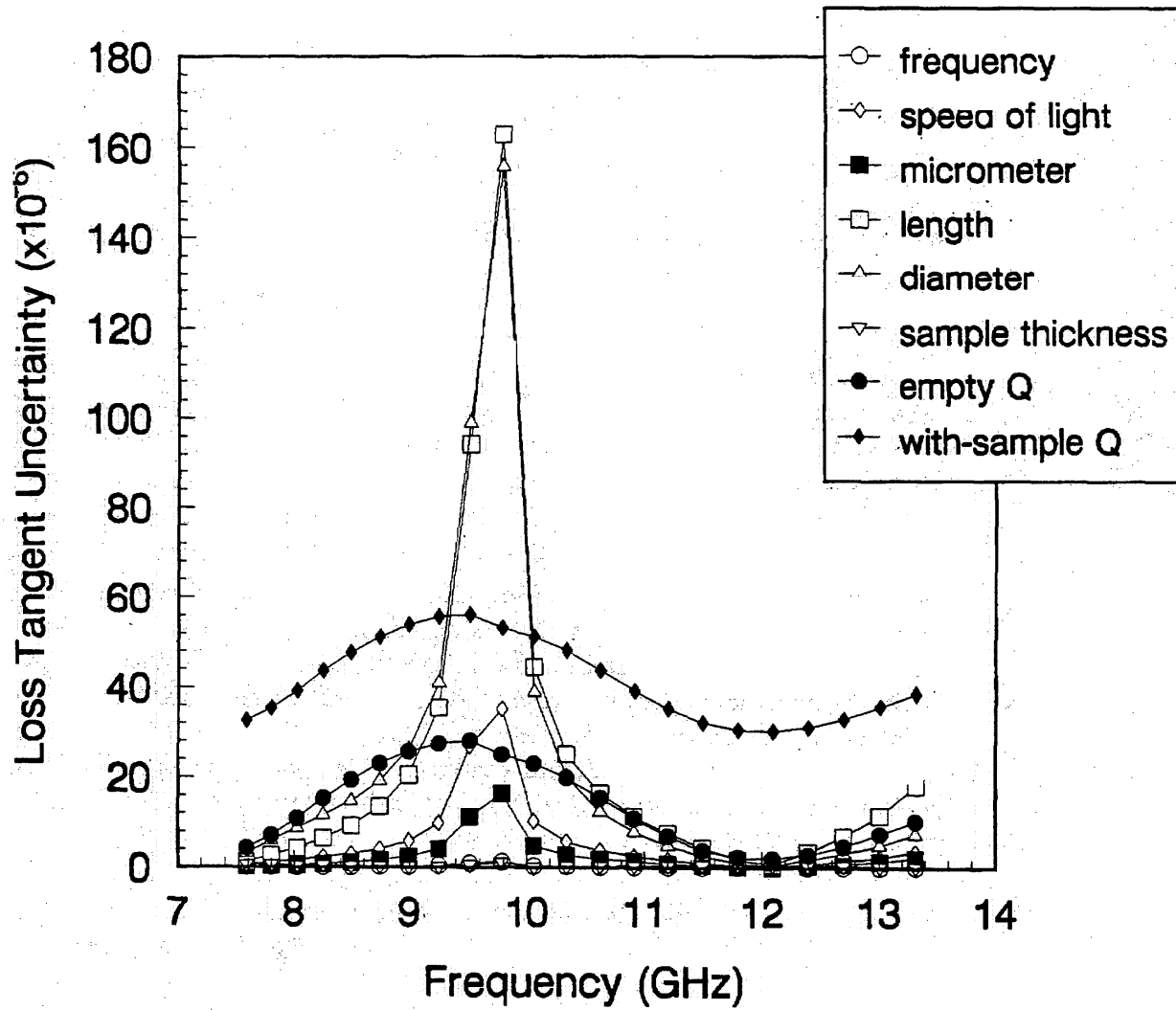


Figure 9.2: Loss tangent uncertainty $\Delta(\tan\delta)$ versus frequency due to uncertainties in various measured parameters for a 10-mm thick sample with $\epsilon'_R = 10 - j0.001$.

Chapter 10

Future Work

The change in length of the cavity is monitored by the sensing micrometer. As described in Sec. 4.1.2, the sensing micrometer's tip gradually wears a small pit into the extended finger of the tuner-endplate yoke. The yoke is made of aluminum and is therefore softer than the sensing micrometer's tip. To improve the repeatability of cavity length measurement, we will press fit a hardened steel plug into the yoke for the sensing micrometer to press against.

The uncertainty estimates given in Chapter 9 assume that systematic errors have been accounted for. To determine whether our accuracy estimates are truly correct, we must compare our results with the results of other instruments of similar accuracy. The glass results given in Sec. 8.1 compare favorably with results from an international intercomparison completed in 1974. NIST plans to engage in future international comparisons not only to qualify the NIST resonator for calibration service, but also to qualify our stock of cross-linked polystyrene, fused silica, and alumina as a Standard Reference Material.

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Appendix A

Partial Filling Factor for Loss Determination

Previously we showed that the unloaded quality factor of the cavity with sample inserted is

$$\frac{1}{Q_{0, \text{ sample}}} = \frac{1}{Q_C} + \frac{1}{Q_S}, \quad (\text{A.1})$$

where $\tan \delta = 1/Q_S$. The empty unloaded cavity quality factor is

$$\frac{1}{Q_{0, \text{ empty}}} = \frac{1}{Q_C}. \quad (\text{A.2})$$

In practice, a loaded cavity quality factor is always measured, which includes coupling port losses Q_E

$$\frac{1}{Q_L} = \frac{1}{Q_0} + \frac{1}{Q_E}. \quad (\text{A.3})$$

where Q_0 is either with or without sample insertion. The loaded quality factor is related to the unloaded quality factor by (2.80).

The analysis in Secs. 2.2.5 and 2.2.6 derived the theoretical Q_C and the measurement approach for evaluating total power loss in the cavity walls, as well as a technique for separating sidewall from end plate losses. However, the loss in the sidewall of the cavity is different when the sample is inserted as opposed to when the cavity is empty (air filled). Therefore, it is necessary to determine the energy stored and power dissipated both (1) in the sample of arbitrary thickness having relative complex permittivity $\epsilon'_R - j\epsilon''_R$ and (2) in the air-filled portion of the cavity when the sample is inserted. This allows us to account for the fact that the cavity is only partially filled with dielectric.

Reference [2] shows that the theoretical quality factor Q_0 of a cylindrical cavity of radius a that is partially filled with a sample of thickness b having a relative complex permittivity $\epsilon_R^* = \epsilon'_R - j\epsilon''_R$, is given by

$$Q_0 = \frac{\{F(2b - G) + [2(L - b) - U]/\epsilon'_R\} a\omega^2 \mu'_R \epsilon'_R}{\{\delta_s c^2 (k_c^2 [(2b - G)F + 2(L - b) - U] + 2a [F\beta_1^2 + \beta_0^2]) + a\omega^2 \mu'_R \epsilon'_R F(2b - G) \tan \delta\}} \quad (\text{A.4})$$

where c is the speed of light in air, ω is angular frequency for any TE_{01p} resonant mode, δ_s is the effective penetration depth given by (2.123) of the cavity's wall and end plates, and

$$F = \frac{\sin^2 [\beta_0(L - b)]}{\sin^2 [\beta_1 b]} \quad (\text{A.5})$$

$$G = \frac{\sin [2\beta_1 b]}{\beta_1} \quad (\text{A.6})$$

$$U = \frac{\sin [2\beta_0(L - b)]}{\beta_0} \quad (\text{A.7})$$

Equation (A.4) is valid for a dielectrically lossy sample and takes into account the effect of finite wall and end plate loss intrinsic to the cavity. The measured unloaded quality factor $Q_{0,m}$ with the dielectric sample is generally somewhat lower than the theoretical Q_0 in (A.4). From (2.123) and (A.4) Cook [2] has further shown that the loss tangent may be precisely calculated from

$$\tan \delta = \left\{ \frac{F(2b - G) + [2(L - b) - U]/\epsilon'_R}{F(2b - G)} \right\} \left[\frac{1}{Q_{0,m}} - \frac{1}{\bar{Q}} \right] \quad (\text{A.8})$$

where

$$\bar{Q} = \frac{\{F(2b - G) + [2(L - b) - U]/\epsilon'_R\} a\omega^2 \mu'_R \epsilon'_R}{\delta_{s,m} c^2 [k_c^2 [F(2b - G) + 2(L - b) - U] + 2a (F\beta_1^2 + \beta_0^2)]} \quad (\text{A.9})$$

Equation (A.8) is similar to previous analysis except it now contains a filling factor because the specimen only partly fills the cavity.

Appendix B

Radial Air Gap Correction

One possible source of error in dielectric property measurements at microwave frequencies is that due to the air gap between the sample under test and the wall of the waveguide or resonator fixture. The solid dielectric under test is usually machined into the shape of a parallel-sided disk for a right-circular cylindrical resonator, and clearance must be left between the sides of the disk and the wall of the resonator so that the disk may move freely in the resonator. Exactly how much clearance is permissible before the error in dielectric characterization of the sample becomes significant is a problem that must be examined.

Consider two situations for our cavity resonator in which the walls are modelled as perfect electrical conductors. The first case (Fig. B.1) illustrates a dielectric sample in a cylindrical resonator with no air gap. The second case (Fig. B.2) shows the actual situation in which an air gap exists. The presence of this gap results in an increased length of $L_1 + \Delta L$ at resonance, where $L_1 = L - b$ for the ideal case. This, in turn, leads to a measured permittivity of the sample under test that is too small.

One approach to correct for the presence of an air gap is to consider the change in resonant frequency due to the gap with perturbation theory, where the ideal situation shown in Fig. B.1 is perturbed. Since we know the measured length $L_1 + \Delta L$ at measured resonant frequency f_m and the radii of both the sample under test and dielectric resonator, we can estimate the resonant frequency for the resonator with no air gap but of length $L_1 + \Delta L$ for use in the transcendental equation for permittivity calculation. Here a method described by Grivet [44] and Bussey [45] is outlined.

We noted that at resonance the total energy interior to the cavity is constant, as a function of time and that the average values of the energy stored in the

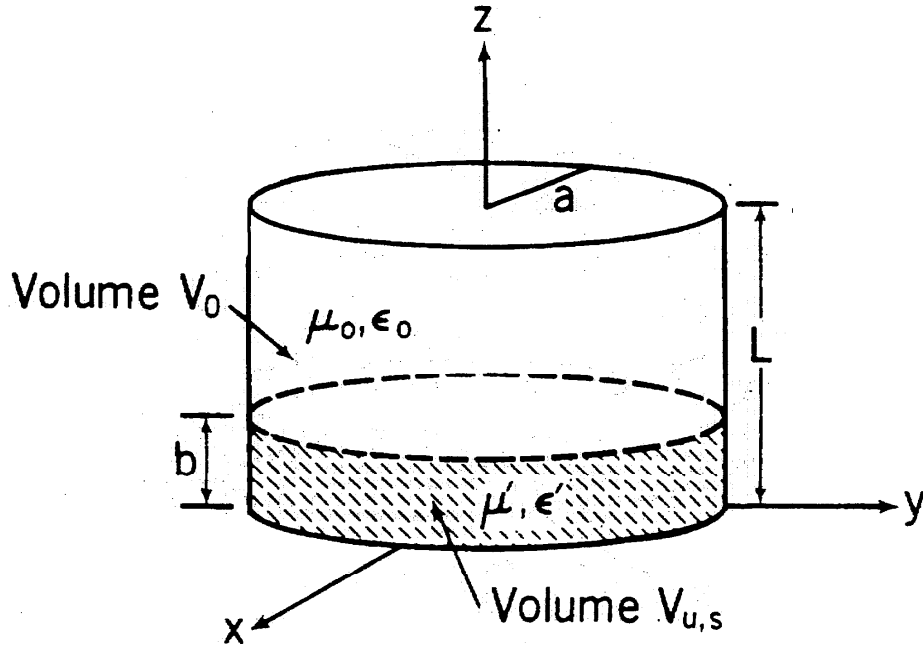


Figure B.1: Unperturbed cylindrical cavity resonator with no air gap present between sample under test and resonator side wall.

electric or magnetic field are equal. Hence, for TE_{01p} mode structure,

$$\iiint_{\text{cavity}} \epsilon' |E_o E_o^*| dV = \iiint_{\text{cavity}} \mu |\vec{H} \vec{H}^*| dV, \quad (\text{B.1})$$

where the * denotes complex conjugate and the integrals are taken over the entire volume of the cavity. For a magnetically impermeable sample under test, μ may be taken outside of the above integral and since, by Faraday's law,

$$\nabla \times \vec{E} = -j2\pi f \mu \vec{H}. \quad (\text{B.2})$$

(B.1) can be written

$$\iiint_{\text{cavity}} \epsilon' |\vec{E}|^2 dV = -\frac{1}{4\pi^2 f^2 \mu} \iiint_{\text{cavity}} [\nabla \times \vec{E}]^2 dV. \quad (\text{B.3})$$

If we now make the assumption that the electric field in the perturbed cavity is approximately that of the unperturbed, we may write

$$\vec{E} \approx \vec{E}_u. \quad (\text{B.4})$$

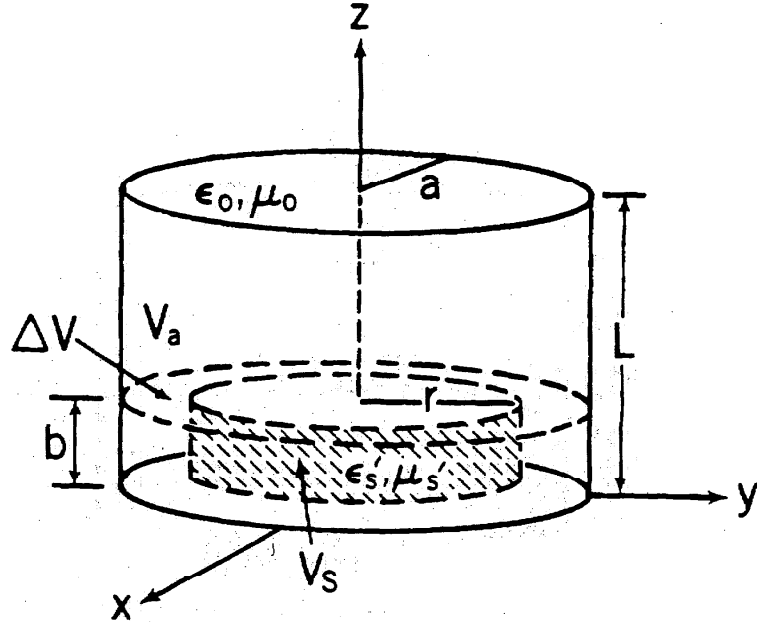


Figure B.2: Perturbed cylindrical cavity with uniform air gap between sample and resonator wall.

since we know that \vec{E}_u satisfies the correct boundary conditions at the conducting walls of the cavity and is continuous at the boundary of the sample under test. This approximation is valid only if the air gap is not too large relative to λ_g and if the air gap is uniform, so azimuthal symmetry is maintained. From (B.4),

$$\nabla \times \vec{E}_u = -j2\pi f_u \mu \vec{H}_u, \quad (\text{B.5})$$

where f_u denotes the unperturbed resonant frequency. Hence,

$$\begin{aligned} f_m^2 \left\{ \iiint_{V_0} \epsilon_0 |\vec{E}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s} - \Delta V} \epsilon_s |\vec{E}_{u,s}|^2 dV \right. \\ \left. + \iiint_{\Delta V} \epsilon_0 |\vec{E}_{u,s}|^2 dV \right\} = \\ f_u^2 \left\{ \iiint_{V_0} \mu_0 |\vec{H}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s} - \Delta V} \mu_0 |\vec{H}_{u,s}|^2 dV \right. \\ \left. + \iiint_{\Delta V} \mu_0 |\vec{H}_{u,s}|^2 dV \right\}, \quad (\text{B.6}) \end{aligned}$$

where f_m represents the measured resonant frequency, the subscripts u and s denote unperturbed and sample, and where the integration over ΔV is an integration over the volume of the air gap which is

$$\Delta V = \pi a^2 b - \pi r^2 b = \pi b (a^2 - r^2), \quad (\text{B.7})$$

where a and r are the respective radii of the resonator and the sample under test. Equation (B.6) may be simplified as follows:

$$\begin{aligned} & f_m^2 \left\{ \iiint_{V_0} \epsilon_0 |\vec{E}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \epsilon'_s |\vec{E}_{u,s}|^2 dV \right. \\ & \left. + \iiint_{\Delta V} (\epsilon_0 - \epsilon'_s) |\vec{E}_{u,s}|^2 dV \right\} \\ & = f_u^2 \left\{ \iiint_{V_0} \mu |\vec{H}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \mu |\vec{H}_{u,s}|^2 dV \right\}. \end{aligned} \quad (\text{B.8})$$

We now note that

$$\begin{aligned} & f_u^2 \left\{ \iiint_{V_0} \mu |\vec{H}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \mu |\vec{H}_{u,s}|^2 dV \right\} \\ & \equiv f_m^2 \left\{ \iiint_{V_0} \epsilon_0 |\vec{E}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \epsilon'_s |\vec{E}_{u,s}|^2 dV \right\}, \end{aligned} \quad (\text{B.9})$$

from (B.1), so

$$\begin{aligned} & f_m^2 \left\{ \iiint_{V_0} \epsilon_0 |\vec{E}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \epsilon'_s |\vec{E}_{u,s}|^2 dV \right. \\ & \left. - \iiint_{\Delta V} (\epsilon'_s - \epsilon_0) |\vec{E}_{u,s}|^2 dV \right\} \\ & = f_u^2 \left\{ \iiint_{V_0} \epsilon_0 |\vec{E}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \epsilon'_s |\vec{E}_{u,s}|^2 dV \right\}. \end{aligned} \quad (\text{B.10})$$

We now have a correction formula for the unperturbed resonant frequency f_u in terms of the measured resonant frequency f_m in the presence of an air gap of volume ΔV ; that is,

$$f_u = f_m \left\{ 1 - \frac{\iiint_{\Delta V} (\epsilon'_s - \epsilon_0) |\vec{E}_{u,s}|^2 dV}{\iiint_{V_0} \epsilon_0 |\vec{E}_{u, \text{air}}|^2 dV + \iiint_{V_{u,s}} \epsilon'_s |\vec{E}_{u,s}|^2 dV} \right\}^{1/2} \quad (\text{B.11})$$

Of course, when there is no air gap $\Delta V = 0$ and $f_u = f_m$. For the mode-filtered TE_{01p} cavity of concern here, we have only an azimuthal component of the electric field that is of concern in the above integrations.

If an air gap is present between the sample under test and wall of the waveguide of the cavity resonator, the cavity will be at resonance at an increased length (compared to the situation where no gap exists). This leads to a measured dielectric constant that is too small. The procedure from the analysis above, which is valid if the air gap is not too large relative to the guide wavelength and if the air gap is uniform so that azimuthal symmetry is maintained, is as follows:

1. Measure the resonant frequency f_m and measured length of the cavity at resonance.
2. Compute the unperturbed resonant frequency f_u for a resonator with no air gap but with the length measured in step 1 above.
3. Use the unperturbed frequency derived in step 2 for ϵ'_R determination in (2.46).

An alternate form of (B.11) for computing the correction to the measured relative dielectric constant in the presence of a radial air gap for TE_{01p} mode is

$$\Delta\epsilon'_R = (\epsilon'_R - 1) \frac{\int_r^a \int_0^{2\pi} [J_1(k_c r_1)]^2 \sin^2(\beta_1 z) r_1 dr_1 d\phi}{\int_0^a \int_0^{2\pi} [J_1(k_c r_1)]^2 \sin^2(\beta_1 z) r_1 dr_1 d\phi}, \quad (\text{B.12})$$

which reduces to

$$\Delta\epsilon'_R = \frac{(\epsilon'_R - 1)}{2} \left\{ 1 - \left(\frac{r}{a}\right)^2 \frac{\Psi}{J_0^2(k_c a)} \right\}, \quad (\text{B.13})$$

with

$$\Psi = J_0^2(k_c r) - \frac{2}{k_c r} J_0(k_c r) J_1(k_c r) + J_1^2(k_c r). \quad (\text{B.14})$$

Appendix C

60-mm Cavity Drawings

To be consistent with shop tooling, dimensions for resonator parts were specified in inches. Apologies are offered for our deviation from Standard International units. In the following drawings, dimensions are in inches unless otherwise specified.

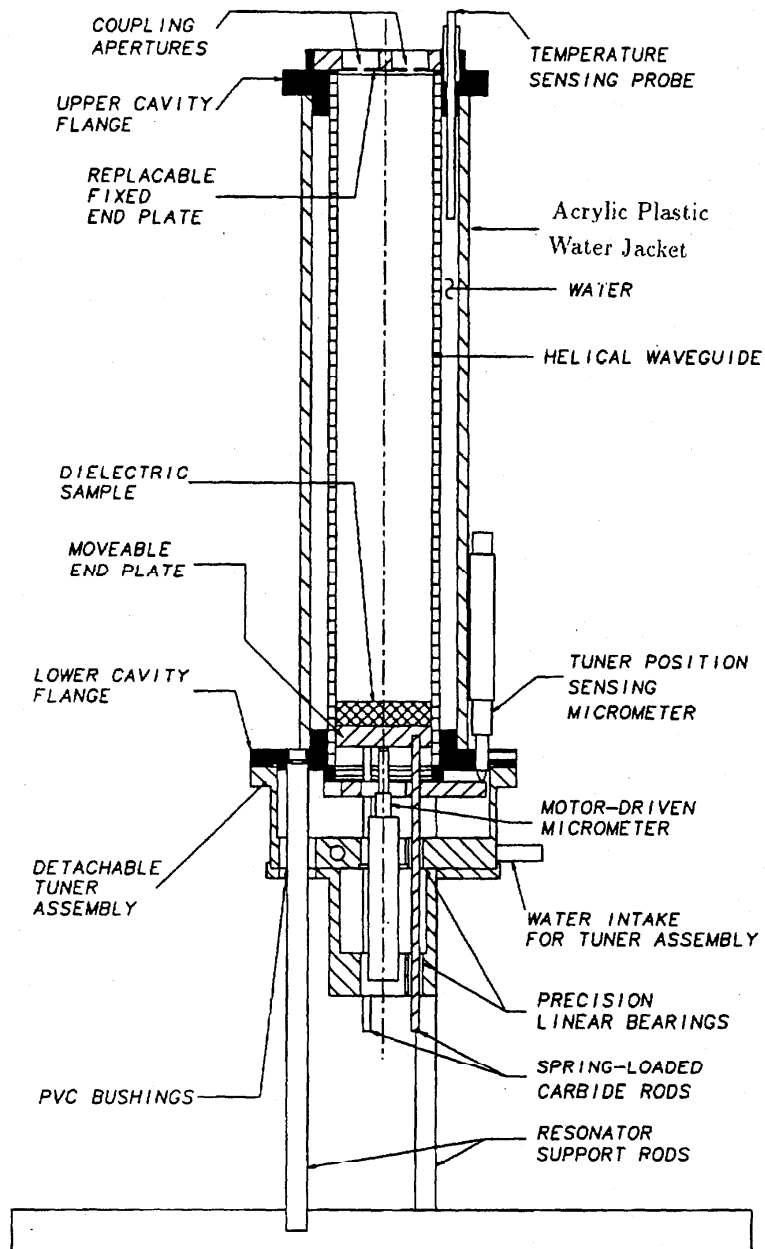


Figure C.1: Cutaway assembly drawing of NIST cavity resonator.

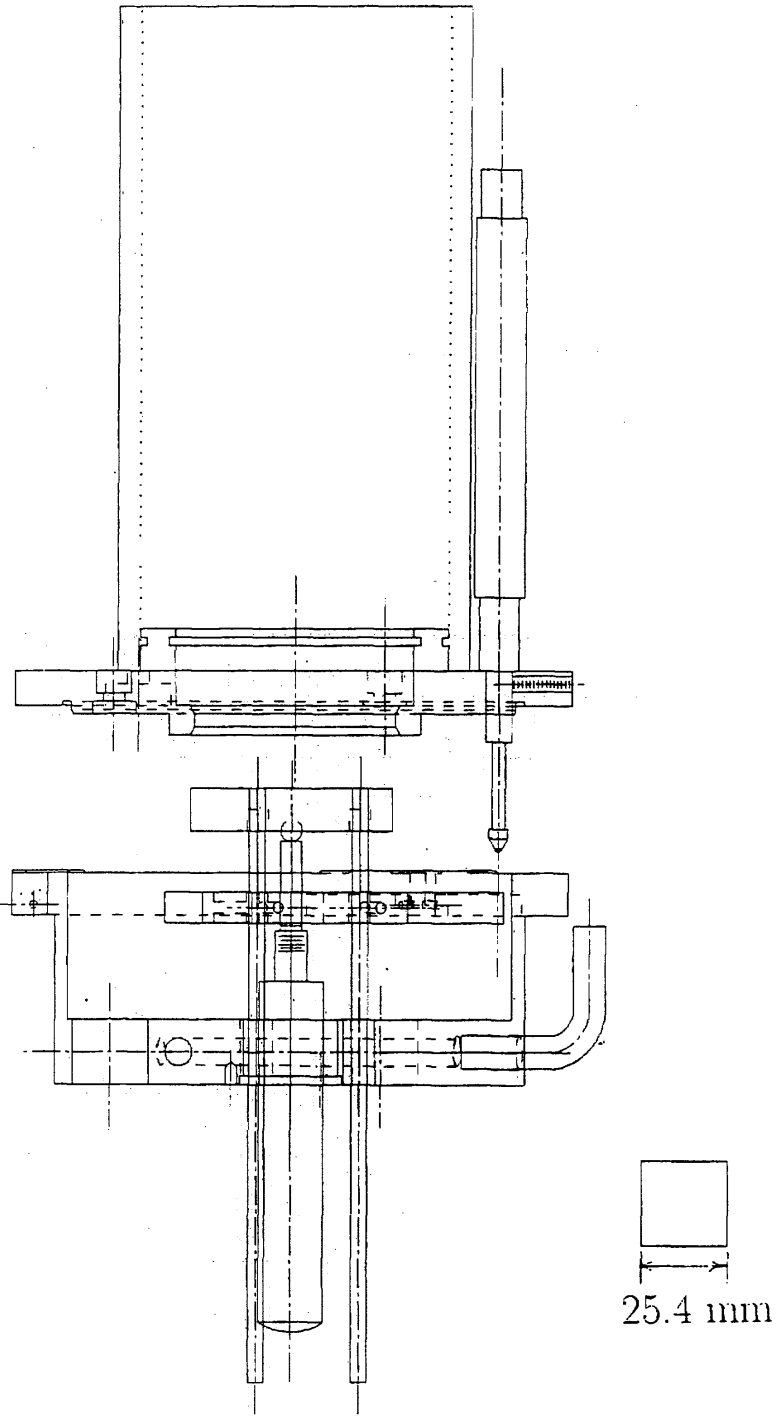


Figure C.2: Detail of tuner assembly and flange.

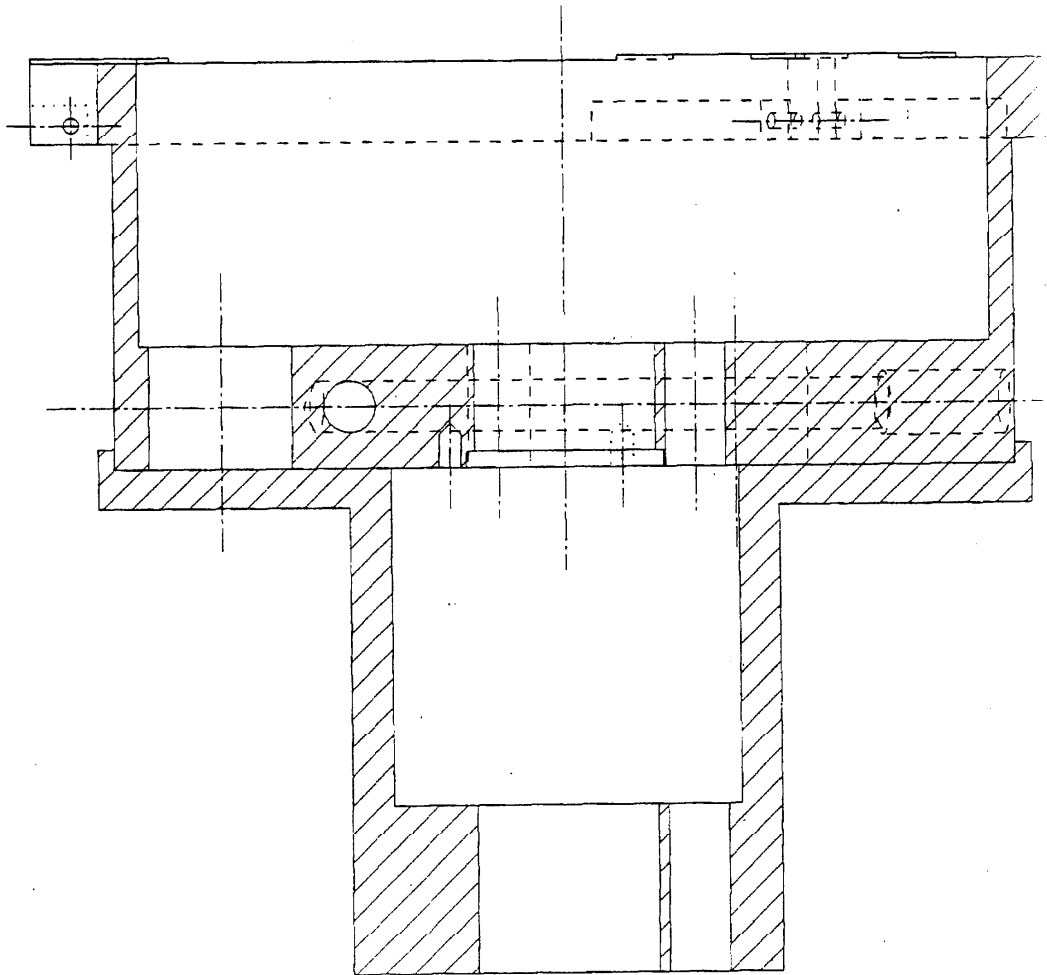


Figure C.3: Outaway side view of tuner-base and extension.

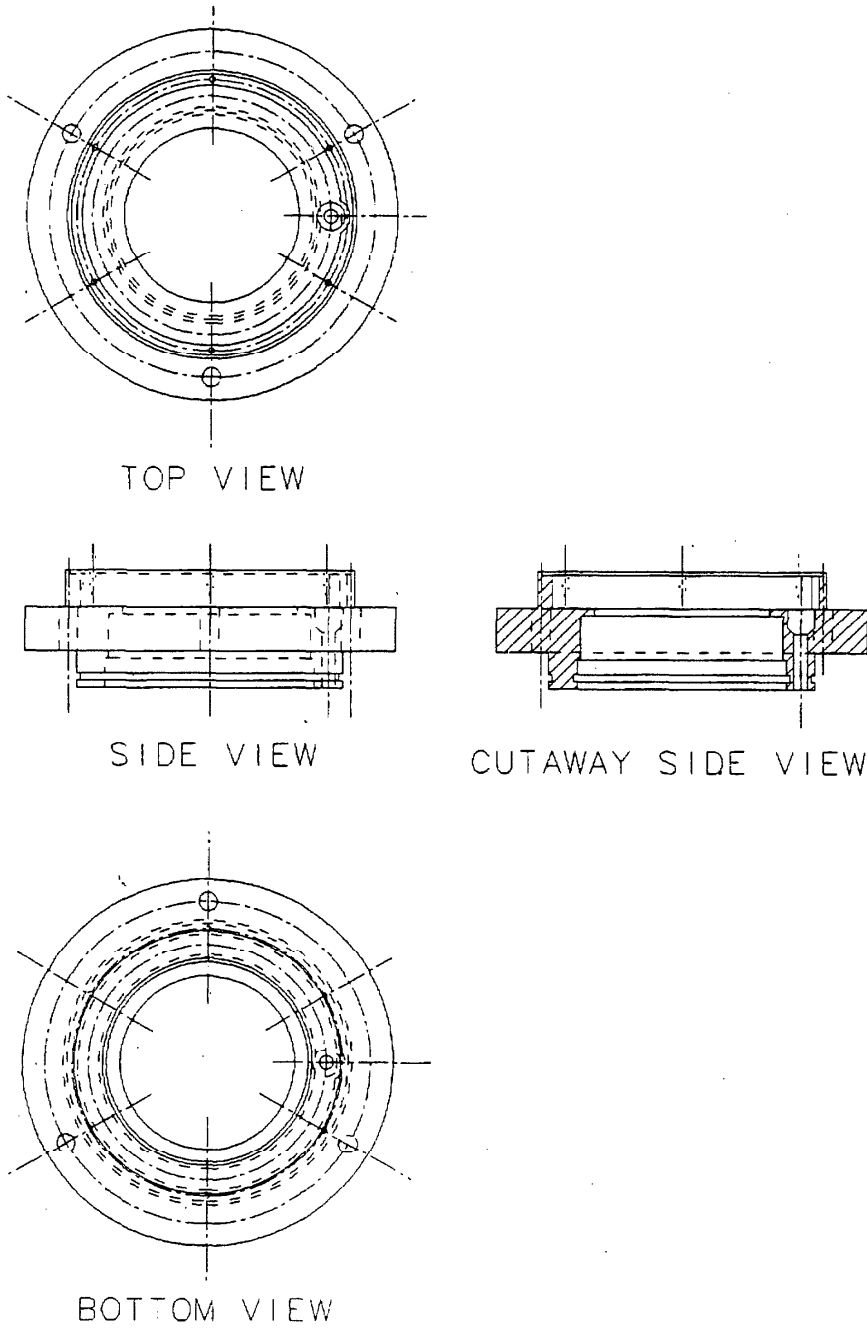


Figure C.4: Orthogonal views of coupling flange.

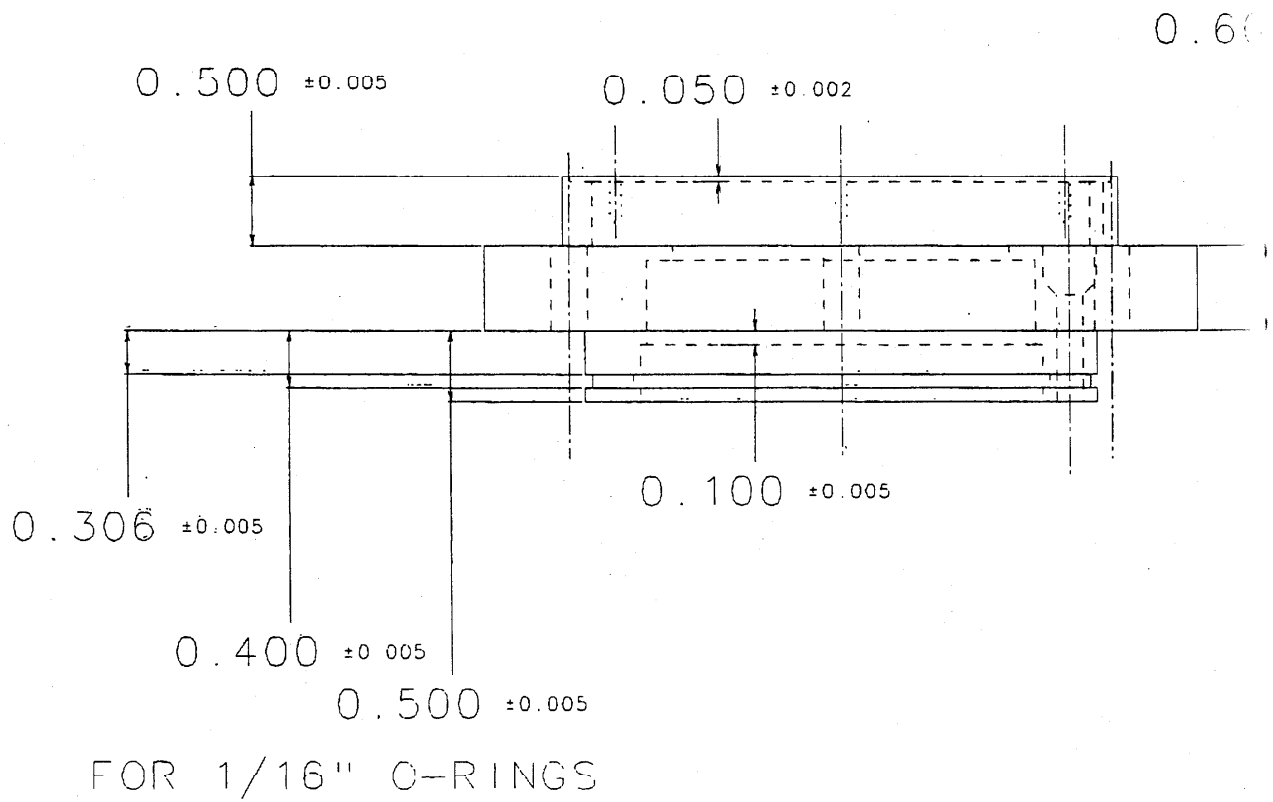


Figure C.5: Side view of coupling flange.

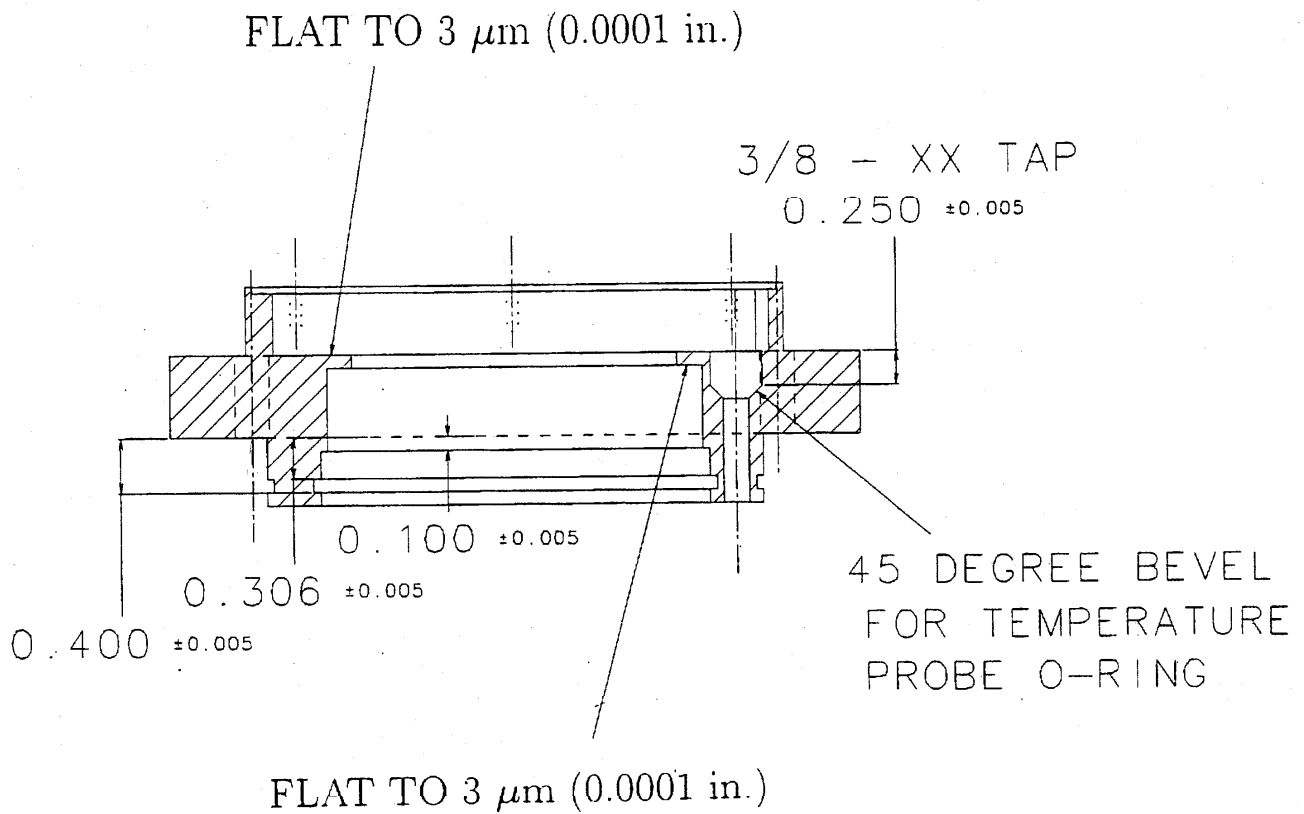


Figure C.6: Cutaway side view of coupling flange.

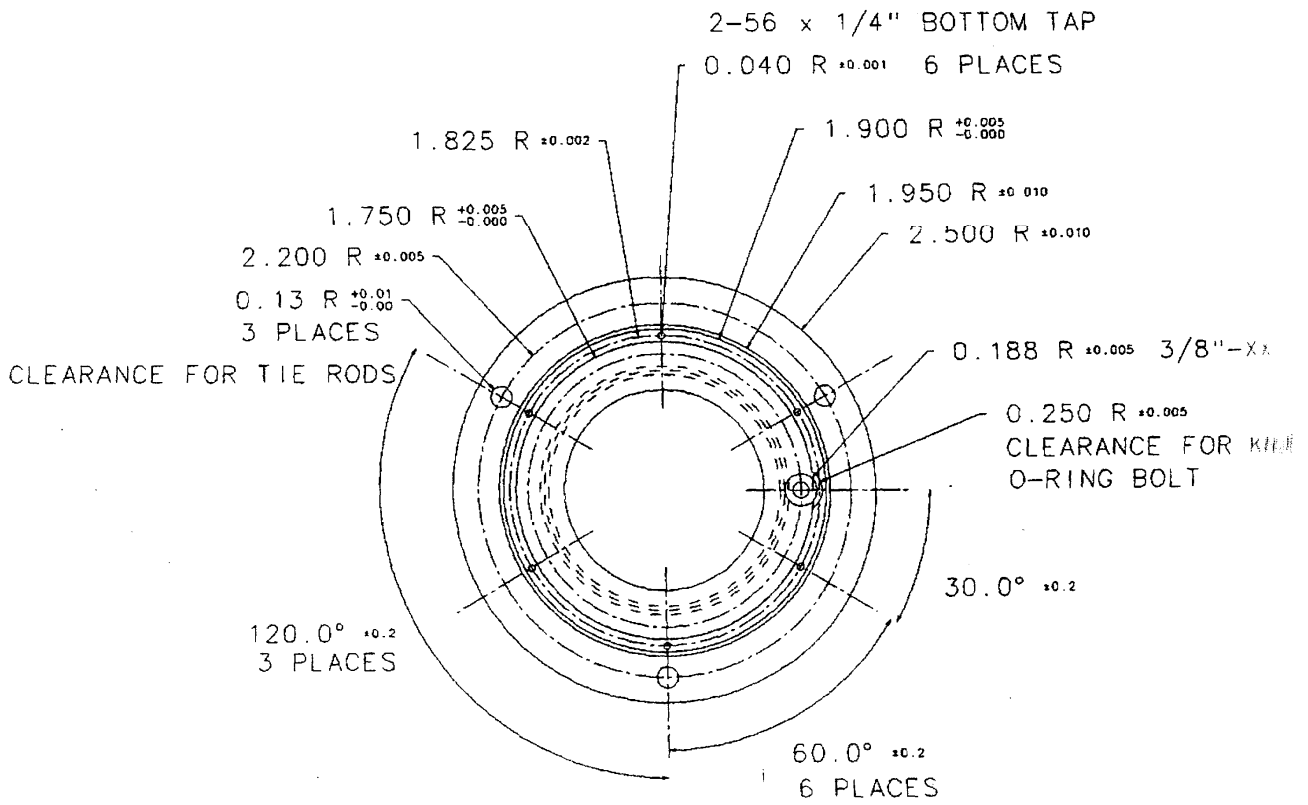


Figure C.7: Top view of coupling flange.

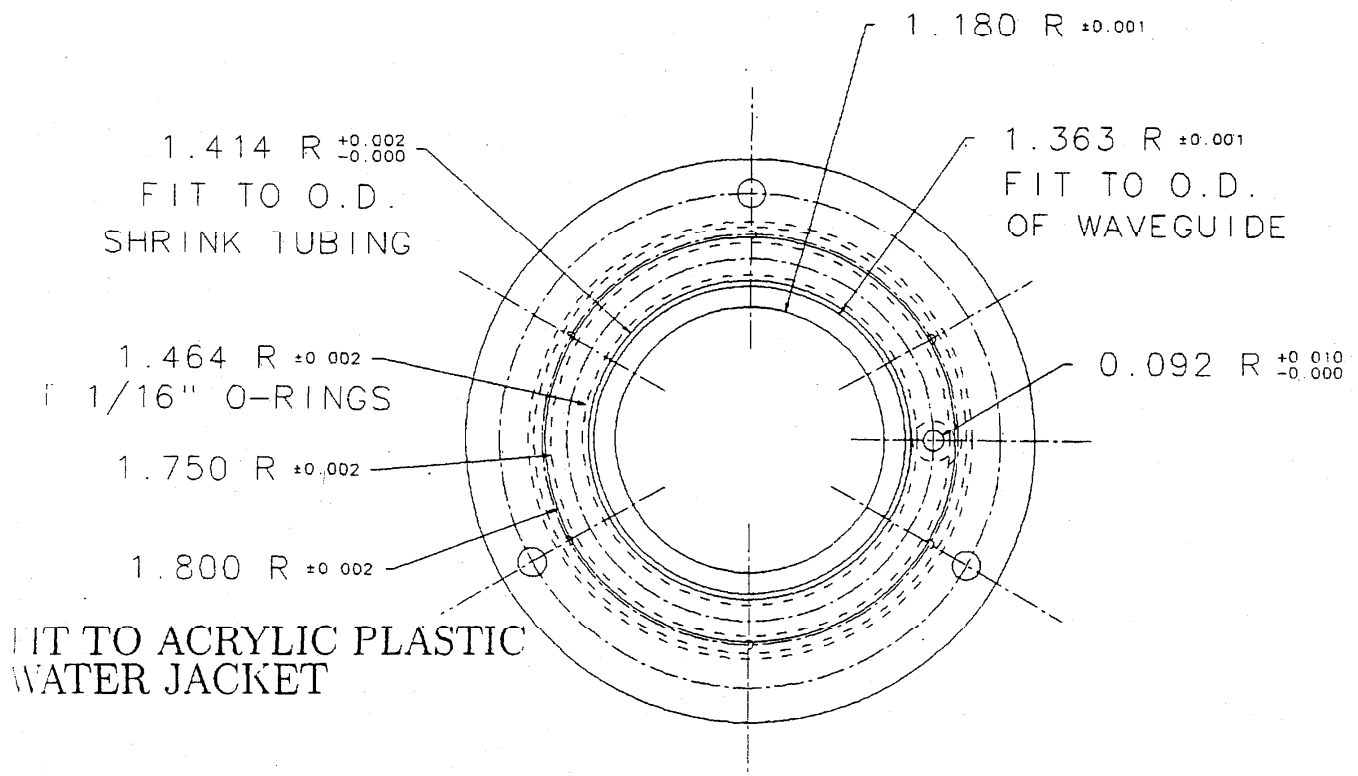
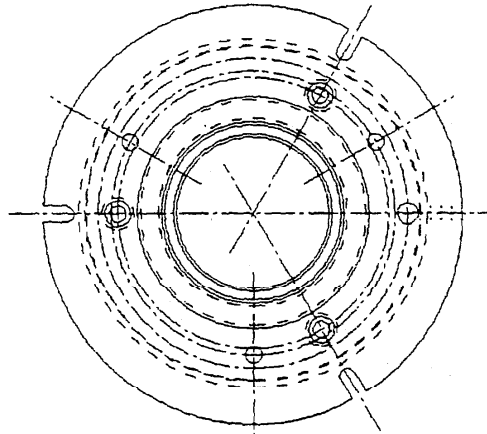
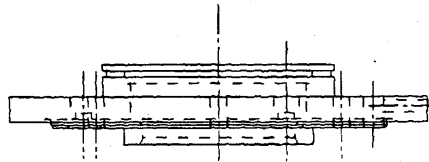


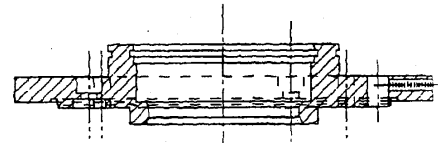
Figure C.8: Bottom view of coupling flange.



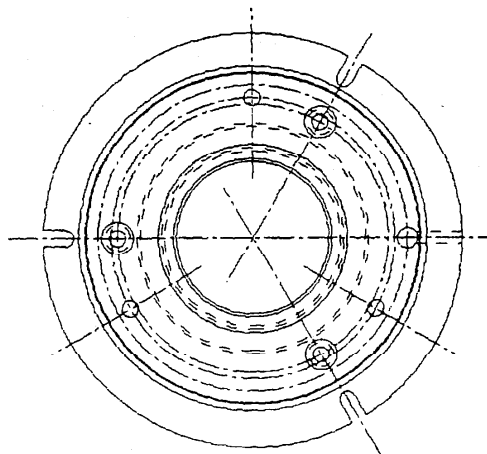
TOP VIEW



SIDE VIEW



CUTAWAY SIDE VIEW



BOTTOM VIEW

Figure C.9: Orthogonal views of tuner flange.

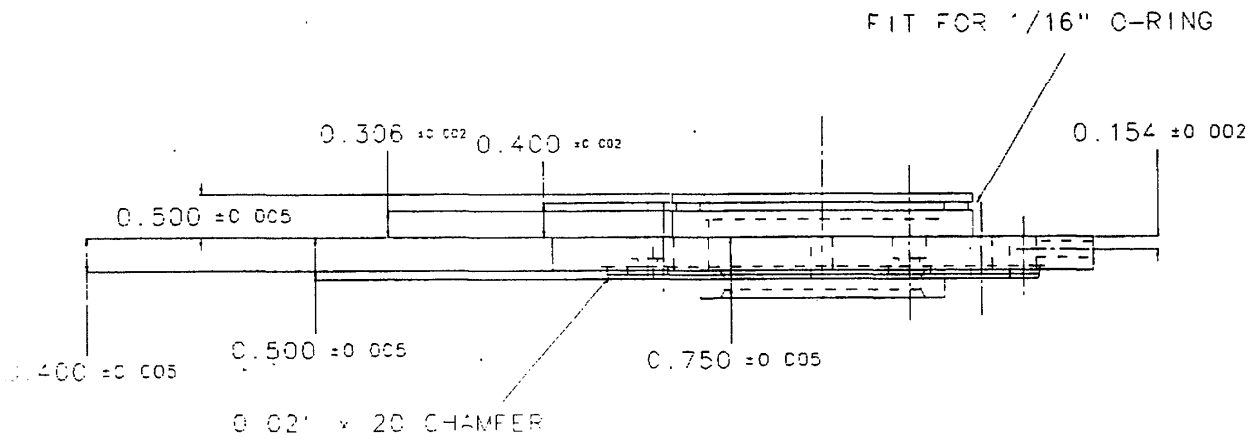


Figure C.10: Side view of coupling flange.

FIT TO O.D. OF SHRINK TUBING
AND HELICAL WAVEGUIDE

FLAT TO $3 \mu\text{m}$ (0.0001 in.) 001 INCH

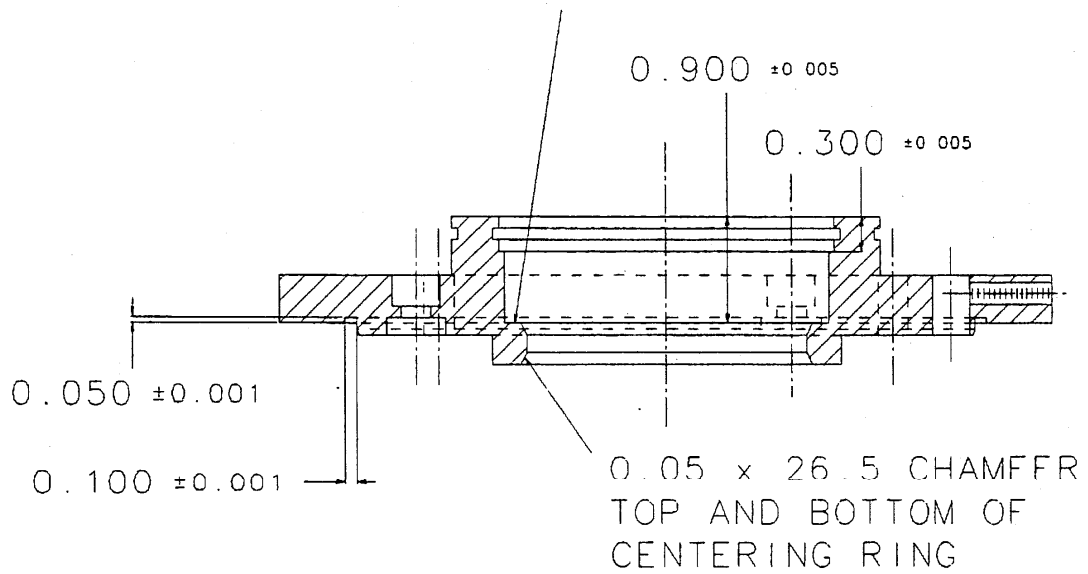


Figure C.11: Cutaway side view of coupling flange.

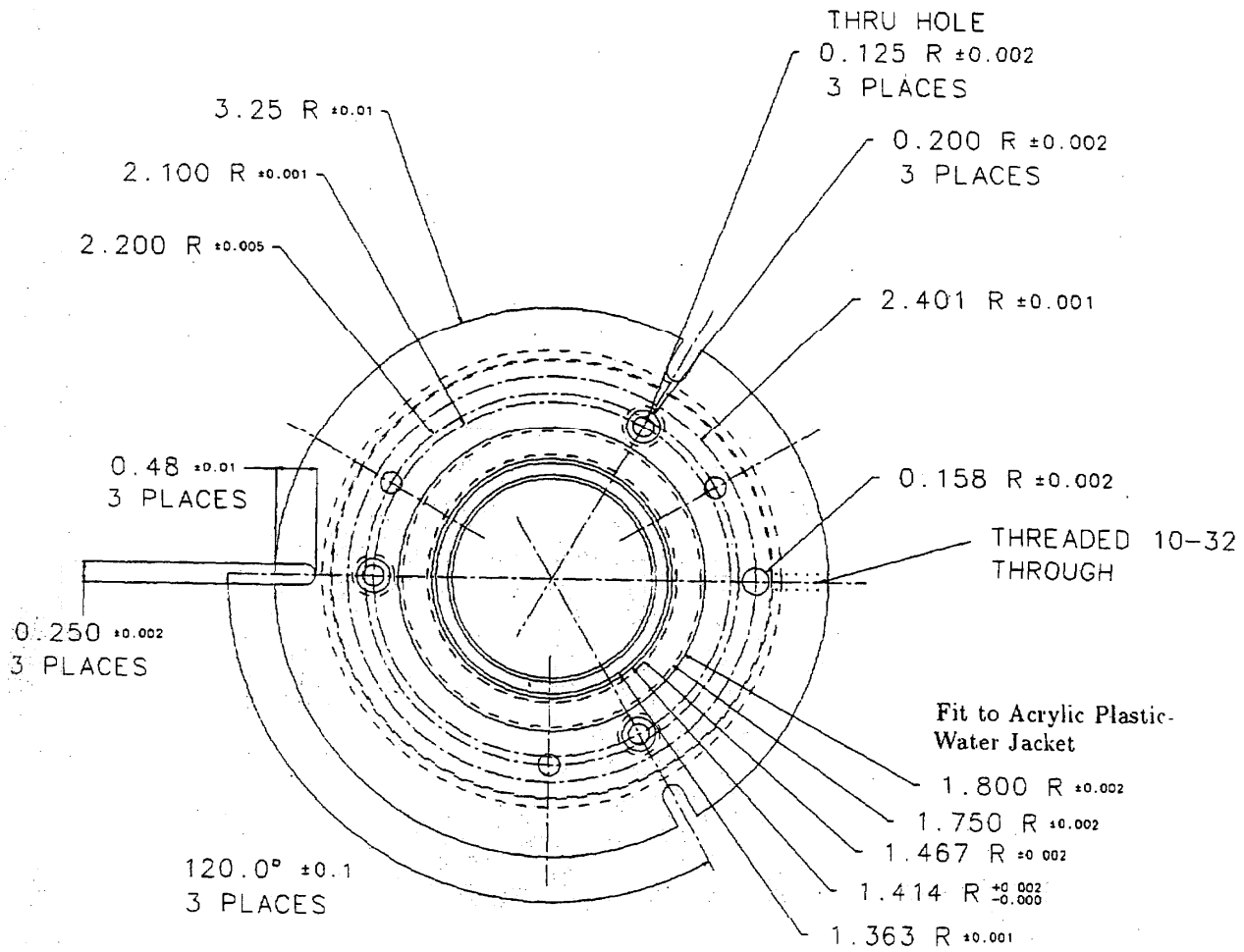


Figure C.12: Top view of tuner flange.

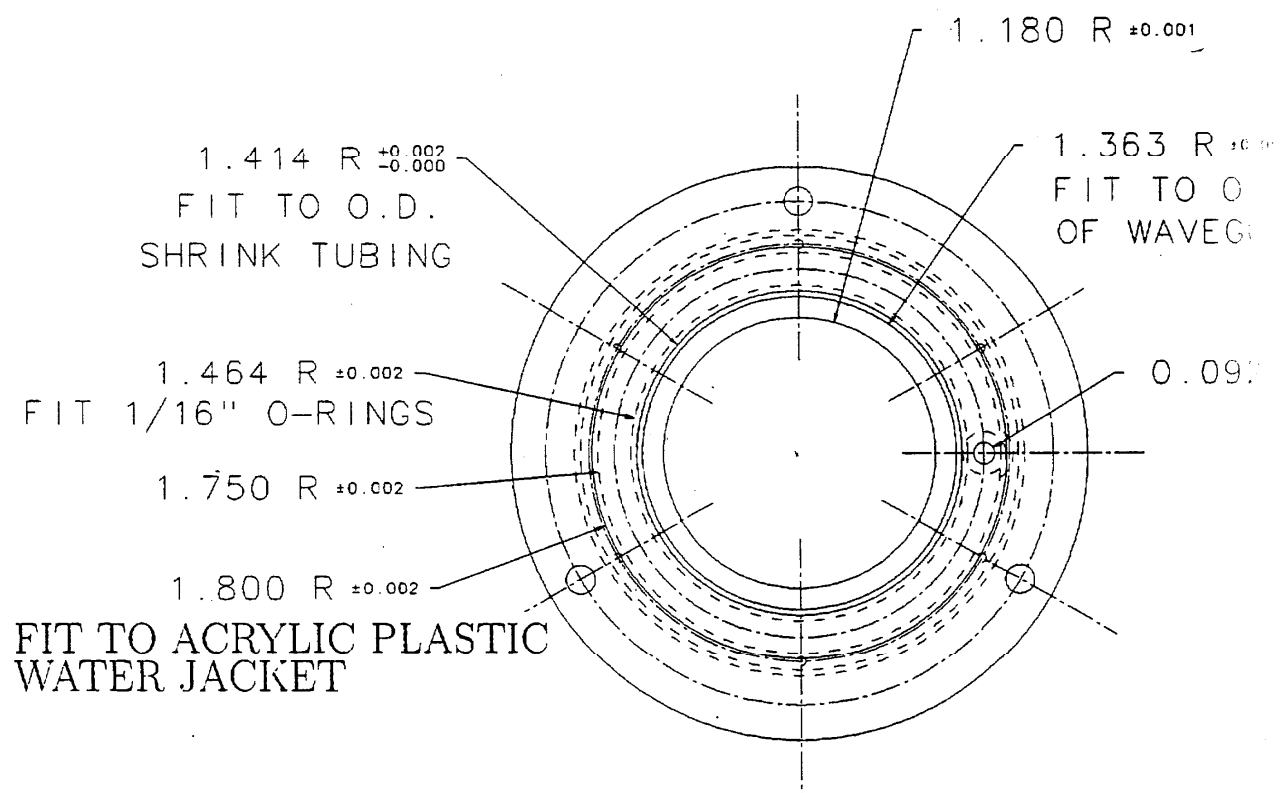
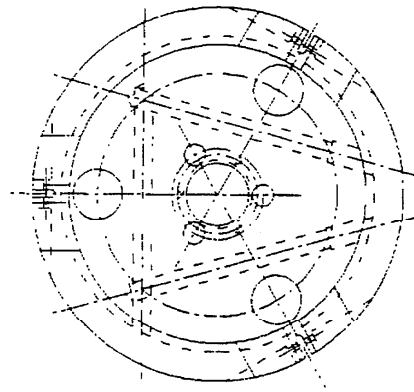
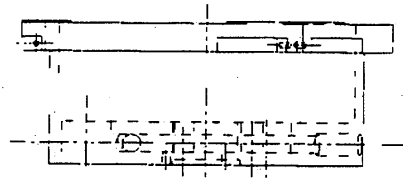


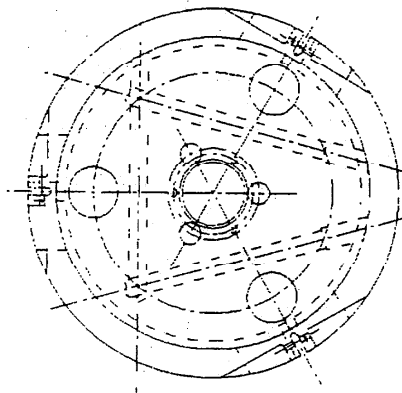
Figure C.13: Bottom view of tuner flange.



TOP VIEW



SIDE VIEW



BOTTOM VIEW

Figure C.14: Orthogonal views of tuner base.

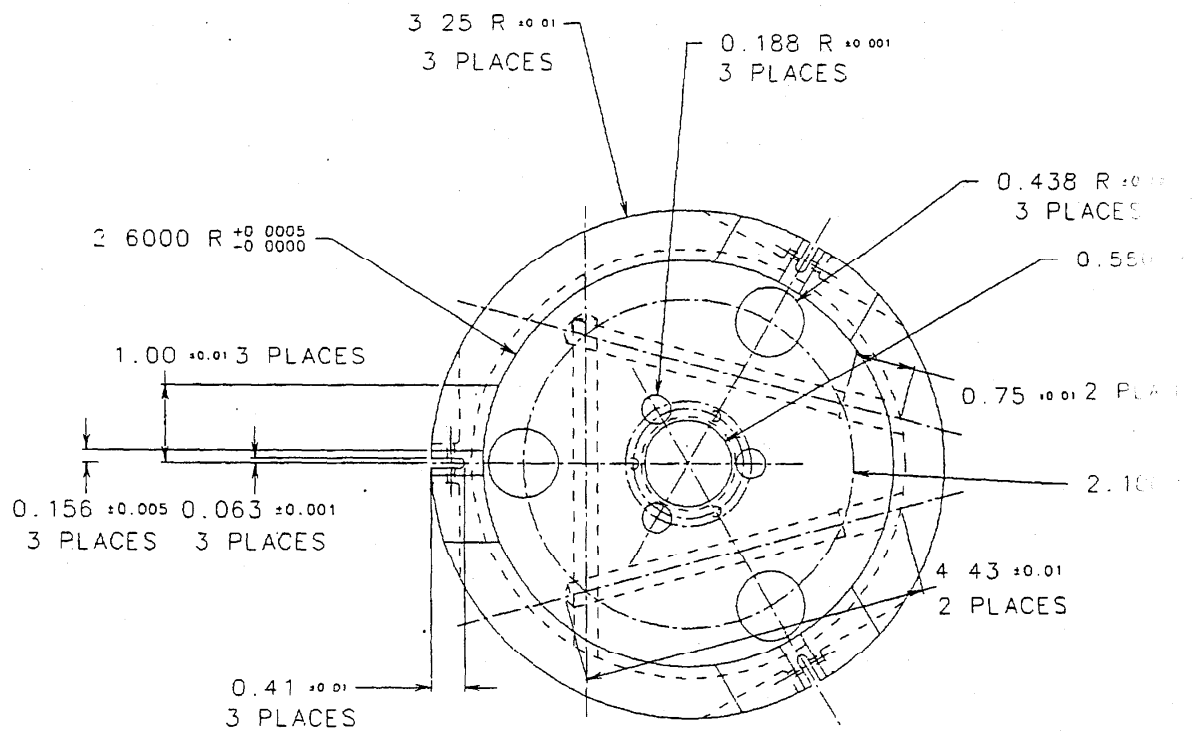


Figure C.15: Top view of tuner base.

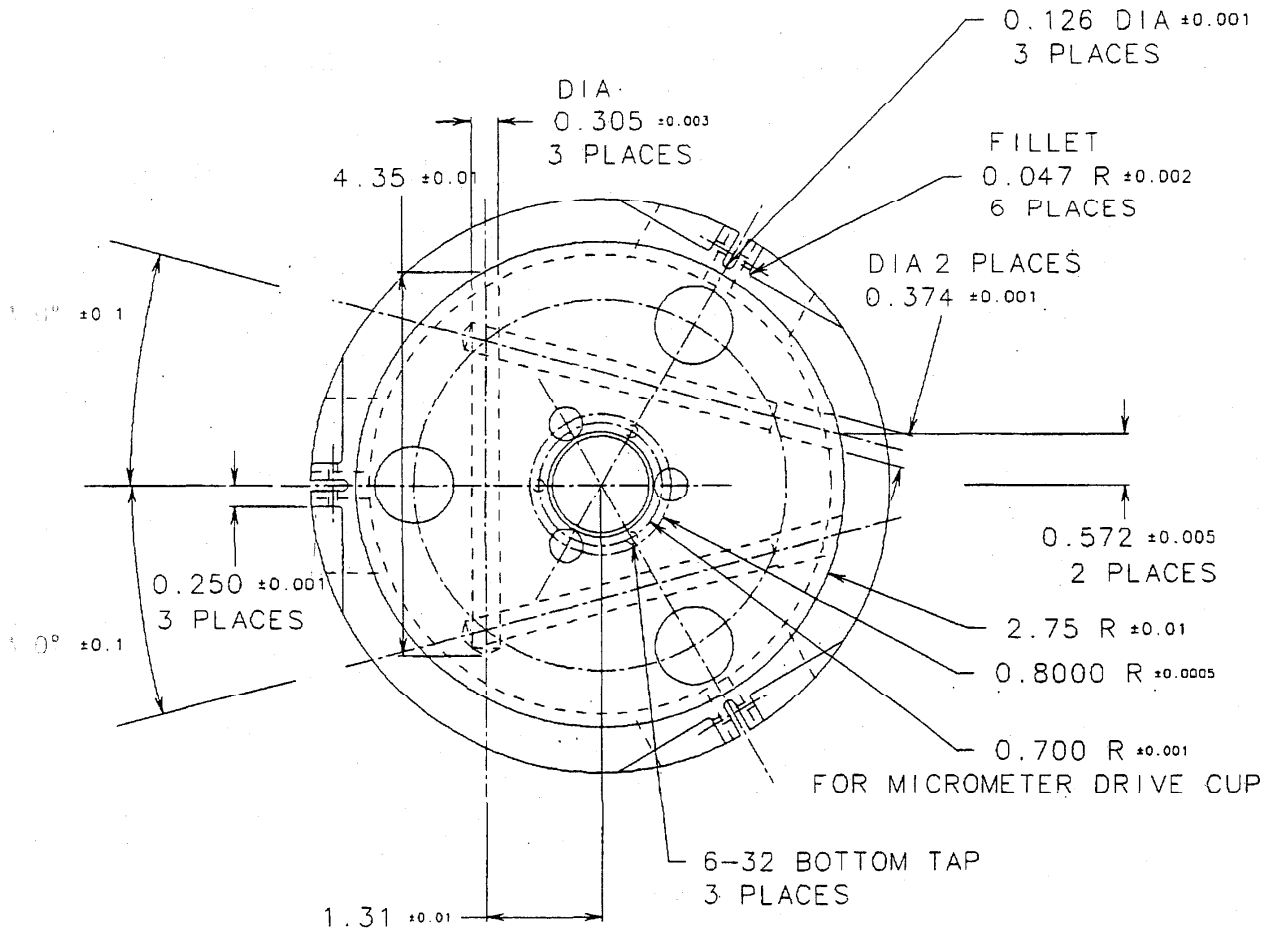


Figure C.16: Bottom view of tuner base.

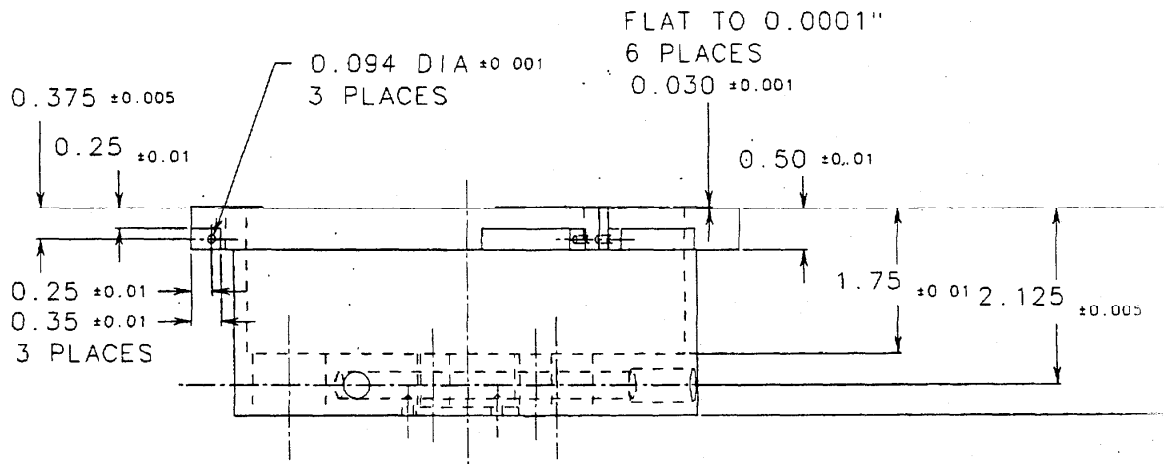


Figure C.17: Side view of tuner base.

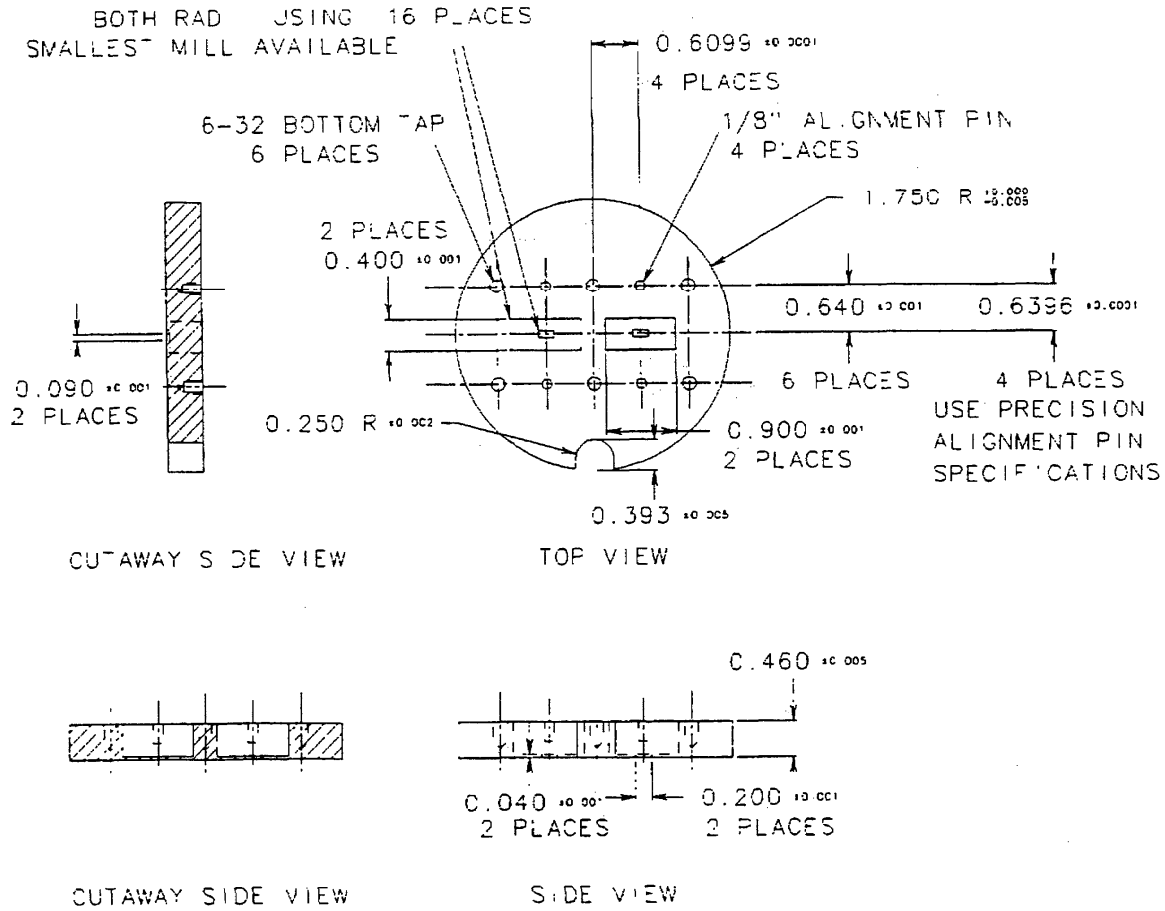


Figure C.18: Coupling endplate.

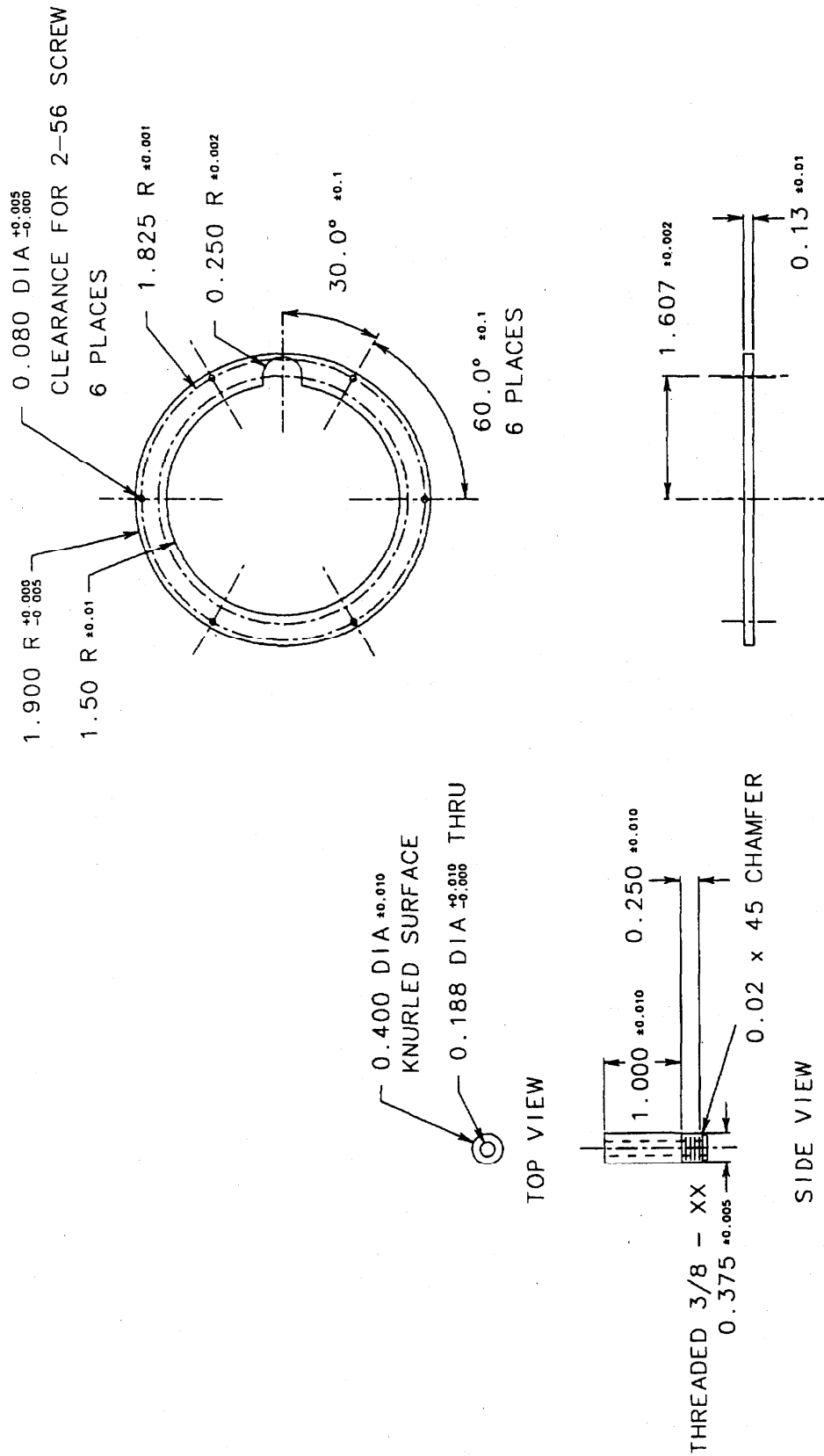


Figure C.19: Temperature-probe bolt and coupling-endplate pressure ring.

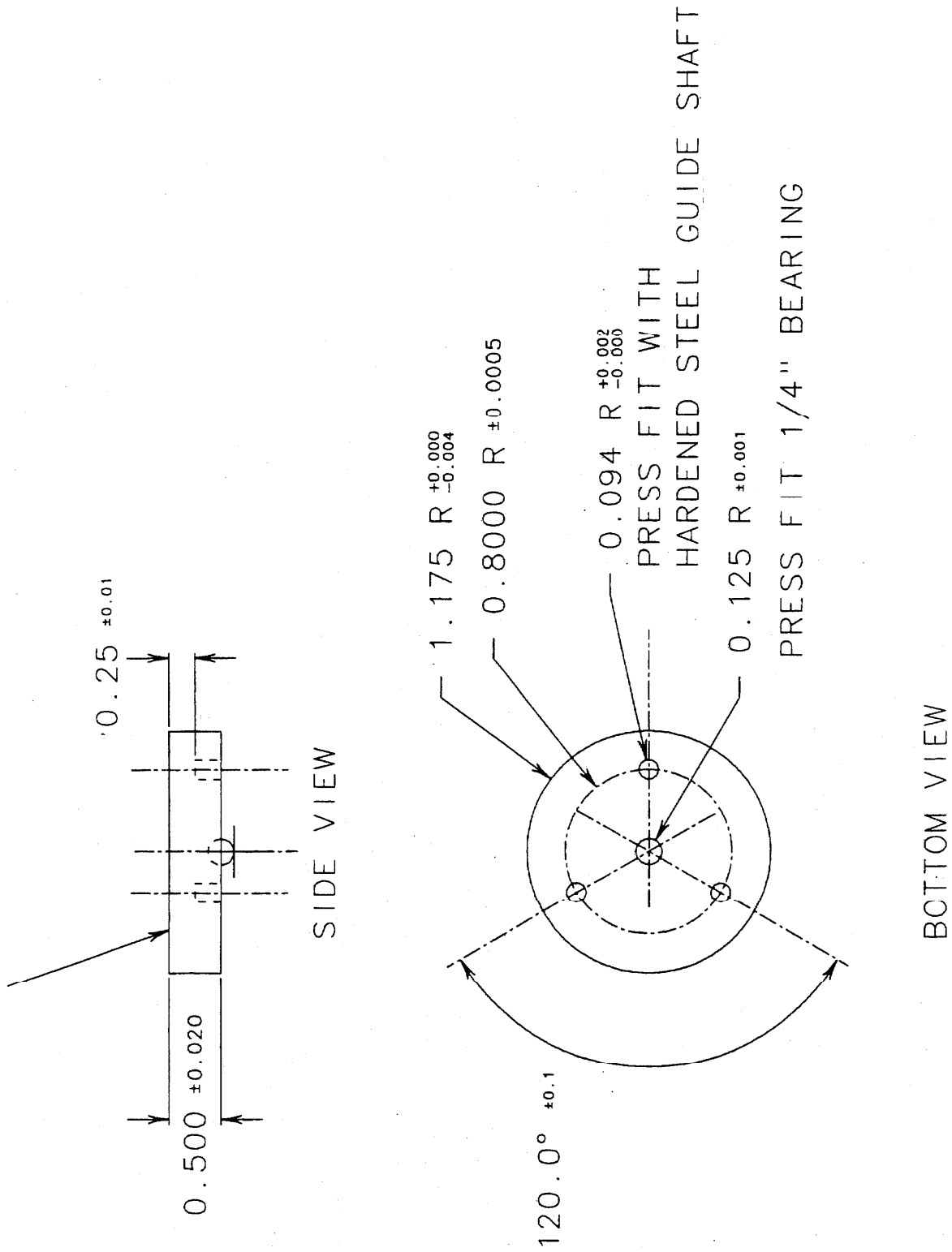
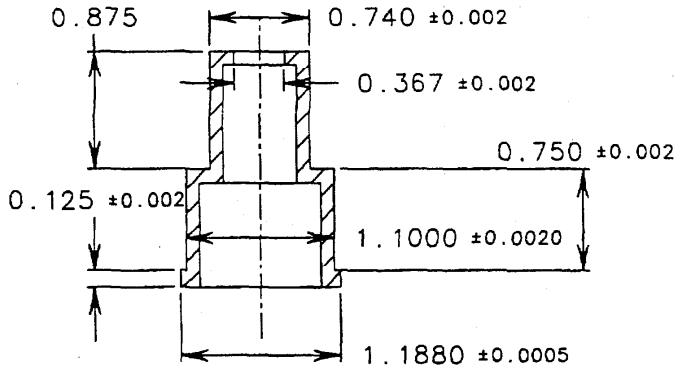
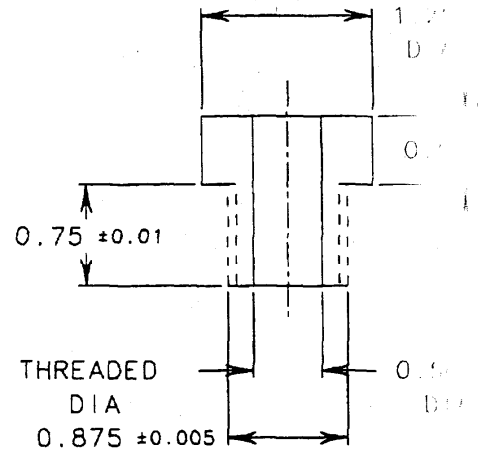


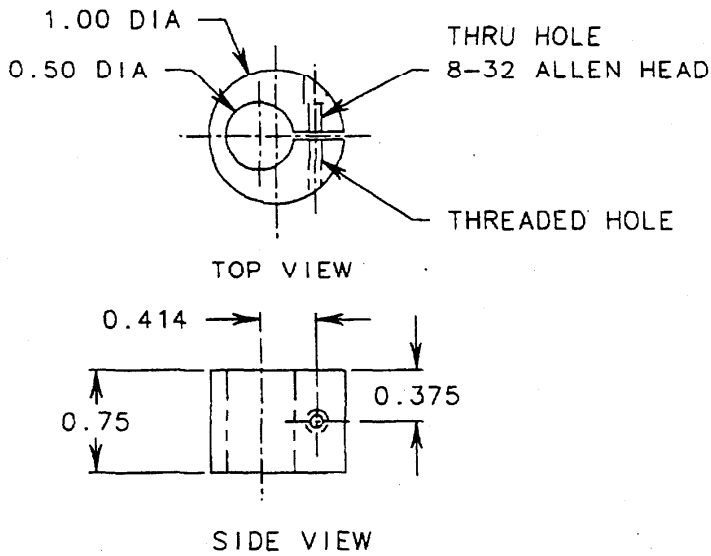
Figure C-20: Tuner-endplate.



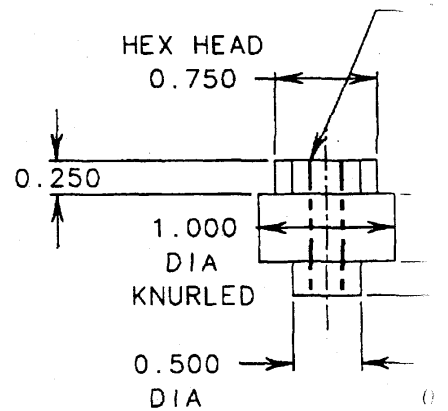
TYPICAL DESIGN FOR MICROMETER DRIVE HOLDER



BUSHING FOR TUNER BASE



BACKSTOP FOR TUNER ASSEMBLY



TUNER BASE TORQUE NUT

Figure C.21: Tuner assembly backstop, torque nut, micrometer-drive holder and sliding bushing.

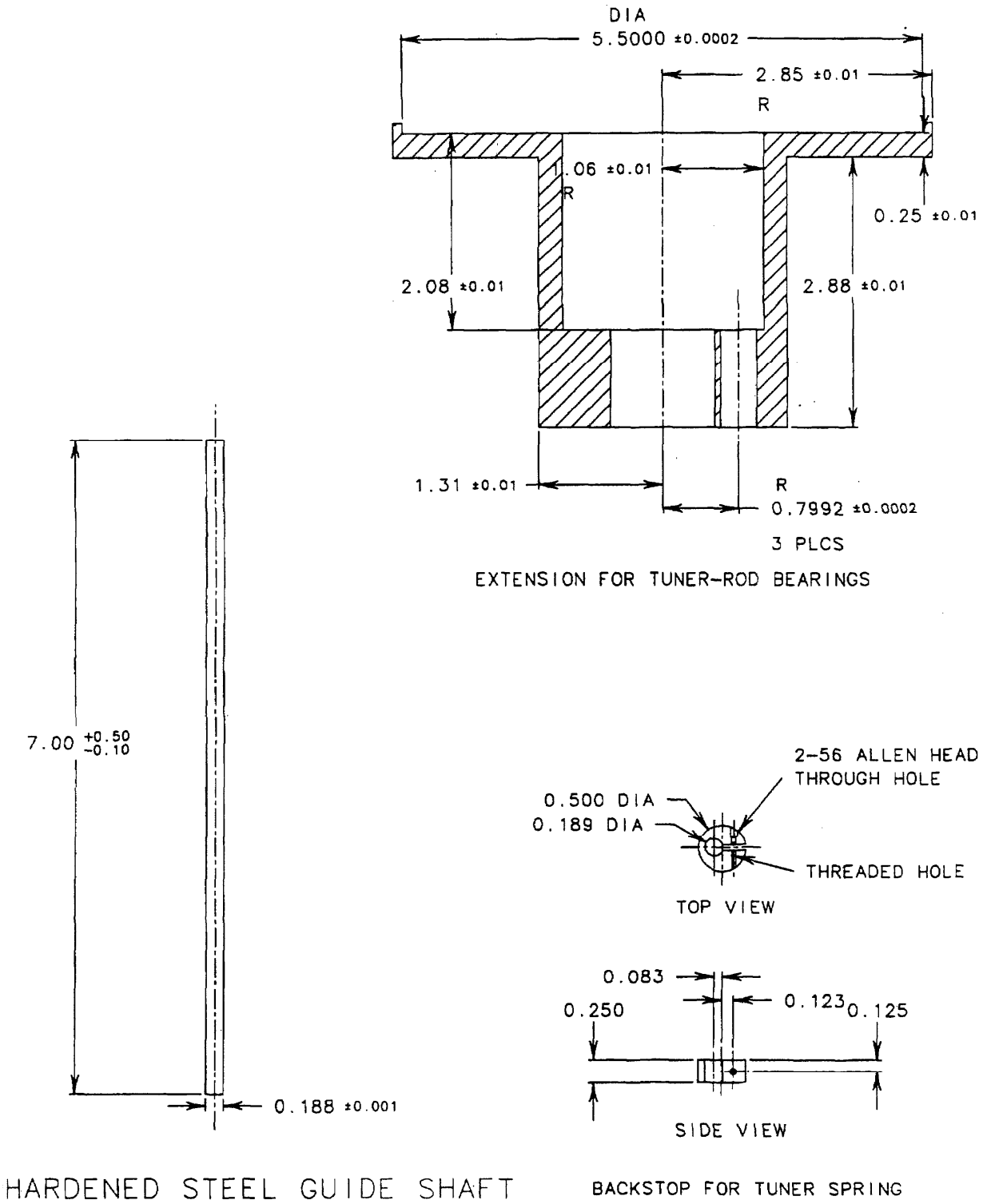


Figure C.22: Tuner guide shaft, spring backstop and extension.

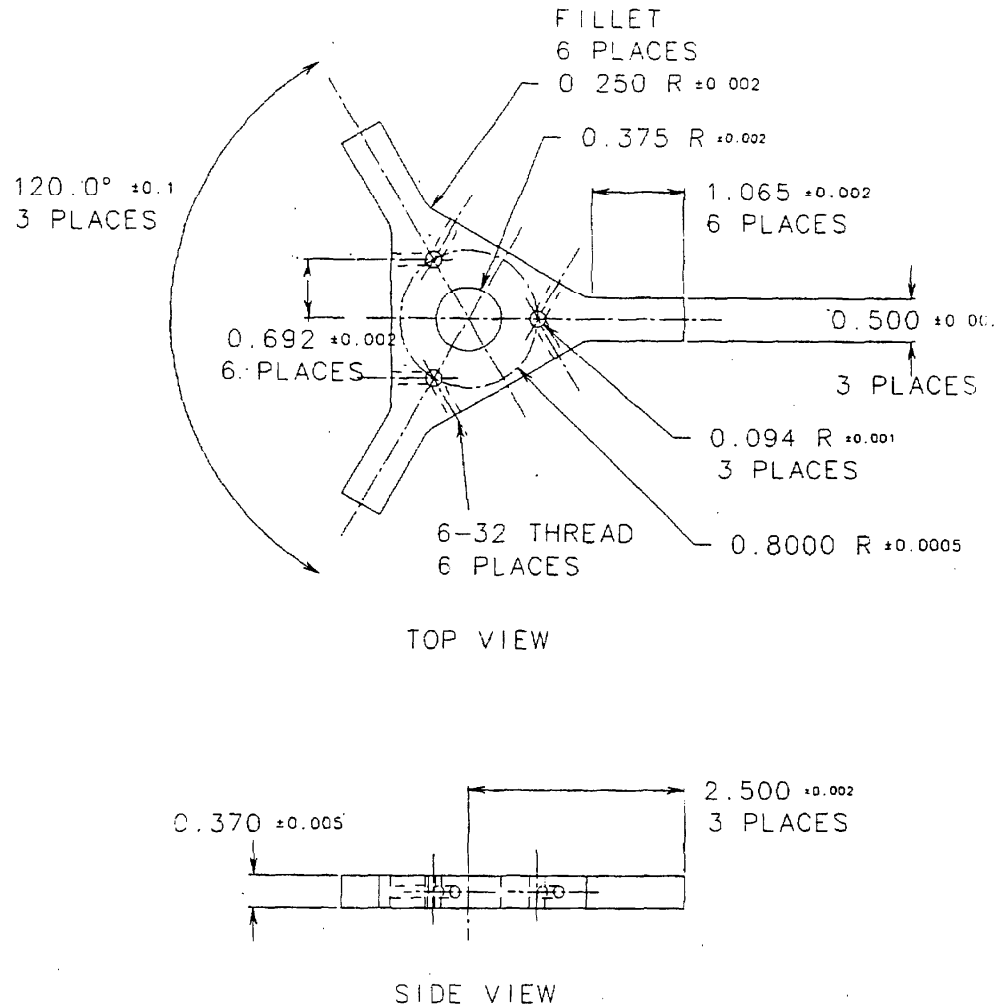


Figure C.23: Measurement-micrometer reference yoke.

Other resonator parts not shown are:

1. 1 ea. 60-mm diameter helical waveguide.
faced off perpendicular to axis.
outside diameter turned to be concentric with inner diameter.
ends flat and perpendicular to axis to within $3 \mu\text{m}$ (0.0001 in).
2. 1 ea. 4 in outside-diameter, 3.5 in inside-diameter acrylic plastic tubing with tangential inlet and outlet ducts at opposite ends.
3. 1 ea. plug for tuner-base water channel.
4. 6 ea. 2-56 machine screws for coupling-endplate pressure ring.
5. 3 ea. 2-56 set-screws for tuner-endplate.
6. 3 ea. Allen-head $\frac{1}{4}$ - 20 bolts to attach stand to tuner flange.
7. 1 ea. 10-32 knurled screw for micrometer attachment.
8. 2 ea. $\frac{1}{16}$ -in inside flange o-rings.
9. 2 ea. $\frac{1}{16}$ -in outside flange o-rings.
10. 1 ea. temperature probe o-ring.
11. 3 ea. $\frac{1}{4}$ -in diameter tie rods. Both ends threaded $\frac{1}{4}$ - 20. For attaching coupling and tuner flanges.
12. 3 ea. washers and $\frac{1}{4}$ - 20 tie rods.
13. 1 ea. base for resonator stand. 18-in diameter, 1-in thick aluminum with 3-in diameter center-hole and $3\frac{1}{2}$ -in diameter holes located 120° apart at a 2.1-in radius.
14. 3 ea. stands to connect from base to tuner flange. 12-in long, $\frac{1}{2} \pm 0.01$ in diameter stainless steel rods.

Appendix D

MEAS_RES Program Listing

```

105! PURGE "MEAS RES"
110! RE-STORE "MEAS_RES"
115 GOSUB Init_com
120 GOSUB Whatkindofmeas
125! Whatkindofmeas="LIST", "SWEEP", or "TRIGGER", "DISKDATA"
130 SELECT Whatkindofmeas$
135 CASE "LIST"
140 GOSUB Measurelist
145 CASE "SWEEP"
150 GOSUB Measuresweep
155 CASE "TRIGGER"
160 GOSUB Measuretrigger
165 CASE "LISTTRIGGER"
170 GOSUB Measurelisttrig
175 CASE "DUMMY"
180 GOSUB Measuredundata
185 CASE "DISKDATA"
190 GOSUB Measurediskdata
195 CASE "BAILOUT"
200 CASE ELSE
205 DISP "hey! ";Whatkindofmeas$;" isn't a measurement option. PROGRAM ST
=> OPED "
210 BEEP
215 STOP
220 END SELECT
225 GOTO End_program
230
235!
240!
245! Whatkindofmeas:
250 Prty=VAL(SYSTEM$( "SYSTEM PRIORITY" ))+1
255 CLEAR SCREEN
260 PRINT "This program allows you to read resonance data from a network anal
=> yzer."
265 PRINT "There are three ways for you to do this:"
270 PRINT
275 PRINT "1) Measure resonances at frequencies given in DATA statements."
280 PRINT "2) Edit line label 'Resonance_data' and enter the frequencies befo
re running."
285 PRINT "3) Sweep a band of frequencies, searching for resonances."
290 PRINT "4) Edit line label 'Init_sweep' to change the sweep parameters."
295 PRINT "5) Measure resonances that you set up manually on the network anal
=> yzer."
300 PRINT "Simply find a resonance on the screen and press a softkey to me
=>asure."
301 Menu: Reset keys=0
303 IF Printer ON THEN
304 PRINT TABXY(1,12);"Printer is ON "
305 ELSE
306 PRINT TABXY(1,12);"Printer is Off"
307 END IF
308 DISP "Which kind of measurement do you want to perform?"
310 ON KEY 1 LABEL "End_program",Prty GOSUB Bailout
315 ON KEY 2 LABEL "DATA LIST",Prty GOSUB Datalist
320 ON KEY 3 LABEL "SWEEP NWA",Prty GOSUB Sweep
325 ON KEY 4 LABEL "MEAS ONTRIGGER",Prty GOSUB Trigger
330 ON KEY 5 LABEL "DUMMY DATA",Prty GOSUB Dummydata
335 ON KEY 6 LABEL "DISK DATA",Prty GOSUB Diskdata
340 ON KEY 7 LABEL "LIST ->TRIGGER",Prty GOSUB Measurelisttrig
345 IF NOT Printer ON THEN
346 ON KEY 8 LABEL "Turn ONPrinter",Prty GOSUB Toggle_printer
ELSE
349 ON KEY 8 LABEL "TurnOffPrinter",Prty GOSUB Toggle_printer
350 END IF
351 Selected=0
352

```

```

355 LOOP
360 EXIT IF Selected
361 IF Reset_keys THEN GOTO Menu
365 END LOOP
370 OFF KEY
375 CLEAR SCREEN
380 RETURN
385!
390!
395!
396 Abort: Abort=1
397 RETURN
400 Bailout: Selected=1
405 Whatkindofmeas$="BAILOUT"
410 RETURN
415!
420 Datalist: Selected=1
425 Whatkindofmeas$="LIST"
430 RETURN
435!
440 Sweep: Selected=1
445 Whatkindofmeas$="SWEEP"
450 RETURN
455!
460 Trigger: Selected=1
465 Whatkindofmeas$="TRIGGE"
470 RETURN
475!
480 Dummydata: Selected=1
485 Whatkindofmeas$="DUMMY"
490 RETURN
495 Diskdata: Selected=1
500 Whatkindofmeas$="DISKDATA"
505 RETURN
510 Oelsub_modesub:
515 DISP "PLEASE WAIT WHILE THE ROUTINES ARE DELETED....."
520 DELSUB Resonance_data
525 DELSUB Empty_cavity
530 DELSUB Loaded_cavity
535 DELSUB Var_freq_empty
540 DELSUB Var_freq_loaded
545 DELSUB Var_len_empty
550 DELSUB Var_len_loaded
555 DELSUB Calc_c
560 DELSUB Calc_mode_freq
565 DISP
570 CLEAR SCREEN
575 RETURN
576!
577!
578!
579 Toggle_printer: Off key 8
580 IF Printer ON THEN
581 PRINT TABXY(1,12);"Printer is Off"
582 Printer_on=0
584 ELSE
585 PRINT TABXY(1,12);"Printer is ON"
586 Printer_on=1
588 END IF
589 Reset_keys=1
590 RETURN
591!
592!
593!
594 Measurelist: Off key
595 Sss=1
600

```

```

605 Speed=0
610 MAT Te01= (0)
615 MAT Te02= (0)
620 MAT Te12= (0)
625 MAT Te13= (0)
630 MAT Freq= (0)
635 GOSUB Resonance_data
640 FOR I=1 TO Num_Peaks
645 GOSUB Get_Resonance
650 CALL Zoom_on_peak
655 GOSUB Measure_sparms
660 CALL Optimizedisplay
665 GOSUB Set_startstop
670 CALL Fit_sparms
675 CALL Phase_delay
680 CALL Fit_sparms
685 NEXT I
690 IF Printer_on THEN OUTPUT Print_addr;CHR$(12);
695 RETURN
700
705
710
715 Measurelisttrig:OFF KEY
720 Sse=1
725 RESTORE Resonance_data
730 READ Num_peaks
735 FOR I=1 TO Num_peaks
740 READ Freq(I)
745 NEXT I
750 FOR I=1 TO Num_peaks
755 GOSUB Give_display
760 CALL Zoom_on_peak
765 GOSUB Measure_sparms
770 CALL Optimizedisplay
775 GOSUB Set_startstop
780 CALL Fit_sparms
785 NEXT I
790 IF Printer_on THEN OUTPUT Print_addr;CHR$(12);
795 RETURN
800
805
810
815 Give display: !
820 CALL Set_nwa("SWEEP=RAMP", "S=21")
825 CALL Set_nwa("SPAN=1E5", "SCALE=10", "CENTER=gVAL$(Freq(1))")
830 CALL Set_nwa("REFV=-40", "AVER=OFF", "NUMG=1")
835 CALL Optimizedisplay
840 CALL Set_nwa("SWEEP=CONT")
845 Prty=VAL(SYSMS("SYSTEM PRIORITY"))+1
850 ON KEY 6 LABEL "READY TO MEAS", Prty GOTO Get_back
855 Sitnspin:GOTO Sitnspin
860 Get_back:OFF KEY
865 CALL Set_nwa("SPAN=3E7")
870 RETURN
875
880
885
890 Get resonance: !
895 CALL Set_nwa("SWEEP=CONT", "SWEEP=RAMP", "S=21")
900 CALL Set_nwa("SPAN=3E7", "SCALE=5", "CENTER=gVAL$(Freq(1))")
905 CALL Set_nwa("REFV=-40", "AVER=16", "NUMG=17")
910 RETURN
915
920 Find resonance: !
925 CALL Set_nwa("SWEEP=CONT", "SWEEP=RAMP", "S=21")
930
935 CALL Set_nwa("SPAN=3E7", "SCALE=5", "CENTER=gVAL$(Freq(1))")
940 CALL Set_nwa("REFV=-40", "AVER=16", "NUMG=17")
945 Set_nwa("MARK=CONT", "MARK=1", "MARK=MAX")
950 Read_nwa("MARK=1")
955 Delta=3
960 Set_nwa("DEL REF=1", "MARKER TARGET=gVAL$(DelTa)")
965 Set_nwa("MARK=2", "MARK=MAX", "MARKER=TARGET", "MARKER=SEARCH LEFT")
970 Set_nwa("MARK=3", "MARK=MAX", "MARKER=TARGET")
975 Read_nwa("MARK=3", "MARK=MAX", "MARKER=TARGET")
980 Set_nwa("DELTA Off=")
985 IF ABS(Marker(2,1))>5000 AND ABS(Marker(3,1))>5000 THEN
990 IF ABS(Marker(2,1))<ABS(Marker(3,1)) THEN
995 SpanF2=ABS(Marker(3,1))
ELSE
SpanF2=ABS(Marker(2,1))
1000 ELSE
Span=2*ABS(Marker(2,1))
1005 END IF
1010 CALL Zoom_on_peak
1015 ELSE
Span=9999999
1020 BEEP
1025 PRINT "Span cannot be determined from marker information."
1030 PRINT "No resonance was found at ", Freq(1)/1.E+9, " GHz."
1035 IF Printer_on THEN OUTPUT Print_addr; "No resonance was found at ", Freq(1)/1.E+9, " GHz."
1040 => q(1)/1.E+9, " GHz."
1050 DISP "No resonance was found at ", Freq(1)/1.E+9, " GHz."
1055 PAUSE
1060 END IF
1065 RETURN
1070
1075
1080
1085 Measuresweep:OFF KEY
1090 CALL Sweep_nwa
1095 RETURN
1100
1105
1110
1115 Measuretrig:OFF KEY
1120 Prty=VAL(SYSMS("SYSTEM PRIORITY"))+1
1125 DISP "When you have a resonance on the network analyzer CRT, press 'READ
PEAK'."
1130 ON KEY 0 LABEL "END PROG", Prty GOTO Triggerdone
1135 ON KEY 1 LABEL "READ S11", Prty GOSUB Read_s11res
1140 ON KEY 2 LABEL "READ S21", Prty GOSUB Read_s21res
1145 ON KEY 3 LABEL "READ S12", Prty GOSUB Read_s12res
1150 ON KEY 4 LABEL "READ S22", Prty GOSUB Read_s22res
1155 ON KEY 5 LABEL "READ S11/S22", Prty GOSUB Read_s11s22res
1160 ON KEY 6 LABEL "ZOOM S11", Prty GOSUB Zoomread_s11res
1165 ON KEY 7 LABEL "ZOOM S21", Prty GOSUB Zoomread_s21res
1170 ON KEY 8 LABEL "ZOOM S12", Prty GOSUB Zoomread_s12res
1175 ON KEY 9 LABEL "ZOOM S22", Prty GOSUB Zoomread_s22res
1180 Reset_keys=0
1185 LOOP
1190 IF Reset_keys THEN GOTO Measuretrigger
1195 END LOOP
1200 Triggerdone:OFF KEY
1205 RETURN
1210
1215
1220
1225 Measurediskdata:OFF KEY
1230 Get_data: !
1235 DISP "Diskdata is in new or old format?"
1240 Prty=VAL(SYSMS("SYSTEM PRIORITY"))+1

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1245 ON KEY 0 LABEL "ABORT THIS" Prty GOSUB Done_return
1250 ON KEY 2 LABEL "OLD SPARMS" Prty GOSUB Get_old_sparms
1255 ON KEY 3 LABEL "NEW SPARMS" Prty GOSUB Get_new_sparms
1260 Done=0
1265 Reset_keys=0
1270 REPEAT
1275 IF Reset_keys THEN GOTO Get_data
1280 UNTIL Done
1285 RETURN
1290 Done_return:Done=1
1300 RETURN
1310 ///////////////////////////////////////////////////////////////////
1315 ///////////////////////////////////////////////////////////////////
1320 Get_old_sparms:OFF KEY
1325 CALL Load_sparms
1330 CALL Phase_delay
1335 RESTORE Diskdatas
1340 Speed=0
1345 Select=1 IS21
1350 Start_Index=1
1355 Stop_Index=Dcount
1360 GOSUB Diskdata_keys
1365 Done=1
1370 RETURN
1375 ///////////////////////////////////////////////////////////////////
1380 ///////////////////////////////////////////////////////////////////
1385 ///////////////////////////////////////////////////////////////////
1390 Get_new_sparms:OFF KEY
1395 ALLOCATE Descs(160)
1400 REPEAT
1405 LINPUT "Enter a description of this data. When finished enter a null
=> string," Descs
1410 PRINT Descs
1415 IF Printer_on THEN OUTPUT Print_addr;Descs
1420 UNTIL Descs=""
1425 DEALLOCATE Descs
1430 N=6
1435 OUTPUT 2 USING "k,#":N
1440 INPUT "Please enter the number of files",N
1441 ON ERROR GOSUB Change_disk
1443 RESTORE Mode_numbers
1444 READ Num_settings
1445 FOR J=1 TO Num_settings
1446 READ Micr_Nummodes
1447 IF Printer_on THEN OUTPUT Print_addr;"Micr:","Micr:"," Nummodes:","Nummo
=> des," 30 Nov 1990 Test of micrometer travel accuracy"
1448 PRINT "Micr:","Micr:"," Nummodes:"," 30 Nov 1990 Test of micr
ometer travel accuracy"
FOR I=1 TO Nummodes
1449 Filenames=VAL$(Micr)&"30Nov90"&VAL$(I)
1455 DiskDrives="1,1400,1"
1460 CALL New_load_sparms
1465 Sselect=1
1470 Start_Index=1
1475 Stop_Index=Dcount
1480 CALL Fit_sparms
1485 Peak(1,1)=A(4)
1490 Peak(1,2)=SORT(A(1)-2*A(2)-2)
1495 Peak(1,3)=A(4)/A(3)
1500 NEXT I
1505 IF Printer_on THEN OUTPUT Print_addr;CHR$(12);
N=Nummodes
1506 ALLOCATE Basket(N,2),ld$(40)
1510 MAT Basket(*,1)= Peak(1:N,1)
1515

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```

1520 MAT Basket(*,2)= Peak(1:N,3)
1525 CLEAR SCREEN
1530 Filenames=VAL$(Micr)&"30Nov90"&VAL$(J)
1535 J=J+1
1540 CALL Select_disk
1545 IF DiskDrives="NO DISK" THEN GOTO Skip_peak_save
1550 CALL Enterfilename("ABORT","SAVING resonance parameters in a file")
1555 IF Filenames="" THEN GOTO Skip_peak_save
1560 CALL Enter_id(lds,"resonance parameter file")
1565 IF lds="" THEN GOTO Skip_peak_save
1570 CALL Save file(Basket(*),Nummodes,ld$(*)
1575 Skip_peak_save:DEALLOCATE lds,Basket(*)
1580 NEXT J
1585 OFF ERROR
1590 Reset_keys=1
1595 RETURN
1600 Change_disk:PRINT "Please insert next disk"
1605 DISP ERRMS
1610 BEEP
1615 PAUSE
1620 RETURN
1625 ///////////////////////////////////////////////////////////////////
1630 ///////////////////////////////////////////////////////////////////
1635 ///////////////////////////////////////////////////////////////////
1640 Diskdata_keys:OFF KEY
1645 GOSUB Diskdata_menu
1650 DISP "Which S-parameter would you like to fit?"
1655 Prty=VAL$(SYSTEMS("SYSTEM PRIORITY"))+1
1660 ON KEY 1 LABEL "Bailout" Prty GOSUB Bailout
1665 ON KEY 6 LABEL "FIT S11" Prty GOSUB Fit_s11
1670 ON KEY 7 LABEL "FIT S12" Prty GOSUB Fit_s12
1675 ON KEY 8 LABEL "FIT S21" Prty GOSUB Fit_s21
1680 ON KEY 9 LABEL "FIT S22" Prty GOSUB Fit_s22
1685 ON KEY 3 LABEL "SAVE SPARMS" Prty GOSUB Save_new_sparms
1690 ON KEY 4 LABEL "START INDEX" Prty GOSUB Fix_start_index
1695 ON KEY 5 LABEL "STOP INDEX" Prty GOSUB Fix_stop_index
1700 Done=0
1705 Selected=0
1710 LOOP
1715 IF Done THEN GOTO Diskdata_keys
1720 EXIT IF Selected
1725 END LOOP
1730 OFF KEY
1735 CLEAR SCREEN
1740 RETURN
1745 Fit_s11: I
1750 SSS=2
1755 SparmS="S11"
1760 MAT Sparm= S11
1765 MAT Frequency= Freq
1770 CALL Fit_sparms
1775 Done=1
1780 RETURN
1785 Fit_s12: I
1790 SSS=1
1795 SparmS="S12"
1800 MAT Sparm= S12
1805 MAT Frequency= Freq
1810 CALL Fit_sparms
1815 Done=1
1820 RETURN
1825 Fit_s21: I
1830 SSS=1
1835 SparmS="S21"
1840 MAT Sparm= S21
1845 MAT Frequency= Freq
1850 CALL Fit_sparms
1855 Done=1
1860 RETURN
1865 Fit_s22: I
1870 SSS=1
1875 SparmS="S22"
1880 MAT Sparm= S22
1885 MAT Frequency= Freq
1890 CALL Fit_sparms
1895 Done=1
1900 RETURN

```



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2030 END IF
2031 Reset_keys=1
2032 RETURN
2033
2034
2035
2036 Diskdata_menu:PRINT TAB(1,1)
2037 PRINT "S PARAMETER FITTING FROM DISK DATA"
2038 PRINT
2039 PRINT USING Image1;Filename$&Diskdrives$
2040 PRINT USING Image2;Title$
2041 PRINT USING Image3;DATE$(DATE$(Measure_time$))
2042 PRINT USING Image3;Dcount
2043 PRINT
2044 PRINT USING Image4;Start_index
2045 PRINT USING Image5;Stop_index
2046 PRINT
2047 Image1: IMAGE:FILENAME : "K
2048 Image2: IMAGE:DESCRIPTION : "K
2049 Image3: IMAGE:DATAACCOUNT : "3D
2050 Image4: IMAGE:START INDEX : "3D
2051 Image5: IMAGE:STOP INDEX : "3D
2052 Image6: IMAGE:DATE : "K
2053 RETURN
2054
2055
2056 Measuredata:OFF KEY
2057 Span=2.E+5
2058 Fb=(1.-200E+10)
2059 00=90000
2060 01=70000
2061 G0=F0/D0
2062 G1=F0/D1
2063 Mag=7
2064 Degress=58
2065 Num_points=50
2066 Start_index=1
2067 Stop_index=50
2068 ALLOCATE X(Num_points),COMPLEX S21(Num_points),
2069 FOR I=1 TO Num_points
2070 X(I)=Span*I/Num_points+F0.(Span/2)
2071 Frequency(I)=X(I)
2072 012f=2*q1*(X(I)-F0)/F0
2073 P=CMPLX(COS(Degress),SIN(Degress))
2074 S21(I)=Mag*(1-G0/G1)*P/CMPLX(1,2*(X(I)-F0)/G1)
2075 S11(I)=Mag*CMPLX(01/00+012f*q12f,012f*(1-01/00))*P/(1+012f*012f)
2076 S21(I)=(1-(01/00))/CMPLX(1,2*q1*(F(I)-F0)/F0)
2077 S11(I)=Mag*CMPLX(01/00,2*q1*(F(I)-F0)/F0)/CMPLX(1,2*q1*(F(I)-F0)/F0)*
2078 Rndmag*(001*(RND-.5)
2079 Rndphase=.001*(RND-.5)
2080 Sparm(I)=S21(I)+Rndmag*(CMPLX(COS(Rndphase),SIN(Rndphase)))
2081 NEXT I
2082 CALL Fit_sparms
2083 PRINT USING "2(K,MD,3DE,2X),K,MD,9DE";"00=";A(1),"01=";A(2),"F0=";A(3)
2084 PRINT USING "2(K,MD,3DE,2X),K,MD,9DE";"000=";U0,"001=";U0,"010=";U0,"011=";U0,"012=";U0,"013=";U0
2085 RETURN
2086
2087
2088
2089 Read_s11s22res:OFF KEY
2090 INPUT "ENTER THE NUMBER OF AVERAGES",Avs
2091 OUTPUT 716;"AVERON "&VAL$(Avs)
2092
2093
2094 OUTPUT 716;"S11:"
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2116 Read_allsparms:OFF KEY
2117 GOSUB Read_s11res
2118 GOSUB Read_s21res
2119 GOSUB Read_s12res
2120 GOSUB Read_s22res
2121 GOSUB Save_new_sparms
2122 Reset_keys=1
2123 RETURN
2124
2125
2126
2127 Zoomread_s11res:OFF KEY
2128 CALL Zoom_on_peak
2129 Read_s11res:OFF KEY
2130 Sparm="S11"
2131 Sss=3
2132 GOSUB Measure_sparms
2133 CALL Optimizedisplay
2134 GOSUB Set_startstop
2135 CALL Fit_sparms
2136 GOSUB Give_results
2137 Reset_keys=1
2138 RETURN
2139
2140
2141
2142 Zoomread_s21res:OFF KEY
2143 CALL Zoom_on_peak
2144 Read_s21res:OFF KEY
2145 Sparm="S21"
2146 Sss=1
2147 GOSUB Measure_sparms
2148 CALL Optimizedisplay
2149 GOSUB Set_startstop
2150 CALL Set_nwa("SPAN=1000MHz","SWEEP=RAMP","AVER-OFF","SWEEP=STING")
2151 CALL Set_nwa("SWEEP=CONT","AUTOSCALE=","LOCAL=")
2152 CALL Fit_sparms
2153 GOSUB Give_results
2154 Reset_keys=1
2155 RETURN
2156
2157
2158
2159 Zoomread_s12res:OFF KEY
2160 CALL Zoom_on_peak
2161 Read_s12res:OFF KEY
2162 Sparm="S12"
2163 Sss=1

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2164 GOSUB Measure_sparms
2165 CALL Optimizedisplay
2166 CALL Fit_sparms
2167 GOSUB Set_startstop
2168 GOSUB Give_results
2169 Reset_keys=1
2170 RETURN
2171
2172
2173
2174 Zoomread_s22res:OFF KEY
2175 CALL Zoom_on_peak
2176 Read_s22res:OFF KEY
2177 Sparm8="S22"
2178 Sss=3
2179 GOSUB Measure_sparms
2180 CALL Optimizedisplay
2181 CALL Fit_sparms
2182 GOSUB Give_results
2183 Reset_keys=1
2184 RETURN
2185
2186
2187
2188 Measure_sparms:
2189 SELECT Sparm8
2190 CASE "S11"
2191   Set nwa("S=11", "LOG MAG=")
2192 CASE "S21"
2193   Set nwa("S=21", "LOG MAG=")
2194 CASE "S12"
2195   Set nwa("S=12", "LOG MAG=")
2196 CASE "S22"
2197   Set nwa("S=22", "LOG MAG=")
2198 END SELECT
2199 CALL Set_nwa("SWEEP=STEP")
2200 CALL Set_nwa("NUMBER OF GROUPS=1")
2201 CALL Read_nwa("S=")
2202 RETURN
2203
2204
2205
2206 Set_startstop:
2207 Start_index=1
2208 Stop_index=0count
2209 RETURN
2210
2211
2212
2213
2214 Give_results:
2215 SELECT Sss
2216 CASE 1
2217   IF Printer_on THEN
2218     PRINTER IS Print_addr
2219     PRINT USING "K,D,SUESZ,K,D,90E",resonance 3dB Span: "A(
2220     4)/A(3)," Unloaded resonant frequency: "A(4)+A(4)/A(3)+2)
2221     PRINT
2222     PRINTER IS 1
2223   END IF
2224 END SELECT
2225 RETURN
2226
2227 Send_to_local:
2228
2164 RETURN
2165
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2284 COM /Modes/ REAL Te01(50,2),Te02(50,2)
2285 COM /Modes/ REAL Te12(50,2),Te13(50,2)
2286 COM /Output/ INTEGER Print_addr,Plotter,Printer_on,Plotter_on
2287 INTEGER Num_settings,Mtr, NumModes,Ssave
2288 !
2289 Init_var: !
2290 Speed=1
2291 Printers=701
2292 Print_addr=Printer
2293 Plotter=705
2294 Printer_on=1
2295 Plotter_on=0
2296 COMPLEX Cnum,P
2297 Ssave=0
2298 RETURN
2299
2300 !
2301 Data: !
2302 Resonance_data: !
2303 DISP "PLEASE WAIT WHILE ROUTINES ARE LOADED....."
2304 LOADSUB Resonance_data FROM "MODESUB: 1400,0,3"
2305 LOADSUB Empty_cavity FROM "MODESUB: 1400,0,3"
2306 LOADSUB Loaded_cavity FROM "MODESUB: 1400,0,3"
2307 LOADSUB Var_freq_empty FROM "MODESUB: 1400,0,3"
2308 LOADSUB Var_freq_loaded FROM "MODESUB: 1400,0,3"
2309 LOADSUB Var_len_empty FROM "MODESUB: 1400,0,3"
2310 LOADSUB Var_len_loaded FROM "MODESUB: 1400,0,3"
2311 LOADSUB Calc_c FROM "MODESUB: 1400,0,3"
2312 LOADSUB Calc_mode_freq FROM "MODESUB: 1400,0,3"
2313 DISP Resonance_data
2314 I=1
2315 J=0
2316 WHILE Te01(I,1)=0
2317 I=I+1
2318 END WHILE
2319 WHILE Te01(I,1)<0
2320 J=J+1
2321 Freq(J)=Te01(I,1)
2322 I=I+1
2323 END WHILE
2324 Num_peaks=J
2325 DISP "PLEASE WAIT WHILE ROUTINES ARE DELETED....."
2326 DELSUB Resonance_data
2327 DELSUB Empty_cavity
2328 DELSUB Loaded_cavity
2329 DELSUB Var_freq_empty
2330 DELSUB Var_freq_loaded
2331 DELSUB Var_len_empty
2332 DELSUB Var_len_loaded
2333 DELSUB Calc_c
2334 DELSUB Calc_mode_freq
2335 DISP
2336 RETURN
2337
2338 !
2339 End_program: !
2340 DISP "the program has ended."
2341
2342 Mode_numbers:DATA 29,0,16,2,16,4,16,6,16,8,17,10,16,12,15,14,16,16,17,18,14
2343 ,20,14,22,14,24,15,21,15,19,14,17,16,15,15,13,16,11,16
2344 DATA 6,15,7,15,5,16,3,16,1,17,2,16,4,16,6,16,8,16
2345 =>
2346 END
2347 !
2348 ! *****

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```

2349 ! SUB Set_nwa(Cmd1$,OPTIONAL Cmd2$,Cmd3$,Cmd4$,Cmd5$,Cmd6$)
2350 Set_nwa: !
2351 OPTION BASE 1
2352 ASSIGN @nwa TO 715
2353 ALLOCATE C$(1:6)$(5),Hp_ib$(35)
2354 ON NPAT GOTO C1,C2,C3,C4,C5,C6
2355 C$(6)=Cmd6$
2356 C6:
2357 C$(5)=Cmd5$
2358 C5:
2359 C$(4)=Cmd4$
2360 C4:
2361 C$(3)=Cmd3$
2362 C3:
2363 C$(2)=Cmd2$
2364 C2:
2365 C$(1)=Cmd1$
2366 C1:
2367 FOR I=1 TO NPAT
2368 Functions=C$(I)$(1),POS(C$(I),"=")-1
2369 Values=C$(I)$(2),POS(C$(I),"=")+1
2370 SELECT Function$
2371 CASE "PRES" "PRESET"
2372 Hp_ib$="PRES:"
2373 CASE "START" "START FREQ"
2374 Hp_ib$="START"Value$&","
2375 CASE "STOP" "STOP FREQ"
2376 Hp_ib$="STOP"Value$&","
2377 CASE "CENTER" "CF"
2378 SELECT Value$
2379 CASE "MARK" "MARKER"
2380 Hp_ib$="CENT"Value$&","
2381 CASE ELSE
2382 Hp_ib$="CENT"Value$&","
2383 END SELECT
2384 CASE "SPAN" "SPN"
2385 SELECT Value$
2386 CASE "MARK" "MARKER"
2387 Hp_ib$="SPAN"EQUA;"
2388 CASE ELSE
2389 Hp_ib$="SPAN"Value$&","
2390 END SELECT
2391 Marker: !
2392 CASE "MARKER" "MARK"
2393 SELECT Value$
2394 CASE "MAX" "MAXIMUM" "MAXI"
2395 Hp_ib$="MARKMAXI"Value$&","
2396 CASE "MIN" "MINIMUM" "MINI"
2397 Hp_ib$="MARKMINI"Value$&","
2398 CASE "TAR" "TARGET"
2399 Hp_ib$="MARKTARG;"
2400 CASE "CONT" "CONTINUOUS"
2401 Hp_ib$="MARKCONT;"
2402 CASE "DISC" "DISCRETE"
2403 Hp_ib$="MARKDISC;"
2404 CASE "SEARCH RIGHT"
2405 Hp_ib$="SEAR;"
2406 CASE "SEARCH LEFT"
2407 Hp_ib$="SEAR;"
2408 CASE ELSE
2409 Hp_ib$="SEAL;"
2410 END SELECT
2411 CASE "MARKER TARGET" "TARGET VALUE" "TARV"
2412 Hp_ib$="TARV"Value$&","
2413 CASE "DELTA REF" "DELTA REFERENCE" "DEL REF" "DELR"
2414 Hp_ib$="DELREF"Value$&","
2415 CASE "DEL OFF" "DELTA OFF" "DELO"
2416 Hp_ib$="DELO;"
2417 CASE "AVER" "AVERAGING" "AVERAGE" "AVG"
2418

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2415 IF Value$="OFF" THEN
2416   Hp_ib$="AVEROFF";"
2417 ELSE
2418   Hp_ib$="AVERON"&Value$&";"
2419 END IF
2420 CASE "NUMG", "NUM GROUPS", "NUMBER OF GROUPS"
2421   Hp_ib$="NUMG"&Value$&";"
2422 CASE "REF POS", "REF POSN", "REFERENCE POSITION", "REFP"
2423   Hp_ib$="REFP"&Value$&";"
2424 CASE "REF VAL", "REFERENCE VALUE", "REF VALUE", "REFV"
2425   Hp_ib$="REFV"&Value$&";"
2426 CASE "WAIT"
2427   Hp_ib$="WAIT";"
2428 CASE "UP", "STEP UP"
2429   Hp_ib$="UP";"
2430 CASE "DOWN", "STEP DOWN"
2431   Hp_ib$="DOWN";"
2432 CASE "LOGM", "LOG MAG"
2433   Hp_ib$="LOGM";"
2434 CASE "MAG", "CARTESIAN", "X-Y"
2435   Hp_ib$="MAG";"
2436 CASE "PHASE"
2437   Hp_ib$="PHAS";"
2438 CASE "SMITH CHART", "SMITH"
2439   Hp_ib$="SMIC";"
2440 CASE "SWEEP"
2441   SELECT Value$
2442     CASE "CONT", "CONTINUOUS"
2443       Hp_ib$="CONT";"
2444     CASE "SING", "SINGLE"
2445       Hp_ib$="SING";"
2446     CASE "STEP"
2447       Hp_ib$="STEP";"
2448     CASE "RAMP"
2449       Hp_ib$="RAMP";"
2450     CASE "HOLD"
2451       Hp_ib$="HOLD";"
2452     CASE "AUTO"
2453       Hp_ib$="AUTO";"
2454     CASE "N", "N"
2455       Hp_ib$="FUNCTIONS"&Value$&";"
2456     CASE "SCALE", "SCALE"
2457       Hp_ib$="SCALE"&Value$&";"
2458     CASE "AUTO SCALE", "AUTO"
2459       Hp_ib$="AUTO";"
2460     CASE "LOCAL"
2461       LOCAL @NWA
2462     CASE ELSE
2463       PRINT Functions;" IS NOT DEFINED IN SET_NWA"
2464     END SELECT
2465   IF Hp_ib$="" THEN OUTPUT @NWA;Hp_ib$
2466 NEXT I
2467 ASSIGN @NWA TO *
2468 SUBEND
2469 *****
2470 SUB Read_nwa(Cmd$$, OPTIONAL Cmd2$, Cmd3$, Cmd4$, Cmd5$, Cmd6$)
2471   OPTION BASE 1
2472   COM /NWA_dbr/ Start_freq, Stop_freq, Marker(*)
2473   COM /S_array/ INTEGER %count, S_val(id), REAL S(801, 5, 2)
2474   DIM Sparms(1), Swems(20)
2475   Read_nwa: 1
2476   ASSIGN @NWA_data1 TO 716;FORMAT ON
2477   ASSIGN @NWA_data2 TO 716;FORMAT OFF
2478   REAL Freq
2479   INTEGER Preamble_bytes
2480   ALLOCATE C$(1:6) [35], Hp_ib$(35)
2481   ON NPAR GOTO C1, C2, C3, C4, C5, C6
2482   C$(6)=Cmd$
2483   C$(5)=Cmd$
2484   C$(4)=Cmd$
2485   C$(3)=Cmd$
2486   C$(2)=Cmd$
2487   C$(1)=Cmd$
2488   FOR I=1 TO NPAR
2489     Functions=C$(1) [1, POS(C$(1), "=")-1]
2490     Value$=C$(1) [POS(C$(1), "=")+1]
2491     SELECT Function$
2492       CASE "MARKER", "MARK"
2493         Hp_ib$="MARK"&Value$&";"
2494       OUTPUT @NWA;Hp_ib$
2495       ENTER @NWA_data1;OUTPUTACTI;"
2496       OUTPUT @NWA; "OUTPACTI;"
2497       ENTER @NWA_data1;Marker(VAl(Value$), 1)
2498       CASE "STAR"
2499         OUTPUT @NWA;"STAR; OUTPUTACTI;"
2500         ENTER @NWA_data1;Start_freq
2501       CASE "STOP"
2502         OUTPUT @NWA;"STOP; OUTPUTACTI;"
2503         ENTER @NWA_data1;Stop_freq
2504       CASE "N", "N"
2505         GOSUB Get_sparm
2506       CASE ELSE
2507         PRINT Functions;" NOT DEFINED IN READ_NWA"
2508     END SELECT
2509   NEXT I
2510   ASSIGN @NWA TO *
2511   ASSIGN @NWA_data1 TO *
2512   ASSIGN @NWA_data2 TO *
2513   SUBEND
2514   *****
2515   Get_sparm: This subroutine gets the S-parameter data that is on the NWA CRT
2516   OUTPUT @NWA;"PARA?"
2517   Sparms=Sparms(2, 4)
2518   OUTPUT @NWA;"SWEM?"
2519   Swems=Swems(2, LEN(Swems)-1)
2520   SELECT Sparms
2521     CASE "S11"
2522       N=1
2523     CASE "S21"
2524       N=2
2525     CASE "S12"
2526       N=3
2527     CASE "S22"
2528       N=4
2529     CASE ELSE
2530       BEEP
2531     DISP "Sparms is invalid in Read_nwa"
2532     PAUSE
2533   END SELECT
2534   OUTPUT @NWA;"FOR?"
2535   *****

```

```

2547 ENTER @Nwa_data2:Preamble;Bytes
2548 Dcount=Bytes/16
2549 ALLOCATE S1(1:Dcount,2),F(1:Dcount)
2550 ENTER @Nwa_data2:S1(*)
2551 SELECT SMemS
2552 CASE "RAMP", "STEP"
2553 CALL Read_nwa("START=", "STOP=")
2554 FOR J=1 TO Dcount
2555 F(J)=Start_freq*(J-1)*(Stop_freq-Start_freq)/(Dcount-1)
2556 NEXT J
2557 CASE "FREQUENCY LIST"
2558 OUTPUT @Nwa:"FORM3;OUTPFREL;"
2559 ENTER @Nwa:Preamble,Bytes,F(*)
2560 CASE ELSE
2561 BEEP
2562 DISP "SmemS has not been properly selected in Read_nwa"
2563 PAUSE
2564 END SELECT
2565 MAT S(1:Dcount,N,1)= S1(1:Dcount,1)
2566 MAT S(1:Dcount,N,2)= S1(1:Dcount,2)
2567 MAT S(1:Dcount,5,1)= F
2568 DEALLOCATE F(*) S1(*)
2569 IF NOT BIT(Sval13,N-1) THEN Sval1d=Sval1d*2*(N-1)
2570 RETURN
2571 SUBEND
2572 *****
2573 *****
2574 *****
2575 SUB Zoom on peak
2576 !This subprogram assumes there is a resonance peak on the network analyzer
2577 !display. The routine will zoom in on the center frequency by reducing the
2578 !span until the resonance 'deltav' dB points are shown. Averaging depends
2579 !on signal level. Step frequency mode is used only for the last sweep.
2580 !The S(I,K) array is dimensioned (801,5,2) (1:Dcount,Ns,REAL/IMAG)
2581 !NS=1:511 2:521 3:512 4:522 5:Frequency(for K=1)
2582 !Written by Eric J. Vanzura NIST 725.02
2583 OPTION BASE 1
2584 ASSIGN @Nwa TO 716
2585 COM /Nwa_data/ Start_freq,Stop_freq,Marker(*)
2586 COM /Sparm/ Frequency(801),COMPLEX Sparm(801),SparmS(3),INTEGER Num_p
=> oints
2587 COM /S_array/ INTEGER Dcount,Sval1d,REAL S(801,5,2)
2588 Zoom_on_peak:
2589 Lastscale=.5
2590 Delta=.3
2591 Sparmratio=1.5
2592 !
2593 Scale=20*Lastscale
2594 Prty=VAL(SYSTEMS("SYSTEM PRIORITY"))+1
2595 DISP "Zooming in on peak"
2596 ! Resonance data already exists on NWA CRT.
2597 Averages=4
2598 Nwa_span=5.E+6
2599 GOSUB Take_sweeps
2600 CALL Optimizedisplay
2601 GOSUB Find_deltaspan
2602 Scale=Lastscale*10
2603 GOSUB Set_new_nwaspan
2604 GOSUB Select_averages
2605 Set_nwa("REF_VAL"=@VAL$(Level+Scale),"SCALE"=@VAL$(Scale))
2606 GOSUB Take_sweeps
2607 CALL Optimizedisplay
2608 GOSUB Find_deltaspan
2609 GOSUB Print_peak_info
2610 Scale=Lastscale
2611

```

```

2612 GOSUB Set_new_nwaspan
2613 GOSUB Select_averages
2614 Set_nwa("REF_VAL"=@VAL$(Level+Scale),"SCALE"=@VAL$(Scale))
2615 GOSUB Take_sweeps
2616 CALL Optimizedisplay
2617 GOSUB Find_deltaspan
2618 GOSUB Print_peak_info
2619 SUBEXIT
2620 !
2621 !
2622 !
2623 Set_new_nwaspan:
2624 SELECT Scale
2625 CASE >Lastscale
2626 IF Span<5000 THEN
2627 Span=3.E+6
2628 Nwa_span=Span
2629 ELSE
2630 Nwa_span=Span*3
2631 END IF
2632 CASE =Lastscale
2633 Nwa_span=Spanratio*Span
2634 END SELECT
2635 RETURN
2636 !
2637 !
2638 !
2639 Find_deltaspan:
2640 Set_nwa("MARK=CONT","MARK=1","MARK=MAX")
2641 Read_nwa("MARK=1")
2642 Freq=Marker(1,1)
2643 Level=Marker(1,2)
2644 Set_nwa("DEL RE=E" "MARKER TARGET=@VAL$(Delta))
2645 Set_nwa("MARK=2","MARK=MAX","MARK=TARGET","MARKER=SEARCH LEFT")
2646 Set_nwa("MARK=3","MARK=MAX","MARKER=TARGET")
2647 Read_nwa("MARK=2","MARK=3")
2648 Set_nwa("DELTA OFF=")
2649 IF ABS(Marker(2,1))>1000 AND ABS(Marker(3,1))>1000 THEN
2650 IF ABS(Marker(2,1))<ABS(.9*Marker(3,1)) THEN
2651 Span=2*ABS(Marker(3,1))
2652 ELSE
2653 Span=ABS(Marker(2,1))*ABS(Marker(3,1))
2654 ELSE
2655 Span=2*ABS(Marker(2,1))
2656 END IF
2657 Span=999999
2658 BEEP
2659 PRINT "Span cannot be determined from marker information;"
2660 PRINT " in Find_deltaspan."
2661 PAUSE
2662 END IF
2663 RETURN
2664 !
2665 !
2666 !
2667 Print_peak_info:
2668 IF Freq<1.E+10 THEN
2669 q/1.E+9 "Level=",Level,"Span=",Span,"Freq(Span/2)=", (Freq+Span/2)/1.E+9
2670 ELSE
2671 PRINT USING "K X DD 9DE 2X K S2D 2D 2X K 8D K X DD 9DE" "Freq=", F
2672 req/1.E+9 "Level=",Level,"Span=",Span,"Freq(Span/2)=", (Freq+Span/2)/1.E+9
2673 END IF
2674 RETURN
2675 !

```



```

2806 COM /Background/ GraphType$(12), Margins$(2), (10), PaperSize$(1)
2807 COM /Background/ REAL Pen_speed, INTEGER Backand_pen, Auto_time
2808 COM /Background/ INTEGER Auto_file, REAL X_cross, Y_cross, X
2809 COM /Background/ Xgrid_ticks(4), INTEGER Xmajor, Xminor
2810 COM /Background/ Ygrid_ticks(4), INTEGER Ymajor, Yminor
2811 COM /Background/ REAL Xmin_graph, Xmax_graph, Ymin_graph, Ymax_graph
2812 COM /Axes_labels/ Print_x_labels$(3), Print_y_labels$(3)
2813 COM /Axes_labels/ Sig_digits(2), REAL X_csize, Y_csize
2814 COM /Windowspace/ REAL Xmin, Xmid, Xmax, Ymin, Ymid, Ymax, graph edges UDU
2815 COM /Windowspace/ REAL Xleft, Xright, Ybottom, Ytop, paper edges UDU
2816 COM /Log_scale/ REAL Xcycles, Xbegin, Ycycles, Ybegin
2817 COM /Plot_device/ Plot_angles(10), INTEGER Plot_addr
2818 COM /Hard_space/ REAL View_lift, View_rt, View_btm, View_top
2819 COM /Hard_space/ REAL Viewscale, Aspect_ratio
2820 COM /Frame_onoff/ INTEGER Frame_flag
2821 COM /Clear_space/ REAL Space_lift, Space_rt, Space_btm, Space_top
2822 I
2823 COM /S array/ INTEGER Dcount, Svalld, REAL S(801, 5, 2)
2824 COM /Ftype/ INTEGER Sselect, Start_index, Stop_index
2825 COM /Cdata/ INTEGER Mfct(2), Lista(5, 2), Ma(2), REAL A(5), Ua(5), INTEGER
=> Ndata(2)
2826 COM /Stats/ Alameda, Chsq, INTEGER Nca(2), REAL Alpha(5, 5), Covar(5, 5)
2827 COM /Files/ Disksdrives(20), filenames$(10)
2828 COM /Bugs/ INTEGER Bug1, Bug2, Bug3, Printer
2829 COM /Raw_data/ Start_freq, Stop_freq, Marker(*)
2830 COM /Peak_data/ Peak(*), Upeak(*), INTEGER Peak_index, Num_peaks
2831 COM /Resonance/ Kx, Ky, G, FO, Uky, Uq, J, U
2832 COM /Output/ INTEGER Print_addr, Plotter, Printer_on, Plotter_on
2833 ASSIGN @HWA TO 716
2834 I
2835 Init_vals:
2836 Min_freq=6.55E+9
2837 Max_freq=1.35E+10
2838 Start_freq=8.2E+9
2839 Stop_freq=1.24E+10
2840 Step_freq=8.0E+7
2841 Min_sig_to_noi=-55
2842 Num_peaks=0
2843 Peak_index=0
2844 Save_to_later=1
2845 Init_sweep_menu
2846 Prty=VAL(SYSKEYS("SYSTEM PRIORITY"))+1
2847 ON KEY 0 LABEL "OK/DONE", Prty GOSUB Done_return
2848 ON KEY 1 LABEL "Startfreq", Prty GOSUB Startfreq
2849 ON KEY 2 LABEL "Stopfreq", Prty GOSUB Stopfreq
2850 ON KEY 3 LABEL "Stepfreq", Prty GOSUB Stepfreq
2851 ON KEY 4 LABEL "Min_sig_to_noi", Prty GOSUB Min_sig_to_noi
2852 ON KEY 5 LABEL "Print", Prty GOSUB Print
2853 ON KEY 6 LABEL "Min_sig_to_noi", Prty GOSUB Min_sig_to_noi
2854 ON KEY 7 LABEL "Min_sig_to_noi", Prty GOSUB Min_sig_to_noi
2855 IF Printer ON THEN
2856 ON KEY 8 LABEL "No printer", Prty GOSUB Tog_printer
2857 ELSE
2858 ON KEY 8 LABEL "Printer ON", Prty GOSUB Tog_printer
2859 END IF
2860 IF Save_to_later THEN
2861 ON KEY 9 LABEL "Compute Sparms", Prty GOSUB Tog_savefolater
2862 ELSE
2863 ON KEY 9 LABEL "Save Sparms", Prty GOSUB Tog_savefolater
2864 END IF
2865 Done=0
2866 Reset_keys=0
2867 REPEAT
2868 IF Reset_keys THEN GOTO Init_sweep
2869 UNTIL Done
2870 OFF KEY
2871 RETURN
2872 I
2873 I
2874 I
2875 Tog_printer: OFF KEY 8
2876 IF Printer ON THEN
2877 Printer_on=0
2878 ELSE Printer_on=1
2879 END IF
2880 Reset_keys=1
2881 RETURN
2882 I
2883 I
2884 Done_return: Done=1
2885 RETURN
2886 I
2887 Startfreq: CALL Enter_real(Min_freq, Stop_freq, "Sweep start frequency", Startfre
=> q)
2888 Reset_keys=1
2889 RETURN
2890 I
2891 Stopfreq: CALL Enter_real(Startfreq, Max_freq, "Sweep stop frequency", Stopfreq)
2892 Reset_keys=1
2893 RETURN
2894 I
2895 Stepfreq: CALL Enter_real(1.E+4, 3.E+8, "Sweep step span (Hz)", Stepfreq)
2896 Reset_keys=1
2897 RETURN
2898 I
2899 Min_peak_val: CALL Enter_real(-120, 20, "Minimum resonance peak value (dB)", Mi
=> n_peak_val)
2900 Reset_keys=1
2901 RETURN
2902 I
2903 Min_sig_to_noi: CALL Enter_real(3, 50, "Minimum signal-to-noise ratio", Min_sig
=> _to_noi)
2904 Reset_keys=1
2905 RETURN
2906 I
2907 Tog_savefolater: Reset_keys=1
2908 IF Save_to_later THEN
2909 Save_to_later=0
2910 ELSE Save_to_later=1
2911 END IF
2912 RETURN
2913 I
2914 I
2915 Sleep_menu: PRINT TABXY(1, 1);
2916 IF Startfreq<1.E+10 THEN
2917 PRINT USING "K, D, 60E, K"; Startfreq, " Hz"
2918 ELSE
2919 PRINT USING "K, D, 60E, K"; Startfreq, " Hz"
2920 END IF
2921 IF Stopfreq>1.E+10 THEN
2922 PRINT USING "K, D, 60E, K"; Stopfreq, " Hz"
2923 ELSE
2924 PRINT USING "K, D, 60E, K"; Stopfreq, " Hz"
2925 END IF
2926 PRINT USING "K, D, 20E, K"; Stepfreq, " Hz"
2927 PRINT
2928 PRINT USING "K, M3D, K"; Min_peak_val, " dB"
2929 PRINT USING "K, M3D, K"; Min_sig_to_noi, " dB"
2930 PRINT
2931 IF Printer ON THEN
2932 PRINT TP; Printer IS ON Print_addr, " "
2933 ELSE

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2934 PRINT "Printer is OFF Print_addr=";Print_addr;" "
2935 END IF
2936 PRINT
2937 IF Save_to_later THEN
2938 PRINT "The S-parameters will be saved on disk. Resonance paramete
=> rs won't be computed."
ELSE
2939 PRINT "The resonance parameters will be computed."
2940
2941 END IF
2942 RETURN
2943
2944
2945
2946 Print_info:
2947 ALLOCATE Desc$(160)
2948 REPEAT
2949 INPUT "Enter a description of this sweep. When finished enter a
null string.",Desc$
2950 PRINT Desc$
2951 IF Printer_on THEN OUTPUT Print_addr;Desc$
2952 UNTIL Desc$=""
2953 DEALLOCATE Desc$
2954 RETURN
2955
2956
2957
2958 Add_peak_2_list:
2959 SELECT Sselect
2960 CASE 1,2
2961 Peak(Num_peaks,1)=A(3)
2962 Peak(Num_peaks,2)=SQRT(A(1)^2+A(2)^2)
2963 Peak(Num_peaks,3)=A(3)/A(4)
2964 CASE ELSE
2965 BEEP
2966 DISP "Sweep_nwa needs to have S11 measurements added"
2967 PAUSE
2968 END SELECT
2969 RETURN
2970
2971
2972
2973 Save_sweepfreqs:OFF KEY
2974 IF Save_to_later THEN RETURN
2975 Prty=VAL(SYSTEMS("SYSTEM PRIORITY"))+1
2976 ON KEY 0 LABEL "NO",Prty GOSUB Done_return
2977 ON KEY 2 LABEL "YES",Prty GOSUB Freqs_to_disk
2978 DISP "Do you want to save the resonant frequency (list to disk?"
2979 Done=0
2980 LOOP
2981 EXIT IF Done
2982 END LOOP
2983 RETURN
2984
2985
2986
2987 Freqs_to_disk:OFF KEY
2988 ALLOCATE Basket(Num_peaks,2)
2989 MAT Basket(1:Num_peaks,1)= Peak(1:Num_peaks,1)
2990 MAT Basket(1:Num_peaks,2)= Peak(1:Num_peaks,3)
2991 CALL Enter_id(168,"the resonant frequency (list)")
2992 IF Id$="" THEN GOTO Abort_savefreqs
2993 CALL Select_disk
2994 IF Diskdrive$="" THEN GOTO Abort_savefreqs
2995 CALL Enterfilename("ABORT", " resonant frequency (list)")
2996 IF filename$="" THEN GOTO Abort_savefreqs

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```

2997 CALL Save_file(Basket(*),Num_peaks,Id$)
2998 Abort_savefreqs:
2999 Done=1
3000 RETURN
3001
3002
3003
3004 End_sweep_nwa:
3005 CALL Set_nwa("PRES=")
3006 IF Save_to_later THEN
3007 IF Save_to_later THEN
3008 Peak(Num_peaks,1)=1.063E+10
3009 Peak(Num_peaks,3)=2.E+7
3010 Peak(Num_peaks,2)=25
3011 END IF
3012 Peak_index=1
3013 CALL Set_nwa("CENTER=";VAL$(Peak_index,1))
3014 CALL Set_nwa("SPAN=";VAL$(Peak_index,3))
3015 CALL Set_nwa("REF_VAL=";VAL$(Peak_index,2)*5),"SCALE=5")
3016 LOCAL GNWB
3017 IF Printer_on THEN OUTPUT Print_addr;CHR$(12);
3018 BEEP
3019 GOSUB Save_sweepfreqs
3020 DISP "I am finished with the experiment oh great master."
3021 SUBEND
3022
3023
3024
3025 SUB Pack_data(File(*),OPTIONAL REAL F1(*),INTEGER Data_cnt,Pen,Id$)
3026 Pack_data:
3027 i Original: 01 Jun 1987 G. Keepke
3028 i Revision: 07 Aug 1987
3029 i This routine will take up to 17 independent data files and pack
3030 i the information into File(*) using GRAPH_DATA master file format.
3031 i The Roster(*) information will be generated for use by Plot_all.
3032 Roster(1,1) = Curve number 1,2,3,....,17.
3033 Roster(1,2) = Start address in File(k,*); = x
3034 Roster(1,3) = Datacount for curve k
3035 Roster(1,4) = PEN number for this curve
3036
3037 i Symbols(1)=" or "y" => no symbol, connect pts
3038 i Symbols(1)="ny" => * symbol, connect pts
3039 i Symbols(1)="n" => * symbol, do not connect pts
3040
3041 OPTION BASE 1
3042 COM /Data_param/ INTEGER Datacount,Filesize,Curvecount,Roster(*)
3043 COM /Data_param/ REAL Sym_size,Symbols(*),Curve_id$(*)
3044 COM /Data_param/ REAL Xmin_data,Xmax_data
3045 COM /Data_param/ REAL Ymin_data,Ymax_data
3046 COM /Bugs/ INTEGER Bug1,Bug2,Bug3,Printer
3047 INTEGER I,J,C
3048 Filesize=SIZE(File,1)
3049 IF NPAS=1 THEN ! Clear data parameters and exit
3050 Datacount=0
3051 Curvecount=0
3052 MAT Roster=(0)
3053 MAT Symbols=( "")
3054 MAT Curve_id$=( "")
3055 SUBEXIT
3056 END IF
3057 IF Data_cnt<1 THEN
3058 DISP " NO data in this file (packing aborted...): ";Id$
3059 BEEP
3060 PAUSE
3061 SUBEXIT
3062 END IF

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3063 ! Add new file to the data being plotted.
3064 Timers=TIMERDATE
3065 DISP " Packing .. "; lds
3066 Pen=MIN(MAX(Pen,1),8)
3067 SELECT Curvecount
3068 CASE 0
3069 Curvecount=1
3070 I=Curvecount
3071 Roster(I,1)=1
3072 Roster(I,2)=1
3073 CASE <17
3074 Curvecount=Curvecount+1
3075 I=Curvecount
3076 Roster(I,1)=I
3077 Roster(I,2)=Roster(I-1,2)+Roster(I-1,3)
3078 CASE ELSE
3079 DISP " CURVE limit has been reached, new data discarded. ";
3080 DISP " (continue) "
3081 BEEP
3082 PAUSE
3083 DISP CHR$(12);
3084 SUBEXIT
3085 END SELECT
3086 Roster(I,3)=Data_cnt
3087 Roster(I,4)=Pen
3088 Symbols(I)=""
3089 Curve_05(I)=lds
3090
3091 IF Roster(I,2)+Roster(I,3)-1>Filesize THEN
3092 DISP " DATA FILE overflow, new data discarded. ";
3093 DISP " (continue) "
3094 BEEP
3095 PAUSE
3096 DISP CHR$(12)
3097 Curvecount=Curvecount-1
3098 DISP CHR$(12)
3099 SUBEXIT
3100 END IF
3101 ! Copy data into File(*)
3102 C=1
3103 FOR J=Roster(I,2)+Roster(I,3)-1
3104 File(J,1)=F1(C,1)
3105 File(J,2)=F1(C,2)
3106 C=C+1
3107 NEXT J
3108 LOOP
3109 EXIT IF TIMERDATE-Timer>1
3110 END LOOP
3111 DISP CHR$(12)
3112 SUBEXIT
3113
3114 *****
3115 SUB Init_graphics(Label1$,Label2$,Label3$)
3116 ! Original: 01 Jun 1987 G. Koepke
3117 ! Revision: 07 Aug 1987
3118 ! Roster(i,1) = Curve number 1,2,3,...,17
3119 ! Roster(i,2) = Start address in File(x,*); = x
3120 ! Roster(i,3) = Datecount for curve i
3121 ! Roster(i,4) = PEN number for this curve
3122
3123 ! *****
3124 !
3125 !
3126 !
3127 !
3128 !
3129 ! Symbols(i)="m" or "y" => no symbol, connect pts
3130 ! Symbols(i)="n" => * symbol, connect pts
3131 ! Symbols(i)="n" => * symbol, do not connect pts
3132 ! Lbl_addr: x, y, pen, size, LDIR, LORG
3133
3134 COM /Labels/ Labels$,INTEGER Lbl_count,REAL Lbl_addr(*)
3135 COM /Data param/ INTEGER Datacount,Filesize,Curvecount,Roster(*)
3136 COM /Data param/ REAL Sym_size,Symbols(*),Curve_ids(*)
3137 COM /Data param/ REAL Ymin_data,Ymax_data
3138 COM /Data param/ REAL Ymin_data,Ymax_data
3139 COM /Background/ GraphType$,Margins$,Papersize$
3140 COM /Background/ REAL Pen_speed,INTEGER Background_pen,Auto_time
3141 COM /Background/ INTEGER Auto_file,REAL X_cross,Y_cross,X
3142 COM /Background/ Xgrid_ticks$,INTEGER Xmajor,Xminor
3143 COM /Background/ Ygrid_ticks$,INTEGER Ymajor,Yminor
3144 COM /Background/ REAL Xmin_graph,Xmax_graph,Ymin_graph,Ymax_graph
3145 COM /Axes_labels/ Print_labels$,Print_Ylabels$
3146 COM /Axes_labels/ Sig_digits$,REAL Xlcsize,Ylcsize
3147
3148 COM /Windowspace/ REAL Xmin,Xmid,Xmax,Ymin,Ymid,Ymax,graph edges,UDU
3149 COM /Windowspace/ REAL Xleft,Xright,Ybottom,Ytop, paper edges,UCUS
3150 COM /Plot_device/ Plot_lang$,INTEGER Plot_addr
3151 COM /Hard_space/ REAL Xview_lft,Xview_rt,Yview_btm,Yview_top
3152 COM /Hard_space/ REAL ViewScale$,Aspect_ratio
3153 COM /Frame_onoff/ INTEGER Frame_flag
3154 COM /Clear_space/ REAL Space_lft,Space_rt,Space_btm,Space_top
3155 COM /Bugs/ INTEGER Bug1,Bug2,Bug3,Printer
3156
3157 !
3158 !
3159 !
3160 !
3161 Ginit_clear:GINIT ! Clear all graphics
3162 GLEAR
3163 OUTPUT 2 USING "#,K,," K" ! Clear the screen
3164 Initial_values: ! DEFINE ALL INITIAL VALUES
3165 Frame_flag=1 ! Completely frame the data area.
3166 Print_xlabels$="YES" ! AND print labels.
3167 Print_ylabels$="YES"
3168 Sig_digits$="fff" ! Select free format for X & Y labels.
3169 Xlcsize=.035 ! Label size factor for axes labels.
3170 Ylcsize=.035 ! graph data default size = .032
3171 Plot_lang$="INTERNAL" OR "HPGL"
3172 Plot_addr=3 ! OR 705 (HP1B ADDR OF PLOTTER)
3173 Margins$(1)="HORIZONTAL" OR "VERTICAL" ! for plotter only
3174 Margins$(2)="FULL" ! OR "BOUND LEFT"
3175 ! OR "BOUND TOP"
3176 ! OR "SQUARE"
3177 ! OR "USER"
3178 Papersize$="4" ! 8.5x11
3179 ! OR "3", 11x17
3180 Background_pen=1 ! 1 to 8
3181 Auto_time=1 ! Automatically label time/date on graphs.
3182 Auto_file=0 ! Automatically label last file name.
3183 Pen_speed=20.0 ! .38 cm/s TO 36.1 cm/s
3184
3185 !
3186 ! SET VIEWPORT PARAMETERS
3187 Xview_lft=34.4
3188 Xview_rt=100*RATIO
3189 Yview_btm=24
3190 Yview_top=100
3191 ! Clear space around labels in terms of frac of scale
3192 Space_lft=.21
3193 Space_rt=.02

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3195 Space_btm=.11
3196 Space_top=.10
3197
3198 ! Background parameters.....
3199
3200 GraphType$="LINEAR"
3201 Xgrid_ticks$="TICK"
3202 Ygrid_ticks$="TICK"
3203 Xmajor=6
3204 Ymajor=1
3205 Xminor=3
3206 Yminor=1
3207 Xmax_graph=1.8E+10
3208 Xmin_graph=0
3209 Ymax_graph=0
3210 Ymin_graph=1
3211 X_cross_y=Ymin_graph
3212 Y_cross_x=Xmin_graph
3213
3214 Xmin=Xmin_graph
3215 Xmax=Xmax_graph
3216 Xmid=Xmin*(Xmax-Xmin)/2
3217 Ymin=Ymin_graph
3218 Ymax=Ymax_graph
3219 Ymid=Ymin*(Ymax-Ymin)/2
3220
3221 ! Label parameters.....
3222
3223 Lbl_count=0
3224 MAT Labels$="( "
3225 MAT Lbl_addr$=(0)
3226
3227 ! Insert desired labels here using UDUS defined above.
3228
3229 Lbl_count=3
3230 Labels(1)=Label1$
3231 Lbl_addr(1,1)=9.E+9 ! X location UDUS
3232 Lbl_addr(1,2)=.05 ! Y location UDUS
3233
3234 Lbl_addr(1,3)=5 ! PEN
3235 Lbl_addr(1,4)=-.035 ! SIZE factor
3236 Lbl_addr(1,5)=0 ! LDIR
3237 Lbl_addr(1,6)=6 ! LONG
3238
3239 Labels(2)="Relative amplitude (dB)"
3240 Lbl_addr(2,1)=2.E+9 ! X location UDUS
3241 Lbl_addr(2,2)=-5 ! Y location UDUS
3242
3243 Lbl_addr(2,3)=5 ! PEN
3244 Lbl_addr(2,4)=-.035 ! SIZE factor
3245 Lbl_addr(2,5)=0 ! LDIR
3246 Lbl_addr(2,6)=4 ! LONG
3247
3248 Labels(3)="Probe Calibration Data"
3249 Lbl_addr(3,1)=9.E+9 ! X location UDUS
3250 Lbl_addr(3,2)=1.02 ! Y location UDUS
3251
3252 Lbl_addr(3,3)=5 ! PEN
3253 Lbl_addr(3,4)=.040 ! SIZE factor
3254 Lbl_addr(3,5)=0 ! LDIR
3255 Lbl_addr(3,6)=4 ! LONG
3256
3257 SUBEXIT
3258
3259 ! *****
3260
3261 SUB Plot_all(File$)
3262 Plot_all: Original: 13 Nov 1984
3263 Revision: 06 Aug 1987
3264
3265 OPTION BASE 1
3266 DEG
3267
3268 COM /SYS/ Sys_ids[10]
3269 COM /Labels/ Labels$(*),INTEGER Lbl_count,REAL Lbl_addr(*)
3270 COM /Data_param/ INTEGER Datacount,Filesize,CurveCount,Roster(*)
3271 COM /Data_param/ REAL Sym_size,Symbol$(*),Curve_ids$(*)
3272 COM /Data_param/ REAL Xmin_data,Xmax_data
3273 COM /Data_param/ REAL Ymin_data,Ymax_data
3274 COM /Background/ GraphType$,Margins$(*),PaperSize$
3275 COM /Background/ REAL Pen_speed,INTEGER Backand_pen,Auto_time
3276 COM /Background/ INTEGER Auto_file,REAL X_cross,Y_cross,X
3277 COM /Background/ Xgrid_ticks$,INTEGER Xmajor,Xminor
3278 COM /Background/ Ygrid_ticks$,INTEGER Ymajor,Yminor
3279 COM /Background/ REAL Xmin_graph,Xmax_graph,Ymin_graph,Ymax_graph
3280
3281 COM /Axes_labels/ Print_xlabel$,Print_ylabel$
3282 COM /Axes_labels/ Sig_digits$,REAL Xlsize,Ylsize
3283
3284 COM /WindowSpace/ REAL Xmin,Xmid,Xmax,Ymin,Ymid,Ymax,graph edges UDUS
3285 COM /WindowSpace/ REAL Xleft,Xright,Ybottom,Ytop, paper edges UDUS
3286 COM /Plot device/ Plot_lang$,INTEGER Plot_addr
3287 COM /Hard_space/ REAL Xview_lft,Xview_rt,Yview_btm,Yview_top
3288 COM /Hard_space/ REAL ViewScale,Aspect_ratio
3289 COM /Frame on/off/ INTEGER Frame_flag
3290 COM /Clear_space/ REAL Space_lft,Space_rt,Space_btm,Space_top
3291 COM /Bugs/ INTEGER Bug1,Bug2,Bug3,Printer
3292 COM /Files/ Diskdrive$,Filename$
3293 INTEGER I,J,Local_prty,Suspended,Outofbounds
3294
3295 Aspect_ratio=(Xview_rt-Xview_lft)/((Yview_top-Yview_btm)*2)
3296 ViewScale=MIN(Xview_rt-Xview_lft,(Yview_top-Yview_btm))
3297 Aspect_ratio=MIN(MAX(Aspect_ratio,.35),.65)
3298
3299 Local_prty=VAL(SYSSTEM$( "SYSTEM PRIORITY" ))+1
3300 IF Lbl_count=0 THEN
3301 ALLOCATE Lbl_ratio(Lbl_count,2)
3302 GOSUB Update_lbl_addr ! includes find_mid_pt
3303 DEALLOCATE Lbl_ratio(*)
3304
3305 ELSE
3306 GOSUB Find_mid_pt
3307 END IF
3308 CALL Background(Suspended)
3309 IF Suspended THEN GOTO Nomoregraph
3310 IF Plot_addr<>3 THEN
3311 ON KEY 0 LABEL "ABORT GRAPH",Local_prty+1 G Nomoregraph
3312 END IF
3313 CLIP OFF
3314 GOSUB Label_labels
3315 GOSUB Plot_data
3316 PENUP
3317 PEN 0
3318 CLIP OFF
3319 SUBEXIT
3320
3321 ! *****
3322 Update_lbl_addr: Make sure that the label addresses are current.
3323 FOR I=1 TO Lbl_count
3324 Lbl_ratio(I,1)=(Lbl_addr(I,1)-Xmin)/(Xmax-Xmin)
3325 Lbl_ratio(I,2)=(Lbl_addr(I,2)-Ymin)/(Ymax-Ymin)
3326

```

```

3327 NEXT I
3328 GOSUB Find_mid_pt
3329 FOR I=1 TO Lbl_count
3330 Lbl_addr(1,1)=Lbl_ratio(1,1)*(Xmax-Xmin)+Xmin
3331 Lbl_addr(1,2)=Lbl_ratio(1,2)*(Ymax-Ymin)+Ymin
3332 NEXT I
3333 RETURN
3334
3335
3336
3337 find_mid_pt:
3338 ! Interpret the graph type and the scaling done.
3339 CASE "LINEAR"
3340 Xmins=Xmin_graph
3341 Xmaxs=Xmax_graph
3342 Xmid=(Xmax-Xmin)/2+Xmin
3343 Ymins=Ymin_graph
3344 Ymaxs=Ymax_graph
3345 Ymid=(Ymax-Ymin)/2+Ymin
3346 CASE "SEMILOG X"
3347 Xmins=0
3348 Xmaxs=100
3349 Xmid=50
3350 Ymins=Ymin_graph
3351 Ymaxs=Ymax_graph
3352 Ymid=(Ymax-Ymin)/2+Ymin
3353 CASE "SEMILOG Y"
3354 Xmins=Xmin_graph
3355 Xmaxs=Xmax_graph
3356 Xmid=(Xmax-Xmin)/2+Xmin
3357 Ymins=0
3358 Ymaxs=100
3359 Ymid=50
3360 CASE "LOG LOG"
3361 Xmins=0
3362 Xmaxs=100
3363 Xmid=50
3364 Ymins=0
3365 Ymaxs=100
3366 Ymid=50
3367 CASE "POLAR"
3368 !
3369 BEEP
3370 DISP "POLAR PARAMETERS ARE NOT YET IMPLEMENTED!!!"
3371 PAUSE
3372 END SELECT
3373 RETURN
3374
3375
3376
3377 Label_labels:
3378 ! ALL LABELS ARE APPLIED
3379 PEN Backgnd_pen
3380 CSIZE DROUND(.025*Viewscale,3),Aspect_ratio
3381 IF Auto_time THEN
3382 LOG 1
3383 MOVE Xleft,Ybottom
3384 LABEL " @DATES(TIMEDATE)& ", "&TIMES(TIMEDATE)
3385 PENUP
3386 END IF
3387 IF Auto_file THEN
3388 LOG 7
3389 MOVE Xright,Ybottom
3390 LABEL " @FILENAME>0 THEN
3391 LABEL "file: "&filename&" "
3392 END IF
3393 PENUP

```

```

3393
3394
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3409
3410
3411 Plot_date:
3412 ! PLOT ALL CURVES
3413 Disptime=0.
3414 CLIP Xmin,Xmax,Ymin,Ymax
3415 LOG 5
3416 LDIR 0
3417 CSIZE DROUND(.025*Viewscale,3),Aspect_ratio
3418 FOR I=1 TO Curvecount
3419 Outofbounds=0
3420 PENUP
3421 PEN Roster(I,4)
3422 Xpos=File(Roster(I,2),1)
3423 Ypos=File(Roster(I,2),2)
3424 GOSUB Adjust_xy_pos
3425 IF Xpos>Xmax THEN
3426 GOSUB Skipcurvej
3427 GOTO Skipcurvej
3428 END IF
3429 Xpos=MAX(MIN(Xpos,Xmax),Xmin)
3430 Ypos=MAX(MIN(Ypos,Ymax),Ymin)
3431 MOVE Xpos,Ypos
3432 SELECT LEN(Symbol$(Roster(I,1)))
3433 CASE 1
3434 Mark$=Symbol$(Roster(I,1))
3435 Connect$=Symbol$(Roster(I,1))
3436 CASE 2
3437 Mark$=Symbol$(Roster(I,1))
3438 Connect$=Symbol$(Roster(I,1))
3439 CASE ELSE
3440 Mark$=""
3441 Connect$=""
3442 END SELECT
3443 IF Mark$="" AND Connect$="" THEN
3444 ! No symbol
3445 FOR J=Roster(I,2) TO Roster(I,2)+Roster(I,3)-1
3446 Xpos=File(J,1)
3447 Ypos=File(J,2)
3448 GOSUB Adjust_xy_pos
3449 IF Xpos>Xmax THEN
3450 GOSUB Skipcurvej
3451 GOTO Skipcurvej
3452 END IF
3453 IF NOT Outofbounds THEN
3454 IF Xpos<Xmin OR Ypos<Ymin OR Ypos>Ymax THEN
3455 Outofbounds=1
3456 GOSUB Skipcurvej
3457 END IF
3458 END IF

```

```

3459 Xpos=MAX(MIN(Xpos,Xmax),Xmin)
3460 Ypos=MAX(MIN(Ypos,Ymax),Ymin)
3461 DRAW Xpos,Ypos
3462 NEXT J
3463
3464 !put in symbol
3465 SELECT Connects
3466 CASE "H"
3467   FOR J=ROSTER(1,2) TO ROSTER(1,2)+ROSTER(1,3)-1
3468     Xpos=File(J,1)
3469     Ypos=File(J,2)
3470     GOSUB Adjust_xy_pos
3471     IF Xpos>Xmax THEN
3472       GOSUB Skipcurvej
3473     ELSE
3474       GOTO Skipcurvej
3475     END IF
3476     IF NOT Outofbounds THEN
3477       IF Xpos<Xmin OR Ypos<Ymin OR Ypos>Ymax THEN
3478         Outofbounds=1
3479         GOSUB Skipcurvej
3480       END IF
3481     ELSE
3482       Xpos=MAX(MIN(Xpos,Xmax),Xmin)
3483       Ypos=MAX(MIN(Ypos,Ymax),Ymin)
3484       DRAW Xpos,Ypos
3485       LABEL USING "#,K","Mark$
3486     PENUP
3487     MOVE Xpos,Ypos
3488     NEXT J
3489   CASE "M"
3490     FOR J=ROSTER(1,2) TO ROSTER(1,2)+ROSTER(1,3)-1
3491       Xpos=File(J,1)
3492       Ypos=File(J,2)
3493       GOSUB Adjust_xy_pos
3494       IF Xpos>Xmax THEN
3495         GOSUB Skipcurvej
3496       ELSE
3497         GOTO Skipcurvej
3498       END IF
3499       IF NOT Outofbounds THEN
3500         IF Xpos<Xmin OR Ypos<Ymin OR Ypos>Ymax THEN
3501           Outofbounds=1
3502           GOSUB Skipcurvej
3503         END IF
3504       ELSE
3505         Xpos=MAX(MIN(Xpos,Xmax),Xmin)
3506         Ypos=MAX(MIN(Ypos,Ymax),Ymin)
3507         LABEL USING "#,K","Mark$
3508       PENUP
3509       MOVE Xpos,Ypos
3510       NEXT J
3511     END SELECT
3512   END IF
3513   skipcurvej: I
3514   NEXT I
3515   LOOP
3516   EXIT IF TIME-Disptime>1.2
3517   END LOOP
3518   DISP CHR$(12)
3519   END IF
3520   RETURN
3521   I
3522   I
3523   I
3524   skipcurvej: I
3525
3526 DISP " DATA OUT OF RANGE ... NOT PLOTTED "
3527 IF Disptime<1.0E-10 THEN
3528   Disptime=TIMEDATE
3529 END IF
3530 RETURN
3531
3532
3533 Adjust_xy_pos: !CORRECT Xpos, Ypos FOR VARIOUS AXES
3534 SELECT GraphTypes$
3535 CASE "LINEAR"
3536   !NO CHANGE
3537 CASE "SEMILOG X"
3538   Xpos=FNLin_map_logx(Xpos)
3539   IF Xpos<0 THEN Xpos=0
3540 CASE "SEMILOG Y"
3541   Ypos=FNLin_map_logy(Ypos)
3542   IF Ypos<0 THEN Ypos=0
3543 CASE "LOG LOG"
3544   Xpos=FNLin_map_logx(Xpos)
3545   Ypos=FNLin_map_logy(Ypos)
3546   IF Xpos<0 THEN Xpos=0
3547   IF Ypos<0 THEN Ypos=0
3548 CASE "POLAR"
3549   !NO CHANGE yet.
3550 END SELECT
3551 RETURN
3552
3553
3554
3555 SUBEXIT
3556
3557
3558 *****
3559 SUB Background (INTEGER Suspended)
3560 ! Original: 13 Nov 1984
3561 ! Revision: 06 Aug 1987
3562 ! This SLB program is written to draw the background for the PLOT
3563 ! program. It draws LINEAR, SEMILOG, LOG LOG, and POLAR backgrounds.
3564 ! Parameters are determined in the main program.
3565 !
3566 ! Due to the complexity of LOG and POLAR coordinates, the TICKING is
3567 ! done with MOVES and DRAMS. Also avoided are device dependent codes.
3568 !
3569 ! OPTION BASE 1
3570 DEG
3571 COM /Sys/ Sys_ids
3572 COM /Background/ GraphTypes$,Margins$(*),Papersizes
3573 COM /Background/ REAL Pen_speed,INTEGER Backgnd_pen,Auto_time
3574 COM /Background/ INTEGER Auto_file,REAL X_cross,Y_cross,X
3575 COM /Background/ Xgrid_ticks,INTEGER Xmajor,Xminor
3576 COM /Background/ REAL Y_ticks,INTEGER Ymajor,Yminor
3577 COM /Background/ REAL Xmin_graph,Xmax_graph,Ymin_graph,Ymax_graph
3578 COM /WindowSpace/ REAL Xmin,Xmid,Xmax,Ymin,Ymid,YmaxI,graph_edges
3579 COM /Plot device/ Plot Length(10),INTEGER Plot_addr
3580 COM /Hard space/ REAL Xview_lft,Xview_rt,Yview_btn,Yview_top
3581 COM /Log scale/ REAL Xcycles,Xbegin,Ycycles,Ybegin
3582 COM /Axis labels/ Print_xlabel$,Print_ylabel$$
3583 COM /Axis labels/ size_dflist$,REAL Xlsize,Ylsize
3584 COM /Frame on/off/ INTEGER Frame_flag
3585 COM /Clear space/ REAL Space_lft,Space_rt,Space_btm,Space_top
3586
3587 COM /Bus1/ INTEGER Bus1,Bus2,Bus3,Printer

```

```

3657 END IF
3658 Ytics=(Ymaxg-Yming)/(Ymajor*Yminor)
3659 GOSUB Window_space
3660 IF Print_xLabels$="YES" THEN GOSUB Label_x_log
3661 IF Print_yLabels$="YES" THEN GOSUB Label_y_linear
3662
3663 CASE "SEMILOG Y"
3664   I Set up user units. Outside edges
3665   Yming=0. I range. X axis is unchanged.
3666   Ytic_size=Log_tic_size*(Xmaxg-Xming)*Tic_ratio
3667   Find_log_cycles(Ymin_graph,Ymax_graph,Ycycles,Ybegin,Yend)
3668   IF X_cross_Y<2*Ybegin THEN
3669     Xcrossy=0.
3670   ELSE
3671     Xcrossy=FMLin_map_logY(X_cross_Y)
3672   END IF
3673   Xtics=(Ymaxg-Yming)/(Xmajor*Xminor)
3674   GOSUB Window_space
3675   GOSUB Draw_linear
3676   IF Print_xLabels$="YES" THEN GOSUB Label_x_linear
3677   GOSUB Draw_y_log
3678   IF Print_yLabels$="YES" THEN GOSUB Label_y_log
3679   I
3680 CASE "LOG LOG"
3681   I Set up user units. Outside edges
3682   Xming=0. I All log operations are mapped into this
3683   Ymaxg=100. I range. Both axes.
3684   Yming=0.
3685   Xmaxg=100.
3686   Ytic_size=Log_tic_size*(Ymaxg-Yming)
3687   Find_log_cycles(Xmin_graph,Xmax_graph,Xcycles,Xbegin,Xend)
3688   IF X_cross_Y<2*Ybegin THEN
3689     Xcrossy=0.
3690   ELSE
3691     Xcrossy=FMLin_map_logY(X_cross_Y)
3692   END IF
3693   Y_cross_x<2*Xbegin THEN
3694     Ycrossx=0.
3695   ELSE
3696     Ycrossx=FMLin_map_logX(Y_cross_x)
3697   END IF
3698   GOSUB Window_space
3699   GOSUB Draw_x_log
3700   IF Print_xLabels$="YES" THEN GOSUB Label_x_log
3701   GOSUB Draw_y_log
3702   IF Print_yLabels$="YES" THEN GOSUB Label_y_log
3703   I
3704 CASE "POLAR"
3705   I 177777777777
3706   END SELECT
3707   GOTO Back_done
3708   I End plotting:Suspended=1
3709   Back_done: I
3710   IF Plot_addr<-3 THEN OFF KEY
3711   PEN 0
3712   SUBEXIT
3713   I
3714   I
3715   I
3716   I
3717   I
3718   I
3719   I
3720   I
3721   I
3722   I
3723   I
3724   I
3725   I
3726   I
3727   I
3728   I
3729   I
3730   I
3731   I
3732   I
3733   I
3734   I
3735   I
3736   I
3737   I
3738   I
3739   I
3740   I
3741   I
3742   I
3743   I
3744   I
3745   I
3746   I
3747   I
3748   I
3749   I
3750   I
3751   I
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3753   I
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3787   I
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3789   I
3790   I
3791   I
3792   I
3793   I
3794   I
3795   I
3796   I
3797   I
3798   I
3799   I
3800   I
3801   I
3802   I
3803   I
3804   I
3805   I
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3807   I
3808   I
3809   I
3810   I
3811   I
3812   I
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3830   I
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3834   I
3835   I
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3837   I
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3864   I
3865   I
3866   I
3867   I
3868   I
3869   I
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3873   I
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3875   I
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3877   I
3878   I
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3880   I
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3964   I
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3966   I
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3983   I
3984   I
3985   I
3986   I
3987   I
3988   I
3989   I
3990   I
3991   I
3992   I
3993   I
3994   I
3995   I
3996   I
3997   I
3998   I
3999   I
4000   I

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3591 DIM Digits$(1)
3592 INTEGER Tic_size,Toggle,Log_label_limit,Log_label_step
3593 INTEGER Xminor,Yminor
3594 REAL Xcrossy,Ycrossx,Xmaxg,Ymaxg,Yming,Ymaxg
3595
3596 GRAPHICS ON
3597 IF Print_xLabels$="" THEN Print_xLabels$="YES"
3598 IF Print_yLabels$="" THEN Print_yLabels$="YES"
3599 IF Plot_addr<-3 THEN
3600   Suspended=0
3601   CALL Hp7475a_setup(suspended)
3602   IF Suspended THEN GOTO End_plotting
3603   ON KEY 0 LABEL "ABORT GRAPH",Local_prty+1 GOTO End_plotting
3604 ELSE
3605   VIEWPORT Xview_Lft,Xview_Rt,Yview_Btm,Yview_Top
3606   IF Frame_flg THEN FRAME
3607   END IF
3608   LOG 0
3609   LOG 0
3610   Log_tic_size=.04 IPer cent of total range.
3611   Tic_top=(Yview_top-Yview_btm)/(Xview_Rt-Xview_Lft)
3612   Xming=Xmin_graph
3613   Xmaxg=Xmax_graph
3614   Yming=Ymin_graph
3615   Ymaxg=Ymax_graph
3616   Xcrossx=X_cross_X
3617   Ycrossy=Y_cross_Y
3618   IF Ymajor<1 THEN Ymajor=1
3619   IF Yminor<1 THEN Yminor=1
3620   IF Xmajor<1 THEN Xmajor=1
3621   IF Xminor<1 THEN Xminor=1
3622   IF Xgrid_tick$="GRID" AND Xminor<1 THEN
3623     Xminor=1
3624   ELSE
3625     Yminor=Xminor
3626   END IF
3627   IF Ygrid_tick$="GRID" AND Yminor<1 THEN
3628     Yminor=1
3629   ELSE
3630     Xminor=Yminor
3631   END IF
3632   Log_label_limit=4
3633
3634 PEN Background_pen
3635 SELECT GraphTypes
3636 CASE "LINEAR"
3637   I Set up user units. Outside edges
3638   Xtics=(Xmaxg-Xming)/(Xmajor*Xminor)
3639   Ytics=(Ymaxg-Yming)/(Ymajor*Yminor)
3640   GOSUB Window_space
3641   GOSUB Draw_x_linear
3642   IF Print_xLabels$="YES" THEN GOSUB Label_x_linear
3643   GOSUB Draw_y_linear
3644   IF Print_yLabels$="YES" THEN GOSUB Label_y_linear
3645   I
3646 CASE "SEMILOG X"
3647   I Set up user units. Outside edges
3648   Xming=0. I All log operations are mapped into
3649   Ymaxg=100. I this range. Y axis range is unchanged.
3650   Ytic_size=Log_tic_size*(Ymaxg-Yming)
3651   Find_log_cycles(Xmin_graph,Xmax_graph,Xcycles,Xbegin,Xend)
3652   IF X_cross_x<2*Xbegin THEN
3653     Xcrossx=0.
3654   ELSE
3655     Xcrossx=FMLin_map_logX(Y_cross_x)
3656   END SELECT

```

Allow room outside of curve box for labels.
 I Also scale to curve box and turn on CLIP to these

```

3785 MOVE Xming,Xcrossy,(Xtic_size/2)
3786 RPL0T 0,Xtic_size,-1
3787 FOR C=1 TO Full_cycles
3788   FOR D=2 TO 10
3789     XLoc=Full_in_map_logx(0)*Xbegin*(10^(C-1))
3790     MOVE XLoc,Xcrossy-(Xtic_size/2)
3791     RPL0T 0,Xtic_size,-1
3792   NEXT D
3793 NEXT C
3794 FOR D=1 TO Div_beyond
3795   XLoc=Full_in_map_logx((D+1)*Xbegin*(10^(Full_cycles)))
3796   MOVE XLoc,Xcrossy-(Xtic_size/2)
3797   RPL0T 0,Xtic_size,-1
3798 NEXT D
3799 MOVE Xmaxg,Xcrossy-(Xtic_size/2)
3800 RPL0T 0,Xtic_size,-1
3801 MOVE Xmaxg,Xcrossy
3802 DRAW Xming,Xcrossy
3803 PENUP
3804 !
3805 !CHECK FOR OPPOSITE AXIS.
3806 !
3807 IF Xcrossy=Yming OR Xcrossy=Ymaxg THEN
3808   !Repeat for the opposite axis
3809   IF Xcrossy=Yming THEN
3810     Second_axis=Ymaxg
3811   ELSE
3812     Second_axis=Yming
3813   END IF
3814 MOVE Xming,Second_axis-(Xtic_size/2)
3815 RPL0T 0,Xtic_size,-1
3816 FOR C=1 TO Full_cycles
3817   FOR D=2 TO 10
3818     XLoc=Full_in_map_logx(D*Xbegin*(10^(C-1)))
3819     MOVE XLoc,Second_axis-(Xtic_size/2)
3820     RPL0T 0,Xtic_size,-1
3821   NEXT D
3822 NEXT C
3823 FOR D=1 TO Div_beyond
3824   XLoc=Full_in_map_logx((D+1)*Xbegin*(10^(Full_cycles)))
3825   MOVE XLoc,Second_axis-(Xtic_size/2)
3826   RPL0T 0,Xtic_size,-1
3827 NEXT D
3828 MOVE Xmaxg,Second_axis-(Xtic_size/2)
3829 RPL0T 0,Xtic_size,-1
3830 MOVE Xmaxg,Second_axis
3831 DRAW Xming,Second_axis
3832 PENUP
3833 END IF
3834 CASE "GRID"
3835 MOVE Xming,Yming
3836 DRAW Xming,Ymaxg
3837 Toggles+1 !Toggles between Ymaxg and Yming.
3838   !+1=Ymaxg, -1=Yming
3839 FOR C=1 TO Full_cycles
3840   FOR D=2 TO 10
3841     XLoc=Full_in_map_logx(D*Xbegin*(10^(C-1)))
3842     SELECT Toggle
3843     CASE +1
3844       MOVE XLoc,Ymaxg
3845       DRAW XLoc,Yming
3846     Toggle=Toggle
3847     CASE -1
3848       MOVE XLoc,Yming
3849       DRAW XLoc,Ymaxg
3850     Toggles=Toggle

```

```

3723 !Limits.
3724 IF Margins$(2)="SQUARE" THEN
3725   Edge_lft=Xming-.21*(Xmaxg-Xming)
3726   Edge_rt=Xmaxg+.21*(Xmaxg-Xming)
3727   Edge_btm=Yming-.15*(Ymaxg-Yming)
3728   Edge_top=Ymaxg+.15*(Ymaxg-Yming)
3729 ELSE
3730   Edge_lft=Xming-Space_lft*(Xmaxg-Xming)
3731   Edge_rt=Xmaxg+Space_rt*(Xmaxg-Xming)
3732   Edge_btm=Yming-Space_btm*(Ymaxg-Yming)
3733   Edge_top=Ymaxg+Space_top*(Ymaxg-Yming)
3734 IF
3735 WINDOW,Edge_lft,Edge_rt,Edge_btm,Edge_top
3736 CLIP Xming,Xmaxg,Yming,Ymaxg
3737 IF Frame_flag THEN FRAME
3738 RETURN
3739 !
3740 !
3741 !
3742 Draw_x_linear:
3743   !Draw X (horizontal) axes.
3744   SELECT Xgrid_ticks
3745   CASE "TICK"
3746     AXES Xtics,0,Xming-Xtics,Xminor,Xcrossy,Xminor,Tyminor,Tic_size
3747     !Select opposite axis if necessary
3748     SELECT Xcrossy
3749     CASE <=Yming+1.0E-10*(Ymaxg-Yming)
3750       AXES Xtics,0,Xming-Xtics,Xminor,Ymaxg,Xminor,Tyminor,Tic_size
3751     => ze
3752     CASE >=Ymaxg-1.0E-10*(Ymaxg-Yming)
3753       AXES Xtics,0,Xming-Xtics,Xminor,Yming,Xminor,Tyminor,Tic_size
3754     => ze
3755 END SELECT
3756 CASE "GRID"
3757   GRID Xtics,0,Xcross,Xminor,Tyminor,Tic_size
3758 END SELECT
3759 RETURN
3760 !
3761 !
3762 Draw_y_linear:
3763   !Draw Y (vertical) axes.
3764   SELECT Ygrid_ticks
3765   CASE "TICK"
3766     AXES 0,Ytics,Ycross,Yminor,Ytics*Tyminor,Xminor,Tyminor,Tic_size
3767     !Select the opposite axis if necessary.
3768     SELECT Ycross
3769     CASE <=Xming+1.0E-10*(Xmaxg-Xming)
3770       AXES 0,Ytics,Xmaxg,Yming-Ytics*Yminor,Xminor,Tyminor,Tic_size
3771     => ze
3772     CASE >=Xmaxg-1.0E-10*(Xmaxg-Xming)
3773       AXES 0,Ytics,Xming,Yming-Ytics*Yminor,Xminor,Tyminor,Tic_size
3774     => ze
3775 END SELECT
3776 CASE "GRID"
3777   GRID 0,Ytics,Ycross,Xminor,Tyminor,Tic_size
3778 END SELECT
3779 RETURN
3780 !
3781 !
3782 Draw_x_log:
3783   !Draw log axis according to x parameters.
3784   Full_cycles=INT(Xcycles)
3785   Div_beyond=10-FRACT(Xcycles)
3786   SELECT Xgrid_ticks
3787   CASE "TICK"

```

```

3917 MOVE Second_axis-(Ytic_size/2),Yloc
3918 RPL0T Ytic_size,0,-1
3919 NEXT D
3920 MOVE Second_axis-(Ytic_size/2),Ymaxg
3921 RPL0T Ytic_size,0,-1
3922 MOVE Second_axis,Ymaxg
3923 DRAW Second_axis,Ymaxg
3924 PENUP
3925 END IF
3926 CASE "GRID"
3927 MOVE Xming,Yming
3928 DRAW Xmaxg,Yming
3929 Toggle=+1 !Toggles between Ymaxg and Yming.
3930 FOR C=1 TO Full_cycles
3931   I=I+Ymaxg, -I=Yming
3932   FOR D=2 TO 10
3933     Yloc=FMLin_map_logy(0*Ybegin*(10^(C-1)))
3934     SELECT Toggle
3935     CASE +1
3936       MOVE Xmaxg,Yloc
3937       DRAW Xming,Yloc
3938       Toggle=-Toggle
3939     CASE -1
3940       MOVE Xming,Yloc
3941       DRAW Xmaxg,Yloc
3942       Toggle=-Toggle
3943     END SELECT
3944   NEXT D
3945 NEXT C
3946 FOR D=1 TO Div_beyond+1
3947   Yloc=FMLin_map_logy((D+1)*Ybegin*(10^(Full_cycles)))
3948   SELECT Toggle
3949   CASE +1
3950     MOVE Xmaxg,Yloc
3951     DRAW Xming,Yloc
3952     Toggle=-Toggle
3953   CASE -1
3954     MOVE Xming,Yloc
3955     DRAW Xmaxg,Yloc
3956     Toggle=-Toggle
3957   END SELECT
3958 NEXT D
3959 END SELECT
3960 RETURN
3961 !
3962 !
3963 !
3964 Fixed_sig_digit: ! Assure equal significant digits to right of .
3965 CASE 1.0 TO 9.99
3966   SELECT ABS(Numeric_label)
3967   CASE 1.0 TO 9.99
3968     SELECT Numeric_label
3969     CASE PROUND(Numeric_label,-INT(VAL(Digits)))
3970     SELECT Digits
3971     CASE "1"
3972       LABEL USING "MD.0";Numeric_label
3973     CASE "2"
3974       LABEL USING "MD.DD";Numeric_label
3975     CASE "3"
3976       LABEL USING "MD.0D";Numeric_label
3977     CASE "4"
3978       LABEL USING "MD.0D";Numeric_label
3979     END SELECT
3980   CASE ELSE
3981     LABEL USING "MD.DDESZ";Numeric_label
3982   END SELECT

```

```

3851 END SELECT
3852 NEXT D
3853 NEXT C
3854 FOR D=1 TO Div_beyond+1
3855   Yloc=FMLin_map_logy((D+1)*Xbegin*(10^(Full_cycles)))
3856   SELECT Toggle
3857   CASE +1
3858     MOVE Xloc,Ymaxg
3859     DRAW Xloc,Ymaxg
3860     Toggle=-Toggle
3861   CASE -1
3862     MOVE Xloc,Yming
3863     DRAW Xloc,Ymaxg
3864     Toggle=-Toggle
3865   END SELECT
3866 NEXT D
3867 END SELECT
3868 RETURN
3869 !
3870 !
3871 !
3872 Draw_Y_log:
3873 Full_cycles=INT(Ycycles)
3874 Div_beyond=10*FRACT(Ycycles)
3875 SELECT Ygrid_ticks
3876 CASE "TICK"
3877   MOVE Ycross-(Ytic_size/2),Yming
3878   RPL0T Ytic_size,0,-1
3879   FOR C=1 TO Full_cycles
3880     FOR D=2 TO 10
3881       Yloc=FMLin_map_logy(0*Ybegin*(10^(C-1)))
3882       MOVE Ycross-(Ytic_size/2),Yloc
3883       RPL0T Ytic_size,0,-1
3884     NEXT D
3885   NEXT C
3886 FOR D=1 TO Div_beyond
3887   Yloc=FMLin_map_logy((D+1)*Ybegin*(10^(Full_cycles)))
3888   MOVE Ycross-(Ytic_size/2),Yloc
3889   RPL0T Ytic_size,0,-1
3890 NEXT D
3891 MOVE Ycross-(Ytic_size/2),Ymaxg
3892 RPL0T Xtic_size,0,-1
3893 MOVE Ycross,Ymaxg
3894 DRAW Ycross,Ymaxg
3895 PENUP
3896 !
3897 !
3898 !
3899 !
3900 !
3901 !
3902 !
3903 !
3904 !
3905 !
3906 !
3907 !
3908 !
3909 !
3910 !
3911 !
3912 !
3913 !
3914 !
3915 !
3916 !

```



```

4181 LABEL USING "MD.DDESZ";Numeric_label
4182 END SELECT
4183
4184
4185 CASE ELSE
4186 IF ABS(Numeric_label)>1.0E-99 THEN
4187 IF PROUND(Numeric_label,-9)=Numeric_label THEN
4188 LABEL USING "MD.DDESZ";Numeric_label
4189 ELSE
4190 LABEL USING "MD.DDESZZ";Numeric_label
4191 END IF
4192 ELSE
4193 LABEL USING "D";Numeric_label
4194 END IF
4195 END SELECT
4196 RETURN
4197
4198
4199 Label_x_linear: !Put numeric labels at every MAJOR tick mark.
4200 Digress=Sig_digits$(1); ! Select X significant digits
4201 CSIZE=ROUND(VIEWSIZE*XICSIZE,3),ASPECT_RATIO
4202 CLIP OFF
4203 LDIR 0
4204 Tick=X:ics*Txminor !Divisions for labeling
4205 LOG 6
4206 FOR Numeric_label=Xming TO Xmaxg-.5*Tick STEP Tick
4207 MOVE Numeric_label,Yming-.005*(Ymaxg-Yming)
4208 GOSUB Label_format
4209 NEXT Numeric_label
4210 MOVE Xmaxg,Yming-.005*(Ymaxg-Yming)
4211 LOG 9
4212 Numeric_label=Xmaxg
4213 GOSUB Label_format
4214 PENUP
4215 CLIP ON
4216 RETURN
4217
4218
4219 Label_y_linear: !Put numeric labels at every MAJOR tick mark.
4220 Digress=Sig_digits$(2); ! Select Y significant digits
4221 CSIZE=ROUND(VIEWSIZE*YICSIZE,3),ASPECT_RATIO
4222 CLIP OFF
4223 LDIR 0
4224 Tick=Y:ics*Tyminor !Divisions for labeling
4225 LOG 8
4226 FOR Numeric_label=Yming+Tick TO Ymaxg-.5*Tick STEP Tick
4227 MOVE Ymtg-.005*(Ymaxg-Yming),Numeric_label
4228 GOSUB Label_format
4229 NEXT Numeric_label
4230 MOVE Xmaxg-.005*(Xmaxg-Xming),Ymaxg
4231 LOG 9
4232 Numeric_label=Ymaxg
4233 GOSUB Label_format
4234 PENUP
4235 CLIP ON
4236 RETURN
4237
4238
4239 Label_x_log: !Put numeric labels at every log cycle and end.
4240 !if more than 4 log cycles then thin the labels.
4241
4242
4243
4244
4245
4246

```

```

4115 IF INT(Numeric_label)=ROUND(Numeric_label,5) THEN
4116 LABEL USING "MD";Numeric_label
4117 ELSE
4118 SELECT Numeric_label
4119 CASE PROUND(Numeric_label,-1)
4120 LABEL USING "M2D.D";Numeric_label
4121 CASE PROUND(Numeric_label,-2)
4122 LABEL USING "M2D.DD";Numeric_label
4123 CASE PROUND(Numeric_label,-3)
4124 LABEL USING "M2D.DDD";Numeric_label
4125 CASE ELSE
4126 LABEL USING "MD.DDESZ";Numeric_label
4127 END SELECT
4128
4129 END IF
4130
4131 CASE 99.99 TO 999.99
4132 IF INT(Numeric_label)=ROUND(Numeric_label,6) THEN
4133 LABEL USING "MDD";Numeric_label
4134 ELSE
4135 SELECT Numeric_label
4136 CASE PROUND(Numeric_label,-1)
4137 LABEL USING "M3D.D";Numeric_label
4138 CASE PROUND(Numeric_label,-2)
4139 LABEL USING "M3D.DD";Numeric_label
4140 CASE ELSE
4141 LABEL USING "MD.DDESZ";Numeric_label
4142 END SELECT
4143
4144 END IF
4145
4146 CASE 999.99 TO 9999.99
4147 IF INT(Numeric_label)=ROUND(Numeric_label,7) THEN
4148 LABEL USING "M4D";Numeric_label
4149 ELSE
4150 SELECT Numeric_label
4151 CASE PROUND(Numeric_label,-1)
4152 LABEL USING "M4D.D";Numeric_label
4153 CASE PROUND(Numeric_label,-2)
4154 LABEL USING "M4D.DD";Numeric_label
4155 CASE ELSE
4156 LABEL USING "MD.DDESZ";Numeric_label
4157 END SELECT
4158
4159 END IF
4160
4161 CASE >9999.99
4162 SELECT ABS(Numeric_label)
4163 CASE <1.0E+10
4164 LABEL USING "MD.DDESZ";Numeric_label
4165 CASE <1.0E+100
4166 LABEL USING "MD.DDESZZ";Numeric_label
4167 CASE ELSE
4168 LABEL USING "MD.DDESZZZ";Numeric_label
4169 END SELECT
4170
4171 !+++++All values less than 1.0 ++++++
4172 CASE 0001 TO 1.0
4173 SELECT Numeric_label
4174 CASE PROUND(Numeric_label,-1)
4175 LABEL USING "M2.D";Numeric_label
4176 CASE PROUND(Numeric_label,-2)
4177 LABEL USING "M2D";Numeric_label
4178 CASE PROUND(Numeric_label,-3)
4179 LABEL USING "M2D.D";Numeric_label
4180 CASE PROUND(Numeric_label,-4)
4181 LABEL USING "M2D.DD";Numeric_label
4182 CASE ELSE

```

```

4313 Log Label_steps=INT(DROUND((Ycycles/(Log_Label_limit-1)+.5),1))
4314 IF Log_Label_steps<1 THEN Log_Label_step=1
4315 FOR C=0 TO INT(Ycycles) STEP Log_Label_step
4316 IF C=0 THEN
4317   Numeric_Label=begin*(10^C)
4318   YLoc=FMLin_map_logy(Numeric_Label)
4319   MOVE XLoc,YLoc
4320 IF C=INT(Ycycles) AND INT(Ycycles)=Ycycles THEN LOG 9
4321 GOSUB Label_format
4322 END IF
4323 NEXT C
4324
4325 IF INT(Ycycles)<Ycycles AND INT(Ycycles)<Log_Label_limit-1 THEN
4326   LOG 9
4327   MOVE XLoc,Ymaxg
4328   Numeric_Label=Yend
4329   GOSUB Label_format
4330 END IF
4331 CLIP ON
4332 RETURN
4333
4334 SUBEND
4335
4336 SUB Find_log_cycles(Low,High,Cycles,New_Low,New_High)
4337   ! Determine the number of LOG cycles that will cover
4338   ! the range of MIN (>0) to MAX for either axis. There will
4339   ! be at least ONE cycle and may be stopped at the first
4340   ! 1/10 cycle above the scale MAX.
4341   ! The variable 'cycles' has as the integer part the number
4342   ! of FULL cycles and as the fractional part the number of
4343   ! 1/10 cycle divisions beyond the last FULL cycle.
4344   ! 'New_Low' gives a lower value with one digit.
4345   ! 'New_High' gives the new upper limit value.
4346
4347   ! Original: 13 Nov 1984 - G. Kospke
4348   ! Revision: 06 Aug 1987
4349
4350   COM /Bugs/ INTEGER Bug1,Bug2,Bug3,Printer
4351   INTEGER Exponent,M,N
4352   Exponent=0
4353   IF Low<1.0E-50 THEN
4354     DISP " Range error for LOG plot, begin point too small.";
4355     BEEP
4356     PAUSE
4357     DISP CHR$(12)
4358     LOW=1.0E-50
4359   END IF
4360   N=0
4361   SELECT LOW
4362   CASE <1.0
4363     REPEAT
4364       N=N+1
4365       Test=Low*10.0^(N)
4366       IF Test>=1.0 THEN
4367         Exponent=N
4368         Test=DROUND(Test,12)
4369         New_Low=INT(Test)*10.0^(Exponent)
4370       END IF
4371     UNTIL Exponent<>0
4372   CASE >=1.0
4373     REPEAT
4374       N=N+1
4375     UNTIL Exponent<>0
4376   REPEAT
4377     N=N+1
4378

```

```

4247 !where Log_Label_limit=4
4248 Digits=Sig_digits(1,1) ! Select X significant digits
4249 CSIZE DROUND(Viewscale*(Xsize/5),Aspect_ratio)
4250 CLIP OFF
4251 LOG 6
4252 YLoc=YIning-.005*(Ymaxg-Yming)
4253 !Label left corner
4254 Numeric_Label=Xbegin
4255 XLoc=FMLin_map_logx(Numeric_Label)
4256 MOVE XLoc,YLoc
4257 GOSUB Label_format
4258
4259 IF INT(Ycycles)<Log_Label_limit THEN
4260   FOR C=1 TO INT(Ycycles)
4261     Numeric_Label=Xbegin*(10^C)
4262     XLoc=FMLin_map_logx(Numeric_Label)
4263     MOVE XLoc,YLoc
4264     IF C=INT(Ycycles) AND INT(Ycycles)=Ycycles THEN LOG 9
4265     GOSUB Label_format
4266   NEXT C
4267
4268 ELSE Log_Label_steps=INT(DROUND((Xcycles/(Log_Label_limit-1)+.5),1))
4269 IF Log_Label_steps<1 THEN Log_Label_step=1
4270 FOR C=0 TO INT(Xcycles) STEP Log_Label_step
4271 IF C=0 THEN
4272   Numeric_Label=Xbegin*(10^C)
4273   XLoc=FMLin_map_logx(Numeric_Label)
4274   MOVE XLoc,YLoc
4275 IF C=INT(Xcycles) AND INT(Xcycles)=Xcycles THEN LOG 9
4276 GOSUB Label_format
4277 END IF
4278 NEXT C
4279
4280 IF INT(Xcycles)<Xcycles AND INT(Xcycles)<Log_Label_limit-1 THEN
4281   IF FRACT(Xcycles)>.31 THEN !Only label end point if >3/10 cycle
4282     LOG 9
4283     MOVE Xmaxg,YLoc
4284     Numeric_Label=Xend
4285     GOSUB Label_format
4286   END IF
4287   CLIP ON
4288   RETURN
4289
4290 !
4291 !
4292 !
4293 Label_Y_log: IPut numeric labels at every log cycle, and end.
4294 Digits=Sig_digits(2,1) ! Select Y significant digits
4295 CSIZE DROUND(Viewscale*(Ysize/5),Aspect_ratio)
4296 CLIP OFF
4297 LOG 7
4298 XLoc=XIning-.005*(Xmaxg-Xming)
4299 !Label lower corner
4300 MOVE XLoc,YIning
4301 Numeric_Label=Ybegin
4302 GOSUB Label_format
4303 LOG 8
4304 IF INT(Ycycles)<Log_Label_limit THEN
4305   FOR C=1 TO INT(Ycycles)
4306     Numeric_Label=Ybegin*(10^C)
4307     XLoc=FMLin_map_logy(Numeric_Label)
4308     MOVE XLoc,YLoc
4309     IF C=INT(Ycycles) AND INT(Ycycles)=Ycycles THEN LOG 9
4310     GOSUB Label_format
4311   NEXT C
4312 ELSE

```

```

4445  IR = range of linear scale = 100.0
4446  IC = number of cycles mm,m, nn = whole cycles, m = divisions
4447  I = beyond last whole cycle.
4448
4449  R=100.0
4450  B=Xbegin
4451  C=Xcycles
4452  IF V<8 THEN V=B/10
4453  P=C/(INT(C)+LGT(10*FRACT(C+1)))*LGT(V/B)
4454  RETURN P
4455  FNEND
4456
4457  *****
4458  DEF FNLin_map_Logx(V)
4459  COM /Log_scale/ REAL Xcycles,Xbegin,Ycycles,Ybegin
4460
4461  IP = position, in linear units on that axis, of the LGT(V)
4462  IV = value to be mapped on the LOG scale.
4463  IB = begin LOG value at linear 0.0, ie. 0.1, 0.003, 10, etc.
4464  IR = range of linear scale = 100.0,
4465  IC = number of cycles mm,m, nn = whole cycles, m = divisions
4466  I = beyond last whole cycle.
4467
4468  R=100.0
4469  B=Ybegin
4470  C=Ycycles
4471  IF V<8 THEN V=B/10
4472  P=C/(INT(C)+LGT(10*FRACT(C+1)))*LGT(V/B)
4473  RETURN P
4474  FNEND
4475
4476  *****
4477  SUB Hp7475a_setup(INTEGER Suspended)
4478  Hp7475a_setup: Original: 13 Nov 1984
4479  Revision: 06 Aug 1987
4480
4481  Optimize use of the Hp 7475a plotter to draw
4482  various axes types. This is the first step to
4483  draw the background for the graph.
4484
4485  OPTION BASE 1
4486  DEG
4487  COM /background/ GraphType$,Margins$,PaperSize$
4488  COM /background/ REAL Pen_speed,INTEGER Backgrnd_pen,Auto_time
4489  COM /background/ INTEGER Auto_file,REAL X_cross,Y_cross,X_
4490  COM /background/ Grid_ticks,INTEGER Xmajor,Xminor
4491  COM /background/ Grid_ticks,INTEGER Ymajor,Yminor
4492  COM /background/ REAL Xmin_graph,Xmax_graph,Ymin_graph,Ymax_graph
4493  COM /hard space/ REAL Xview_lift,Xview_rt,Yview_btm,Yview_top
4494  COM /Plot_device/ Plot_lang$,INTEGER Plot_addr
4495
4496  INTEGER Hlty,Hlty,Hurx,Hury
4497  DIM Outputs(80)
4498  ASSIGN Aplotter TO Plot_addr
4499
4500  GOSUB Initial_7475a
4501  IF NOT Suspended THEN GOSUB Select_margin
4502  SUBEXIT
4503
4504  *****
4505  Select margin: Iset margins and VIEWPORT.
4506  IIF GraphType$="POLAR" THEN I must define polar parameters
4507
4508  *****
4509
4510

```

```

4379  Test=Low*10.0^( -N)
4380  IF Test<10. THEN
4381  Exponent=N
4382  Test=DROND(Test,12)
4383  New_Low=INT(Test)*10^(Exponent)
4384  IF Bug3 THEN
4385  PRINT USING "3(MD.16DE)";New_Low,Test,INT(Test)
4386  END IF
4387  END IF
4388  UNTIL Exponent<=0
4389  END SELECT
4390
4391  IF Low>=1.0 AND Low<10.0 THEN
4392  New_Low=INT(Low)
4393  END IF
4394  New_Low=DROND(New_Low,1)
4395
4396  IF High<=10*New_Low THEN
4397  Cycles=1.0_Then
4398  ELSE
4399  N=1
4400  LOOP
4401  Test=DROND(New_Low*10.0^(N),3)
4402  IF Bug3 THEN
4403  PRINT "New_Low:";New_Low;" Full Cycles=";N;
4404  PRINT " Test=";Test;" Max=";High
4405  END IF
4406  N=N+1
4407  EXIT IF High<=Test
4408  END LOOP
4409  IF Test=High THEN Find_range
4410  IF Bug3 THEN PRINT
4411
4412  N=0
4413  I=1/10 Cycles
4414  IF Find the number of divisions above full cycles - 1
4415  necessary to cover the range.
4416  M=M+1
4417  Test=DROND((M+1)*New_Low*10.0^(N-1),3)
4418  IF Bug3 THEN
4419  PRINT "Cycles=";N-1;"1/10s=";M;" Test=";Test;" Max=";High
4420  END IF
4421  END IF
4422  EXIT IF High<=Test
4423  END LOOP
4424  Find_range:
4425  SELECT M
4426  CASE 0, >=9
4427  Cycles=N
4428  CASE <9
4429  Cycles=(N-1)+(M/10)
4430  END SELECT
4431  New_High=New_Low*(10.0*FRACT(Cycles)+1)*10.0^(INT(Cycles))
4432  SUBEXIT
4433  *****
4434  DEF FNLin_map_Logx(V)
4435  COM /Log_scale/ REAL Xcycles,Xbegin,Ycycles,Ybegin
4436
4437  IP = position, in linear units on that axis, of the LGT(V)
4438  IV = value to be mapped on the LOG scale.
4439  IB = begin LOG value at linear 0.0, ie. 0.1, 0.003, 10, etc.
4440
4441
4442
4443
4444

```

```

4577 Initial 7475a:
4578 !Initialize 7475A plotter, and expand P1 & P2 to include maximum
4579 !plotting area.
4580 !
4581 ON TIMEOUT 7,12 GOTO Plotter_dead !Plotter hangs bus!
4582 OUTPUT apPlotter;"IN" !Initialize 7475A plotter.
4583 OUTPUT apPlotter;"PS";papersizes !Select paper size
4584 OUTPUT apPlotter;"VS";gvals(pen_speed) !Set the pen to Pen_speed
4585 !and slow acceleration to .2
4586 IF Margins$(1)="VERTICAL" THEN
4587 OUTPUT apPlotter;"R090;IP;1W"
4588 END IF
4589 IF Margins$(2)="USER DEFN" THEN !Send the P1,P2 coordinates.
4590 OUTPUT apPlotter;"OP"
4591 ENTER apPlotter;HlX,HlY,HlY,HlY,HlY,HlY
4592 OUTPUT apPlotter;"R0"
4593 OUTPUT apPlotter;"SP1"
4594 Test_p1_p2:
4595 Output$="PA";gvals(HlX)&"",gvals(HlY)
4596 OUTPUT apPlotter;Output$
4597 BEEP
4598 DISP " MOVE PEN TO DESIRED LOCATION FOR LOWER LEFT CORNER.";
4599 PAUSE "...press CONTINUE."
4600 ENTER apPlotter;"OK"
4601 OUTPUT apPlotter;HlX,HlY,HlY,HlY,HlY,HlY
4602 !
4603 Output$="PA";gvals(HlX)&"",gvals(HlY)
4604 OUTPUT apPlotter;Output$
4605 BEEP
4606 DISP " MOVE PEN TO DESIRED LOCATION FOR UPPER RIGHT CORNER.";
4607 PAUSE "...press CONTINUE."
4608 ENTER apPlotter;"OK"
4609 OUTPUT apPlotter;"NOA"
4610 !
4611 Output$="PA";gvals(HlX)&"",gvals(HlY)
4612 OUTPUT apPlotter;Output$
4613 BEEP 400;3
4614 GOTO test_p1_p2
4615 END IF
4616 !
4617 DISP "Generating GRAPH on HP7475A plotter."
4618 ELSE
4619 OUTPUT apPlotter;"OH" !Send the HARD CLIP limits.
4620 ENTER apPlotter;HlX,HlY,HlY,HlY,HlY,HlY
4621 !
4622 SELECT Margins$(2) ! SELECT QUADRANTS
4623 CASE "LOW LEFT"
4624 Hurx=Hurx-(Hurx-HlX)/2
4625 Hury=Hury-(Hury-HlY)/2
4626 CASE "UP RIGHT"
4627 HlX=HlX+(Hurx-HlX)/2
4628 HlY=HlY+(Hury-HlY)/2
4629 CASE "UP LEFT"
4630 HlX=HlX+(Hurx-HlX)/2
4631 HlY=HlY+(Hury-HlY)/2
4632 CASE "LOW RIGHT"
4633 Hurx=Hurx-(Hurx-HlX)/2
4634 HlX=HlX+(Hurx-HlX)/2
4635 Hury=Hury-(Hury-HlY)/2
4636 CASE ELSE
4637 ! NO CHANGE
4638 END SELECT
4639 Output$="IP";gvals(HlX)&"",gvals(HlY)&"",gvals(HlY)&"",gvals(HlY)&"",
4640 Output$
4641 !
4642

```

```

4511 Margins$(2)="SQUARE"
4512 !through out program.
4513 IF Margins$(2)="USER DEFN" THEN
4514 GOSUB User
4515 GOTO Skipmargins
4516 END IF
4517 SELECT Margins$(1)
4518 CASE "HORIZONTAL"
4519 GOSUB Horiz_setup
4520 SELECT Margins$(2)
4521 CASE "BOUND TOP"
4522 GOSUB H_boundtop
4523 CASE "BOUND LEFT"
4524 GOSUB H_boundleft
4525 CASE "FULL"
4526 GOSUB H_full
4527 CASE "SQUARE"
4528 GOSUB H_square
4529 CASE ELSE
4530 GOSUB User
4531 END SELECT
4532 CASE "VERTICAL"
4533 GOSUB Vertical_setup
4534 SELECT Margins$(2)
4535 CASE "BOUND TOP"
4536 GOSUB V_boundtop
4537 CASE "BOUND LEFT"
4538 GOSUB V_boundleft
4539 CASE "FULL"
4540 GOSUB V_full
4541 CASE "SQUARE"
4542 GOSUB V_square
4543 CASE ELSE
4544 GOSUB User
4545 END SELECT
4546 Skipmargins:
4547 !
4548 !Scale the Viewscale and Aspect_ratio for the lettering.
4549 Xview_lft=Left_mar
4550 Xview_rt=Right_mar
4551 Yview_btm=Bottom_mar
4552 Yview_top=Top_mar
4553 SELECT Margins$(1)
4554 CASE "HORIZONTAL"
4555 Aspect_ratio=(Xview_rt-Xview_lft)/((Yview_top-Yview_btm)*2)
4556 CASE "VERTICAL"
4557 Aspect_ratio=(Yview_top-Yview_btm)/((Xview_rt-Xview_lft)*2)
4558 END SELECT
4559 Viewscale=MIN(MAX(Aspect_ratio,.3),1.5)
4560 VIEWPORT Left_mar,Right_mar,Bottom_mar,Top_mar
4561 IF Bug2 THEN
4562 PRINTER IS Printer
4563 PRINT "P1, P2 coordinates are (X,Y): ",HlX,HlY,HlY,HlY,HlY,HlY
4564 PRINT "viewport (Xl,Yl,Yr,Yt)",
4565 PRINT Left_mar,Right_mar,Bottom_mar,Top_mar
4566 PRINT USING "5/"
4567 PRINTER IS CRT
4568 BEEP
4569 PAUSE
4570 END IF
4571 RETURN
4572
4573
4574
4575
4576

```

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4643 !
4644 !RESET all graphics operations, read in HARD CLIP limits and
4645 !scale the result to the same dimensions as the screen, which
4646 !is (X longer) 0,100*RATIO,0,100 or (Y longer) 0,100,0,100/RATIO.
4647 !Both VIEWPORT and WINDOW operations are performed.
4648
4649 GINIT
4650 PLOTTER IS Plot_addr,Plot_langs
4651 PENUP
4652 OFF TIMEOUT 7
4653 RETURN
4654 !+++++*****
4655 Plotter dead: !The bus is hung.
4656 DISP "the plotter is NOT responding!"
4657 BEEP
4658 WAIT 1.8
4659 DISP CHR$(12)
4660 Suspended=1
4661 OFF TIMEOUT 7
4662 RETURN
4663
4664 !
4665 !
4666 !User: !This special margins routine will handle the USER DEFN hard
4667 !clip limits.
4668 !The shorter of Hurx-Hltx and Hury-Hlly is scaled to 100 units
4669 !by the PLOTTER IS command. The longer is scaled to 100*RATIO units.
4670
4671 Left_mar=0
4672 Bottom_mar=0
4673 IF Hurx-Hltx>Hury-Hlly THEN
4674 Top_mar=100
4675 ELSE
4676 Right_mar=100
4677 Top_mar=100/RATIO
4678
4679 END IF
4680 RETURN
4681
4682 !
4683 !
4684 !Horizon setup: !Determine the GDUS/cm for the appropriate paper.
4685 IF Papersize$="3" THEN
4686 XLength=100*RATIO
4687 Xgdu_cm=100*RATIO/41.45
4688 YLength=100
4689 Ygdu_cm=100/25.85
4690
4691 ELSE
4692 XLength=100*RATIO
4693 Xgdu_cm=100*RATIO/25.82
4694 YLength=100
4695 Ygdu_cm=100/15.85
4696 END IF
4697 RETURN
4698
4699 !
4700 !
4701 H_boundtop: !set margins for top binding.
4702 IF Papersize$="3" THEN
4703 Left_mar=Xgdu_cm*1.17
4704 Right_mar=XLength-Xgdu_cm*2.25
4705 Bottom_mar=Ygdu_cm*1.65
4706 Top_mar=YLength-Ygdu_cm*2.75
4707
4708

```

```

4709 ELSE
4710 ! 8.5x11 inch paper.
4711 Left_mar=Xgdu_cm*1.6
4712 Right_mar=XLength-Xgdu_cm*1.4
4713 Bottom_mar=Ygdu_cm*2.25
4714 Top_mar=YLength-Ygdu_cm*2.39
4715
4716 END IF
4717 RETURN
4718
4719 !
4720 !
4721 H_boundleft: !set margins for left side binding.
4722 IF Papersize$="3" THEN
4723 Left_mar=Xgdu_cm*2.4
4724 Right_mar=XLength-Xgdu_cm*2.25
4725 Bottom_mar=Ygdu_cm*1.65
4726 Top_mar=YLength-Ygdu_cm*1.45
4727
4728 ELSE
4729 ! 8.5x11 inch paper.
4730 Left_mar=Xgdu_cm*2.89
4731 Right_mar=XLength-Xgdu_cm*1.4
4732 Bottom_mar=Ygdu_cm*2.25
4733 Top_mar=YLength-Ygdu_cm*1.15
4734
4735 END IF
4736 RETURN
4737
4738 !
4739 !
4740 H_full: !set margins to fullest dimensions.
4741 IF Papersize$="3" THEN
4742 Left_mar=Xgdu_cm*0.
4743 Right_mar=XLength-Xgdu_cm*0.
4744 Bottom_mar=Ygdu_cm*15
4745 Top_mar=YLength-Ygdu_cm*0.
4746
4747 ELSE
4748 ! 8.5x11 inch paper.
4749 Left_mar=Xgdu_cm*25
4750 Right_mar=XLength-Xgdu_cm*0.
4751 Bottom_mar=Ygdu_cm*1.12
4752 Top_mar=YLength-Ygdu_cm*0.
4753
4754 END IF
4755 RETURN
4756
4757 !
4758 !
4759 H_square: !set margins for square centered on paper.
4760 IF Papersize$="3" THEN
4761 Left_mar=Xgdu_cm*7.35
4762 Right_mar=XLength-Xgdu_cm*8.45
4763 Bottom_mar=Ygdu_cm*15
4764 Top_mar=YLength-Ygdu_cm*0.
4765
4766 ELSE
4767 ! 8.5x11 inch paper.
4768 Left_mar=Xgdu_cm*3.65
4769 Right_mar=XLength-Xgdu_cm*3.4
4770 Bottom_mar=Ygdu_cm*1.12
4771 Top_mar=YLength-Ygdu_cm*0.
4772
4773
4774

```

```

4841      ! 11x17 inch paper.
4842      Left_mar=Xgdu_cm*.15
4843      Right_mar=Xlength-Xgdu_cm*0.
4844      Bottom_mar=Ygdu_cm*.15
4845      Top_mar=Ylength-Ygdu_cm*0.
4846
4847      ELSE
4848          ! 8.5x11 inch paper.
4849          Left_mar=Xgdu_cm*0.
4850          Right_mar=Xlength-Xgdu_cm*.07
4851          Bottom_mar=Ygdu_cm*.20
4852          Top_mar=Ylength-Ygdu_cm*0.
4853
4854      END IF
4855      RETURN
4856
4857      !
4858      !
4859      !Set margins for square centered on paper.
4860      IF Papersize="3" THEN
4861          ! 11x17 inch paper.
4862          Left_mar=Xgdu_cm*.15
4863          Right_mar=Xlength-Xgdu_cm*0.
4864          Bottom_mar=Ygdu_cm*.15
4865          Top_mar=Ylength-Ygdu_cm*.15
4866
4867      ELSE
4868          ! 8.5x11 inch paper.
4869          Left_mar=Xgdu_cm*0.
4870          Right_mar=Xlength-Xgdu_cm*.07
4871          Bottom_mar=Ygdu_cm*.20
4872          Top_mar=Ylength-Ygdu_cm*.15
4873
4874      END IF
4875      RETURN
4876      SUBEND
4877
4878      *****
4879      SUB Wipe_clean
4880      ! Original: 13 Nov 1984
4881      ! Revision: 06 Aug 1987
4882      ! Clear the CRT and home the cursor.
4883      PRINTER IS CRT
4884      CONTROL KBD 1:0
4885      CONTROL CRT 4:0
4886      OUTPUT 2 USING "#,4A;" K T"
4887      GINIT
4888      GCLEAR
4889      ALPHA ON
4890      GRAPHICS OFF
4891      SUBEND
4892
4893      *****
4894      SUB Autoscale(Files(*),Ax$)
4895      !
4896      !
4897      !
4898      !
4899      !
4900      ! Original: 01 Oct 1985, Eric Vanzura
4901      ! Revision: 13 Aug 1987, 11:00 by G. Koepke
4902
4903      ! This subprogram finds the max and min of X and/or Y axis data
4904      ! then selects the graph type with it's starting and stopping points
4905      ! File(*) contains packed data of one or more curves.
4906

```

```

4775      END IF
4776      RETURN
4777
4778      !
4779      !
4780      !
4781      !Setup the vertical dimensions uniformly.
4782      IF Papersize="3" THEN
4783          ! 11x17 inch paper.
4784          Xlength=100
4785          Ygdu_cm=100/25.85
4786          Ylength=100/RATIO
4787          Ygdu_cm=100/(RATIO*.41.45)
4788      ELSE
4789          ! 8.5x11 inch.
4790          Xlength=100
4791          Ygdu_cm=100/19.85
4792          Ylength=100/RATIO
4793          Ygdu_cm=100/(RATIO*.25.82)
4794      END IF
4795      RETURN
4796
4797      !
4798      !
4799      !Set margins for top binding.
4800      IF Papersize="3" THEN
4801          ! 11x17 inch paper.
4802          Left_mar=Xgdu_cm*.60
4803          Right_mar=Xlength-Xgdu_cm*.45
4804          Bottom_mar=Ygdu_cm*.30
4805          Top_mar=Ylength-Ygdu_cm*.40
4806
4807      ELSE
4808          ! 8.5x11 inch paper.
4809          Left_mar=Xgdu_cm*.11
4810          Right_mar=Xlength-Xgdu_cm*.24
4811          Bottom_mar=Ygdu_cm*.60
4812          Top_mar=Ylength-Ygdu_cm*.67
4813
4814      END IF
4815      RETURN
4816
4817      !
4818      !
4819      !Setup margins for left binding.
4820      IF Papersize="3" THEN
4821          ! 11x17 inch paper.
4822          Left_mar=Xgdu_cm*.38
4823          Right_mar=Xlength-Xgdu_cm*.145
4824          Bottom_mar=Ygdu_cm*.30
4825          Top_mar=Ylength-Ygdu_cm*.12
4826
4827      ELSE
4828          ! 8.5x11 inch paper.
4829          Left_mar=Xgdu_cm*.37
4830          Right_mar=Xlength-Xgdu_cm*.24
4831          Bottom_mar=Ygdu_cm*.60
4832          Top_mar=Ylength-Ygdu_cm*.140
4833
4834      END IF
4835      RETURN
4836
4837      !
4838      !
4839      !Set margins for fulltest dimensions.
4840      IF Papersize="3" THEN

```

```

4907 ! AX$ = "X" if only autoscale X axis.
4908 ! AX$ = "Y" if only autoscale Y axis.
4909 ! AX$ = "XY" for both.
4910 ! IF AX$ is not "XY" then GraphType$ must be set to a valid choice.
4911 ! IF AX$ is "NO" then autoscaling is disabled
4912 !
4913 ! NO SUB Programs are CALLED
4914
4915 ! Roster(i,1) = Curve number 1,2,3,....:17
4916 ! Roster(i,2) = Start address in File(x,*); = x
4917 ! Roster(i,3) = Datecount for curve i
4918 ! Roster(i,4) = PEN number for this curve
4919
4920 ! Symbol$(i)="" or "Y" => no symbol, connect pts
4921 ! Symbol$(i)="XY" => x symbol, connect pts
4922 ! Symbol$(i)="X" => x symbol, do NOT connect pts
4923 ! Lbl_addr: x, Y, pen, size, LDIR, LONG
4924
4925 COM /Labels/ Labels$(*),INTEGER Lbl_count,REAL lbl_addr(*)
4926 COM /data_param/ INTEGER Datecount,Filesize,Curvecount,Roster(*)
4927 COM /data_param/ REAL Sym_size,Symbol$(*),Curve_tds(*)
4928 COM /data_param/ REAL Xmin_data,Xmax_data
4929 COM /data_param/ REAL Ymin_data,Ymax_data
4930 COM /Background/ GraphType$,Margins(*),PaperSize$
4931 COM /Background/ REAL Pen_speed,INTEGER Background_pen,Auto_time
4932 COM /Background/ INTEGER Auto_file,REAL X_cross,Y_cross,X
4933 COM /Background/ Xgrid_ticks,INTEGER Xmajor,Xminor
4934 COM /Background/ Ygrid_ticks,INTEGER Ymajor,Yminor
4935
4936 COM /Axes_labels/ Print_xlabel$,Print_ylabel$,
4937 COM /Axes_labels/ Sig_digits$,REAL Xlsize,Ylsize
4938
4939 COM /WindowSpace/ REAL Xmin,Xmid,Xmax,Ymin,Ymid,Ymax!graph edges JDU
4940 COM /WindowSpace/ REAL Xleft,Xright,Ybottom,Ytop! paper edges LDU$
4941 COM /Plot device/ Plot Leng$,INTEGER Plot_addr
4942 COM /Hard space/ REAL Xview_ft,Xview_rt,Yview_btm,Yview_top
4943 COM /Hard space/ REAL ViewScale,Aspect_ratio
4944 COM /Frame_size/ INTEGER Frame_flag
4945 COM /Clear_space/ INTEGER Space_lift,Space_rt,Space_btm,Space_top
4946 COM /Bugs/ INTEGER Bug1,Bug2,Bug3,Printer
4947
4948 REAL Max_Min
4949 REAL Range,Factor,Base
4950 INTEGER Major_Minor
4951 INTEGER I,J,Axis
4952 DIM Xscales{6},Scales{6}
4953
4954 ! IF Curvecount<1 OR AX$="NO" THEN
4955 GRAPHICS OFF
4956 BEEP
4957 ! IF AX$="NO" THEN
4958 ! DISP " NO DATA AVAILABLE FOR AUTO SCALING "
4959 ELSE
4960 ! DISP " Auto scale is disabled ... see BACKGROUND editor."
4961 WAIT 2.0
4962 DISP CHR$(12)
4963 GRAPHICS ON
4964 SUBEXIT
4965 END IF
4966 ! IF AX$="XY" THEN
4967 SELECT GraphType$
4968 CASE "LINEAR"
4969 CASE "LINEAR"
4970 CASE "LINEAR"
4971 Xscales="LINEAR"
4972
4973 Scales$="LINEAR"
4974 CASE "SEMILOG X"
4975 Xscales="LOG"
4976 CASE "SEMILOG Y"
4977 Xscales="LINEAR"
4978 CASE "LINEAR"
4979 Scales$="LOG"
4980 CASE "LOG LOG"
4981 Xscales="LOG"
4982 Scales$="LOG"
4983 CASE ELSE
4984 DISP "Graph type improperly chosen! ";
4985 DISP " will be set to LINEAR ....continue"
4986 BEEP
4987 PAUSE
4988 DISP CHR$(12)
4989 Xscales="LINEAR"
4990 Scales$="LINEAR"
4991 END SELECT
4992 !
4993 Selective_scale:1
4994 IF AX$="X" OR AX$="Y" OR AX$="XY" THEN
4995 IF AX$="X" OR AX$="XY" THEN
4996 Axis=1
4997 GOSUB Find_max_min
4998 Xmin_data=Min
4999 Xmax_data=Max
5000 GOSUB Choose_scale! Either log or linear parameters
5001 Xscales=Scales
5002 GOSUB Set_X_initial
5003 IF
5004 IF AX$="Y" OR AX$="XY" THEN
5005 Axis=2
5006 GOSUB Find_max_min
5007 Ymin_data=Min
5008 Ymax_data=Max
5009 GOSUB Choose_scale! Either log or linear parameters -> Scales
5010 GOSUB Set_Y_initial
5011 END IF
5012 ELSE
5013 DISP " ERROR is setting axis selector for autoscale. ";
5014 DISP " will be set to XY ... continue "
5015 BEEP
5016 PAUSE
5017 DISP CHR$(12)
5018 AX$="XY"
5019 GOTO Selective_scale
5020 END IF
5021 ! Choose GraphType$
5022 SELECT Xscales
5023 CASE "LOG"
5024 SELECT Scales
5025 CASE "LOG"
5026 GraphType$="LOG LOG"
5027 CASE "LINEAR"
5028 GraphType$="SEMILOG X"
5029 CASE ELSE
5030 GraphType$="LINEAR"
5031 DISP "Auto select graph type error."
5032 BEEP
5033 PAUSE
5034 DISP CHR$(12)
5035 SELECT Scales
5036 CASE "LINEAR"
5037 CASE "LINEAR"
5038

```



```

5039 CASE "LOG"
5040 GraphType$="SEMILOG Y"
5041 CASE "LINEAR"
5042 GraphType$="LINEAR"
5043 CASE ELSE
5044 DISP "Auto select graph type error."
5045 BEEP
5046 PAUSE
5047 END SELECT
5048 CASE ELSE
5049 DISP "Auto select graph type error."
5050 BEEP
5051 PAUSE
5052 END SELECT
5053 SUBEXT1
5054 !
5055 !
5056 Find_max_min: !
5057 Min=File(Roster(1,2),Axis)
5058 Max=File(Roster(1,2),Axis)
5059 FOR I=1 TO Curvcount
5060 FOR J=Roster(I,2) TO Roster(I,3)-1
5061 Max=MAX(Max,File(J,Axis))
5062 Min=MIN(Min,File(J,Axis))
5063 NEXT J
5064 NEXT I
5065 IF Max<Min THEN Max=Min
5066 IF ABS(Max-Min)<1.0E-50 THEN
5067 IF ABS(Max)>1.0E-50 THEN
5068 Max=Max+.1*ABS(Max)
5069 ELSE Max=-.1
5070 END IF
5071 IF ABS(Min)>1.0E-50 THEN
5072 IF ABS(Min)-.1*ABS(Min)
5073 ELSE Min=-.1
5074 END IF
5075 END IF
5076 RETURN
5077 !
5078 !
5079 !
5080 !
5081 !
5082 !
5083 Choose scale: !
5084 Log_cutoff=(Max-Min)*.4+Min
5085 Number_below=0
5086 Total_data_pts=0
5087 FOR I=1 TO Curvcount
5088 Total_data_pts=Total_data_pts+Roster(I,3) \*Datacount
5089 FOR J=Roster(I,2) TO Roster(I,3)-1
5090 IF File(J,Axis)<Log_cutoff THEN
5091 Number_below=Number_below+1
5092 ELSE GOTO Next_crv
5093 NEXT J
5094 END IF
5095 Next_crv: !
5096 NEXT I
5097 IF Number_below/Total_data_pts<.5 OR Min<1.E-25 THEN
5098 Scale$="LINEAR"
5099 ELSE
5100 Scale$="LOG"
5101 RETURN
5102 END IF
5103 !
5104 ! Find parameters for a linear scale (Ideas by Wilber Lambert)
5105
5106 Range=Max-Min
5107 IF Range>1.0E-50 THEN
5108 Base=10 \*INT(LBT(Range))
5109 SELECT Range
5110 CASE <=2*Base
5111 Factor=Base/5
5112 CASE <=5*Base
5113 Factor=Base/2
5114 CASE <=10*Base
5115 Factor=Base
5116 END SELECT
5117 Factor=1
5118 END IF
5119 Min=Factor*\*INT(Min/Factor)
5120 Max=Factor*\*(INT(Max/Factor)+*(Max/Factor->INT(Max/Factor)))
5121 Major=INT((Max-Min)/Factor)
5122 Major=Major+\*(INT((Max-Min)/Factor)<>DROUND((Max-Min)/Factor,4))
5123 IF (INT(Major/2)=Major/2) THEN
5124 Major=Major/2
5125 Minor=2
5126 ELSE
5127 Minor=1
5128 END IF
5129 RETURN
5130 !
5131 !
5132 Set_x_initial: !
5133 Xmajor=Major
5134 Xminor=Minor
5135 Xmin_graph=Min
5136 Xmax_graph=Max
5137 IF Bug2 THEN
5138 PRINTER IS Printer
5139 PRINT "XMIN_GRAPH=";Xmin_graph,"XMAX_GRAPH=";Xmax_graph
5140 PRINT
5141 PRINTER IS CRT
5142 END IF
5143 Y_cross_x=Xmin_graph
5144 RETURN
5145 !
5146 !
5147 !
5148 Set_y_initial: !
5149 Ymajor=Major
5150 Yminor=Minor
5151 Ymin_graph=Min
5152 Ymax_graph=Max
5153 IF Bug2 THEN
5154 PRINTER IS Printer
5155 PRINT "YMIN_GRAPH=";Ymin_graph,"YMAX_GRAPH=";Ymax_graph
5156 PRINT
5157 PRINTER IS CRT
5158 END IF
5159 X_cross_y=Ymin_graph
5160 RETURN
5161 SUBEND
5162 !
5163 !
5164 !
5165 !
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5364 Date orgns=Measure_times(1,1)
5365 MAT S(1:Dcount,5,1)=Local_array(1:Dcount,5,1)
5366 FOR I=1 TO 4
5367 IF BIT(Sval(I,I-1)) THEN MAT S(1:Dcount,I,*)=Local_array(*,I,*)
5368 NEXT I
5369 DEALLOCATE Local_array(*)
5370 SUBEXIT !Load_sparms
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5373
5374 Read_sparm_err:
5375 BEEP
5376 PRINT "Error number: ",ERRN:ERRMS
5377 DISP "something is wrong with S-parameter read"
5378 PAUSE
5379 RETURN
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Date orgns=Measure_times(1,1)
MAT S(1:Dcount,5,1)=Local_array(1:Dcount,5,1)
FOR I=1 TO 4
IF BIT(Sval(I,I-1)) THEN MAT S(1:Dcount,I,*)=Local_array(*,I,*)
NEXT I
DEALLOCATE Local_array(*)
SUBEXIT !Load_sparms
////////////////////
5374 Read_sparm_err:
5375 BEEP
5376 PRINT "Error number: ",ERRN:ERRMS
5377 DISP "something is wrong with S-parameter read"
5378 PAUSE
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////////////////////
SUBEND !Load_sparms
*****
SUB Select_disk(OPTIONAL Prompts)
Select_disk: ! Original: 04 Jul 1987
! Revision: 06 Aug 1987
COM /Files/ Diskdrives(20), Filenames(10)
INTEGER Local_prty,Dd,Pt,Choose(1)
DIM Discs(30),Titles(40),Displs(80)
Local_prty=VAL(SYSTEMS("SYSTEM PRIORITY"))+1
OFF KEY
! Define the disk drives available for this system, reserve the
! first characters for the drive address and the characters after
! the - for a description of the drive.
! Example:
! Discs(1)="",700,0,0 HP 9133H HARD disk, volume 0."
!
!
IF NPAR THEN
Displs="" SELECT DISK DRIVE "8Prompts" ... Abort will cancel. "
Displs=Displs(1:80)
ELSE
Displs="" SELECT DISK DRIVE "... Abort will cancel. "
END IF
Titles="" Available disk drives for this system. "
Pt=1 ! allow only one select
!
IF Diskdrives(1,1)<="" THEN Diskdrives=""
IF Msi_ids(1,1)<="" THEN Msi_ids="" THEN Msi_ids=SYSTEMS("MST")
Diskdrives=TRIMS(Diskdrives)
Msi_ids=TRIMS(Msi_ids)
IF [LEN(Diskdrives)>20 AND LEN(Msi_ids)>0 THEN
Displs(1)=Displs(1)&" " & Msi_ids(1)
Displs(1)=Displs(1)&" " & Msi_ids(1)
DO=1
IF Diskdrives=Msi_ids THEN
Displs(2)=Msi_ids&DRPT(" ",17-LEN(Msi_ids))
Displs(2)=Displs(2)&" " & Msi_ids(1)
DO=DO+1
ELSE
Displs(1)=Displs(1)&" Start-up MSUS."
END IF
END IF
IF LEN(Msi_ids)>0 THEN

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5303 Ref_val=ROUND(((Max+Min)/2)*Scale*5),-2)
5304 RETURN
5305 SUBEND
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5363
Ref_val=ROUND(((Max+Min)/2)*Scale*5),-2)
RETURN
SUBEND
*****
CSUB Mrcmin(XC*), YC*), Sig(*), INTEGER Mdata(*), REAL A(*), INTEGER Ma(*), Lis
ta(*), Mfitt(*), REAL Cover(*), A(pha*), INTEGER Nca(*), REAL Chisq,Alambda, Sss, D
=> ebug)
SUB Load_sparms
OPTION BASE 1
COM /History/ Status$(1), Time orgns(8), Date chngs(11), Descriptions$(160)
COM /Files/ Diskdrives(20), Filenames(10)
COM /Ana_data/ INTEGER Prog_id, Sweep_type, Sweep_mode, Datacount
COM /Ana_data/ REAL Z0, Start_Stop
COM /Ana_data/ Titles$(55), Company$(30), Operator_names$(30), Measure_tin
=> es$(30)
COM /S_array/ INTEGER Dcount,Sval(I),REAL S(801,5,2)
! Sweep_type: 1=Start/stop, 2=center/span
! Sweep_mode: 1=ramp, 2=step, 3=frequency list
DIM Prompts$(30)
INTEGER I,Num_sparms
Load_sparms:
Prompts="8510 S-PARAMETER DATA"
CALL Select_disk(Prompts)
CALL Enterfilename("CAT",Prompts)
IF Filenames="" THEN
CAT Diskdrives
CALL Enterfilename("ABORT",Prompts)
IF Filenames="" THEN
BEEP
DISP "PROGRAM ABORTED"
STOP
END IF
END IF
DISP CHR$(129); Reading S-parameter data: ";CHR$(128);Filenames$;Disk
drives$
ON ERROR CALL Errortrap
ASSIGN ainfile TO Filenames&Diskdrives;FORMAT OFF
OFF ERROR
ON ERROR GOSUB Read_sparm_err
ENTER ainfile;Prog_id
ENTER ainfile;Titles
ENTER ainfile;Company$
ENTER ainfile;User_names$
ENTER ainfile;Measure_times$
ENTER ainfile;Z0
ENTER ainfile;Sweep_type
ENTER ainfile;Sweep_mode
ENTER ainfile;Datacount,Start_Stop,S11_valid,S21_valid,S12_valid,S22_
valid,Ports
Datacount:Datacount:
Datacount=S11_valid+S21_valid+S12_valid+S22_valid
Sval(I)=0
IF S11_valid THEN Sval(I)=Sval(I)+1
IF S21_valid THEN Sval(I)=Sval(I)+2
IF S12_valid THEN Sval(I)=Sval(I)+4
IF S22_valid THEN Sval(I)=Sval(I)+8
ALLOCATE REAL Local_array(Datacount,Num_sparms+1,2)
ENTER ainfile;Local_array(*)
OFF ERROR
ASSIGN ainfile TO *
Descriptions=Titles
Time_orgns=Measure_times$(15,22)

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5430 Disc$(1)=Msi_id$&RPT$( " ",17-LEN(Msi_id$))
5431 Disc$(1)=Disc$(1)&"- Start-up mass storage unit specifier."
5432 Dd=1
5433 ELSE
5434 Dd=0
5435 END IF
5436 END IF
5437 Disk:
5438 ..... customize system drives here .....
5439 Follow format with - after unit specifier, description is
5440 optional but recommended.
5441 .....
5442 Disc$(Dd+1)=": 707,0 - HP 9122 dual microfloppy left drive"
5443 Disc$(Dd+2)=": 707,1 - HP 9122 dual microfloppy right drive"
5444 Disc$(Dd+3)=": 1400,0 - HP 9127 single 5.25 floppy drive"
5445 Disc$(Dd+4)=": 1400,1 - HP 9153B single microfloppy"
5446 Disc$(Dd+5)=": 1400,0,2 - HP 9153H hard disk ANA volume"
5447 Disc$(Dd+6)=": 1400,0,3 - HP 9153B hard disk PROG volume"
5448
5449 Dd=Dd+6 ! add the number of drive specifiers above
5450
5451 IF Sys_id$(1,4) <> "S300" THEN
5452 Disc$(Dd+1)=": 4,1 - LEFT internal series 200"
5453 Disc$(Dd+2)=": 4,0 - RIGHT internal series 200"
5454 Df=Dd+2
5455 END IF
5456
5457 CALL Menu scroll(Disp$,Titles$,Disc$(*),Dd,Pt,Choose(*))
5458 IF Pt=0 THEN
5459 Diskdrives="NO DISK"
5460 ELSE
5461 Dds=POS(Disc$(Choose(Pt)),":")-1 ! find -
5462 IF Dd>5 THEN ! valid msus
5463 Diskdrives=TRIM$(Disc$(Choose(Pt)))[1,Dd]
5464 ELSE
5465 DISP " ERROR in reading MSUS from string, - chr not found. "
5466 .BEEP
5467 .PAUSE
5468 Diskdrives="NO DISK"
5469 END IF
5470 END IF
5471 DiskSelected=OFF KEY
5472 SUBEXIT
5473 SUBEND
5474
5475 *****
5476 !
5477 !
5478 !
5479 SUB Enterfilename(Acs$,OPTIONAL Prompts$)
5480 Enterfilename:
5481 COM /Files/ Diskdrives$,Filename$
5482 DIM Test$(160)
5483 SELECT NPAP
5484 CASE 1
5485 DISP " ENTER the FILE NAME ";
5486 CASE 2
5487 DISP " ENTER the FILE NAME for ",Prompts$;
5488 END SELECT
5489 SELECT Acs$
5490 CASE "CAT"
5491 DISP " ... (ENTER alone to CAT) ";
5492 CASE "ABORT"
5493 DISP " ... (ENTER alone to ABORT) ";
5494 CASE "VALID"
5495

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5496 END SELECT
5497 LINPUT Tests
5498 Test$=TRIM$(Test$)
5499 IF LEN(Test$)=0 THEN
5500 SELECT Acs$
5501 CASE "VALID"
5502 DISP "you MUST enter the FILE NAME now."
5503 .BEEP
5504 WAIT 1.8
5505 GOTO Enterfilename
5506 CASE "ABORT" "CAT"
5507 GOTO Abortline
5508 CASE ELSE "Act$=";Acs$," in SUB Enterfilename"
5509 .BEEP
5510 WAIT 1.8
5511 DISP LEN(Test$); " "
5512 .PAUSE
5513 END SELECT
5514 IF LEN(Test$)>10 THEN
5515 .BEEP
5516 DISP "ERROR in NAME ENTRY--up to 10 chars, you have ";
5517 DISP LEN(Test$); " "
5518 WAIT 1.8
5519 OUTPUT 2 USING "#,K,K,": " #";Test$
5520 GOTO Enterfilename
5521 END IF
5522 FileNames=Test$
5523 FOR I=1 TO LEN(FileNames$)
5524 Ascii_num=NUM(FileNames$(I))
5525 SELECT Ascii_num
5526 CASE 65 TO 97 TO 95,97 TO 122,48 TO 57
5527 ! allowed characters
5528 CASE ELSE
5529 .BEEP
5530 DISP "ERROR in NAME ENTRY--ILLEGAL CHARACTERS, TRY AGAIN. "
5531 WAIT 1.8
5532 OUTPUT 2 USING "#,K,K,": " #";FileNames$
5533 GOTO Enterfilename
5534 END SELECT
5535 NEXT I
5536 SUBEXIT
5537 Abortline:FileNames$=""
5538 SUBEXIT
5539 SUBEND
5540 !
5541 ! *****
5542 !
5543 SUB Errortrap
5544 Errortrap: ! Original: 13 Nov 1984
5545 ! Revision: 02 Dec 1987
5546 ! Trap most errors here
5547 COM /Files/ Diskdrives$(20),FileNames$(10)
5548 DIM File$(20),Test$(160),What$(20),Acs$(5)
5549 SELECT ERRN
5550 CASE 54
5551 DISP "DUPLICATE FILE NAME: ",FileNames$;
5552 DISP "....PURGE old one? (Y/N)";
5553 LINPUT What$
5554 What$=TRIM$(What$)
5555 SELECT What$(1,1)
5556 CASE "y","Y"
5557 PURGE FileNames&Diskdrives$
5558 Acs$="VALID"
5559 CALL Enterfilename(Acs$)
5560 END SELECT
5561

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5562 CASE 52,53
5563 DISP "Improper FILE NAME --- ENTER NEW FILE NAME";
5564 OUTPUT 2 USING "#,K,K;" #;filenames
5565 LINPUT filenames$
5566 filenames$=TRIMS(filenames$)
5567
5568 CASE 56
5569 DISP "FILE: ";filenames$;" is not on this disk, please insert";
5570 DISP " correct disk"
5571 CALL Pause_key_on
5572
5573 CASE 64
5574 DISP "This disk is full, PLEASE insert clean disk"
5575 CALL Pause_key_on
5576
5577 CASE 54
5578 DISP "DATA INPUT disk must be in drive!! ";
5579 DISP "...CONTINUE when ready."
5580 CALL Pause_key_on
5581
5582 CASE 72,73,76
5583 DISP Diskdrives$;
5584 DISP " is not available, type correct";
5585 DISP " unit specifier (ie. ':,707,0');";
5586 OUTPUT 2 USING "K,#";Diskdrives
5587 LINPUT Diskdrives$
5588
5589 CASE 80
5590 DISP "CHECK DISK drive door!"
5591 CALL Pause_key_on
5592
5593 CASE ELSE
5594 DISP ERMS;" CONTINUE" when fixed"
5595 CALL Pause_key_on
5596
5597 END SELECT
5598 DISP CHRS(12)
5599 SUBEXIT
5600 SUBEND
5601
5602 *****
5603 SUB Menu_scroll(D$,I$,Items$(*),INTEGER Item_cnt,To_select,Choose(*))
5604 Menu_scroll: I Original: 22 Jun 1987, Galen Koepke, NBS 723.04
5605 Revision: 30 Jun 1987, 13:55
5606
5607 I A general purpose menu utility for scrolling items and
5608 I selecting a given number of them.
5609 I The items are arranged in screens of 15 items each and
5610 I the user may access screens via softkeys. There may be
5611 I up to 10 screens or 150 items to choose from.
5612 I Items$(*) contains the item descriptions
5613 I Item_cnt is the number of items in Items$(*)
5614 I Choose(*) is dimensioned to the number of required choices
5615 I and will be filled with the item numbers chosen.
5616 I To_select is the number of required choices.
5617
5618 OPTION BASE 1
5619 PRINTER IS CRT
5620 DEG
5621 GOSUB Def_variables
5622 GOSUB Define_screens
5623 GOSUB Make_selections
5624 IF Null_file THEN I reset to zero
5625 Item_cnt=0
5626 Items$(1)="
5627 To_select=0 I no valid selections
5628 END IF
5629 SUBEXIT
5630
5631 I *****
5632 I Def_variables: I
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5694 ELSE Tests=" Select"
5695 END IF
5696 ON KEY K0 LABEL Tests,Local_prtY GOSUB Select_item
5697 ELSE ON KEY K0 LABEL " Exit",Local_prtY GOTO Exit_line
5698 END IF
5699 IF Active_screen<Screen_cnt THEN
5700 ON KEY K0+1 LABEL "Next Screen",Local_prtY GOSUB Next_screen
5701 ELSE OFF KEY K0+1
5702 END IF
5703 IF Active_screen>1 THEN
5704 ON KEY K0+2 LABEL " Last Screen",Local_prtY GOSUB Last_screen
5705 ELSE OFF KEY K0+2
5706 END IF
5707 IF Skips>0 THEN
5708 ON KEY K0+3 LABEL " select Reset",Local_prtY GOSUB Select_reset
5709 ELSE OFF KEY K0+3
5710 END IF
5711 ON KEY K0+4 LABEL " ABORT ",Local_prtY GOTO Escape_line
5712 GOTO Key_loop
5713 Escape_line:Skips=0
5714 MAT Choose= (0)
5715 To_select=0
5716 Exit_line:OFF KEY
5717 OFF KNOB
5718 OFF KBD
5719 OUTPUT KBD;CHRS(255)&CHRS(75);
5720 PRINT CHRS(128);
5721 I=anything cleared, now go back to work.
5722 RETURN
5723 I=
5724 I=
5725 I=
5726 I=
5727 I=
5728 I=
5729 I=
5730 I=
5731 Next_screen:
5732 OFF KBD
5733 OFF KNOB
5734 OFF KEY
5735 IF Active_screen=Screen_cnt THEN RETURN
5736 Active_screen=Active_screen+1
5737 GOSUB Write_screen
5738 RETURN
5739 I=
5740 I=
5741 I=
5742 Last_screen:
5743 OFF KBD
5744 OFF KNOB
5745 OFF KEY
5746 IF Active_screen=1 THEN RETURN
5747 Active_screen=Active_screen-1
5748 GOSUB Write_screen
5749 RETURN
5750 I=
5751 I=
5752 Select_item:
5753 OFF KBD
5754 OFF KNOB
5755 OFF KEY
5756 IF NOT Interactive THEN
5757 DISP "NO additional selections for this screen."
5758
5759
5760 BEEP
5761 WAIT 2
5762 DISP CHRS(12);
5763 RETURN
5764 IF Skips=To_select THEN
5765 To_select=0
5766 DISP "this menu is for information only.";
5767 DISP " no selection allowed."
5768 ELSE
5769 DISP "All selections have been filled.";
5770 END IF
5771 DISP " 'Select Reset' to repeat."
5772 BEEP
5773 WAIT 2
5774 DISP CHRS(12);
5775 RETURN
5776 Skips=Skips+1
5777 Choose(Skips)=first_item(Active_screen)+Pointer-First_line
5778 PRINT CHRS(129); I Inverse video
5779 PRINT TABXY(10,Pointer);Items$(Choose(Skips))
5780 PRINT TABXY(128);
5781 SELECT Pointer
5782 CASE First_line
5783 GOSUB Point_forward
5784 CASE Last_line
5785 GOSUB Point_backward
5786 CASE ELSE
5787 I move forward unless it requires wrapping to beginning.
5788 IF Skips=1+0 THEN I check for selected items.
5789 I=Pointer-First_line
5790 LOOP
5791 FOR J=1 TO Skips
5792 IF First_item(Active_screen)+1=Choose(J) THEN K=1
5793 NEXT J
5794 I=I+1
5795 IF I+First_line>Last_line THEN K=-1
5796 EXIT IF K=-1
5797 END LOOP
5798 IF K=0 THEN
5799 GOSUB Point_forward
5800 ELSE
5801 GOSUB Point_backward
5802 END IF
5803 GOSUB Point_forward
5804 GOSUB Point_backward
5805 END IF
5806 GOSUB Point_forward
5807 END IF
5808 END SELECT
5809 RETURN
5810 I=
5811 I=
5812 I=
5813 I=
5814 I=
5815 I=
5816 Select_reset:
5817 OFF KBD
5818 OFF KNOB
5819 OFF KEY
5820 Skips=0
5821 MAT Choose= (0)
5822 GOSUB Write_screen
5823 RETURN
5824 I=
5825 I=
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5826 Process_kbd: I Allow use of arrows and enter key in addition to soft.
5827 Test:=KBD$
5828 IF LEN(Test)=1 AND Test=1,1 THEN
5829 BEEP 80...1
5830 RETURN
5831 END IF
5832 IF Test=1,1 THEN GOSUB Point_forward
5833 IF Test=1,1 THEN GOSUB Point_backward
5834 SELECT Test$[2,2]
5835 CASE CHR$(255)
5836 I do nothing
5837 CASE "m","M"
5838 GOSUB Point_forward
5839 CASE "u","U"
5840 GOSUB Point_backward
5841 CASE "e"
5842 IF Skips=0 THEN GOSUB Select_item
5843 ELSE I exit routine
5844 Exit_flag=1
5845 END IF
5846 BEEP 80...1
5847 END SELECT
5848 Test="m"
5849 RETURN
5850 I
5851 I
5852 I
5853 I
5854 I
5855 I
5856 Point_forward:Knobcount=5
5857 GOSUB Move_pointer
5858 RETURN
5859 Point_backward:Knobcount=5
5860 GOSUB Move_pointer
5861 RETURN
5862 I
5863 I
5864 I
5865 Jog_pointer: I Move the selection pointer on the active screen.
5866 I Without regard to selected values
5867 IF Knobcount>0 THEN I Move forward
5868 Pointer=Pointer+1
5869 ELSE I Move backward
5870 Pointer=Pointer-1
5871 END IF
5872 IF Pointer<first_line THEN Pointer=first_line
5873 IF Pointer>last_line THEN Pointer=last_line
5874 RETURN
5875 I
5876 I
5877 I
5878 Move_pointer: Control pointer to avoid re-selection of items
5879 IF NOT Pointeractive THEN RETURN I No selections to be made.
5880 Knobcount=Knobcount+KnobX*KnobY
5881 IF ABS(Knobcount)<4 THEN RETURN
5882 Last_pt=Pointer
5883 GOSUB Jog_pointer
5884 IF Skips>0 THEN
5885 LOOP
5886 J=Pointer-First_line
5887 FOR I=1 TO Skips
5888 IF First_item(Active_screen)+J=Choose(I) THEN J=999
5889 IF J=999 AND Pointer=Last_pt THEN Pointeractive=0
5890 I
5891 I

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5892 EXIT IF Pointeractive=0
5893 IF J=999 THEN GOSUB Jog_pointer
5894 EXIT IF J<>999
5895 END LOOP
5896 Knobcount=0
5897 OUTPUT KBD;CHR$(255)&CHR$(84); I Bring screen home
5898 IF Last_pt=Last_line THEN PRINT CHR$(132);
5899 PRINT " ";
5900 IF Pointeractive THEN I Pointer active
5901 IF Pointer=Last_line THEN
5902 PRINT CHR$(132);
5903 ELSE
5904 PRINT CHR$(128);
5905 END IF
5906 PRINT TABXY(1,Pointer);Marker;CHR$(128);
5907 END IF
5908 RETURN
5909 I
5910 I
5911 I
5912 I
5913 Write_screen: I Write the screen pointed to by Active_screen
5914 home and clear screen
5915 OUTPUT KBD;CHR$(255)&CHR$(84)&CHR$(255)&CHR$(75);
5916 PRINT TABXY(1,First_line-1);CHR$(132);" Item #, Screen #";
5917 PRINT USING "#,2D,3A,2D,3A";Active_screen," of ";Screen_cnt;" | "
5918 PRINT I$;RPTS(" ",52-LEN(I$));CHR$(128);
5919 J=0
5920 REPEAT
5921 IF J=Last_item(Active_screen)-first_item(Active_screen) THEN
5922 PRINT CHR$(132);
5923 PRINT TABXY(1,First_line+J);RPTS(" ",80)
5924 ELSE
5925 PRINT CHR$(128);
5926 END IF
5927 PRINT TABXY(5,First_line+J);
5928 PRINT USING "#D,A,#M;First_item(Active_screen)+J,|"
5929 IF Skips=0 THEN I make this line inverse video
5930 FOR I=1 TO Skips
5931 IF First_item(Active_screen)+J=Choose(I) THEN
5932 PRINT CHR$(129);
5933 NEXT I
5934 J=J+1
5935 IF J=Last_item(Active_screen)-first_item(Active_screen)+J
5936 THEN
5937 PRINT TABXY(10,First_line+J);Items$(First_item(Active_screen)+J)
5938 J=J+1
5939 UNTIL J>=Last_item(Active_screen)-first_item(Active_screen)+1
5940 Last_line=Last_item(Active_screen)-first_item(Active_screen)
5941 Last_line=Last_line+first_line
5942 I set marker to first non-selected item.
5943 Pointeractive=0
5944 IF To_select=0 THEN Pointeractive=1
5945 IF Skips>0 AND Pointeractive=1 THEN I find first non-selected item
5946 I
5947 LOOP
5948 Pointer=first_line+J
5949 FOR I=1 TO Skips
5950 IF First_item(Active_screen)+J=Choose(I) THEN Pointer=0
5951 J=J+1
5952 IF Pointer=0 THEN
5953 I
5954 I
5955 I
5956 I
5957 I

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6021 B=.02286
6022 Coaxial_len=1.15
6023 Waveguide_len=.36
6024 Calc_delay:
6025   Cutoff_freq=C/(2*B) | X-BAND WAVEGUIDE
6026   FOR I=1 TO Num_Points
6027     Lambda=C/Freqency(I)
6028     Lambda_wg=Lambda/SQR(1-(Cutoff_freq/Freqency(I))^2)
6029     Phase_delay_cx=4*PI*Coaxial_len/Lambda
6030     Total_delay=Phase_delay_wg+Phase_delay_cx
6031     Real_delay=-SIN(Total_delay)
6032     IF BIT(Sval_id,0) THEN S(J,1,1)=S(J,1,1)*Real_delay
6033     IF BIT(Sval_id,0) THEN S(J,1,2)=S(J,1,2)*Imag_delay
6034     IF BIT(Sval_id,1) THEN S(J,2,1)=S(J,2,1)*Real_delay
6035     IF BIT(Sval_id,1) THEN S(J,2,2)=S(J,2,2)*Imag_delay
6036     IF BIT(Sval_id,2) THEN S(J,3,1)=S(J,3,1)*Real_delay
6037     IF BIT(Sval_id,2) THEN S(J,3,2)=S(J,3,2)*Imag_delay
6038     IF BIT(Sval_id,3) THEN S(J,4,1)=S(J,4,1)*Real_delay
6039     IF BIT(Sval_id,3) THEN S(J,4,2)=S(J,4,2)*Imag_delay
6040   NEXT J
6041   OEG
6042   SUBEND
6043 *****
6044 SUB Enter_id(ids,OPTIONAL Return_tests)
6045   Enter_id:
6046   !LAST REVISION 30/SEPT/86
6047   OPTION BASE 1
6048   COM /Bugs/ INTEGER Bug1,Bug2,Bug3,Printer
6049   DIM Tests(160)
6050   INTEGER N
6051   N=LEN(ids)
6052   Tests=ids
6053   SELECT ids
6054   CASE ""
6055     !OUTPUT NOTHING
6056   CASE ELSE
6057     OUTPUT 2 USING "K, #, Tests"
6058   END SELECT
6059   SELECT NPAR
6060   CASE 1 !NO Return tests given
6061     DISP CHR$(129); "Please ENTER a description (<= 40 chrs).";
6062     DISP CHR$(128);
6063   CASE ELSE
6064     DISP CHR$(129); "Please ENTER a description (<= 40 chrs) ";
6065     DISP CHR$(128);
6066   CASE 105
6067     DISP " for THIS ID";
6068   CASE "ABORT"
6069     DISP " CLR LN/ ENTER to ABORT."
6070   CASE ELSE
6071     DISP " for ";Return_tests;
6072   END SELECT
6073   INPUT Tests
6074   DISP ""
6075   Test$=TRIM$(Tests)
6076   N=LEN(Test$)
6077   SELECT N
6078   CASE >40

```

```

5958   Pointer=First_line
5959   END IF
5960   EXIT IF Pointer<>0
5961   END LOOP
5962   ELSE Pointer=First_line
5963   END IF
5964   IF Pointer=Last_line THEN
5965     PRINT CHR$(132);
5966   ELSE PRINT CHR$(128);
5967   END IF
5968   PRINT TAB$(1,Pointer);Markers;CHR$(128);
5969   END IF
5970   RETURN
5971   SUBEND
5972 *****
5973 SUB Edit_data(Prompts,Variable,OPTIONAL Multiplier,Uvariable)
5974   Edit_data:OFF KEY
5975   IF NPAR=2 THEN
5976     Test=Variable*Multiplier
5977     IF NPAR=4 THEN Uvariable=Uvariable*Multiplier
5978   ELSE Test=Variable
5979   END IF
5980   ON ERROR GOTO Test_again
5981   Test_again:
5982   DISP "Enter the value of ";Prompts;
5983   INPUT Variable
5984   OFF ERROR
5985   IF NPAR=4 THEN
5986     Utest=Uvariable
5987     ON ERROR GOTO Utest_again
5988   Utest_again:
5989   OUTPUT 2 USING "K, #, Utest"
5990   DISP "Enter the uncertainty in ";Prompts;
5991   INPUT Uvariable
5992   OFF ERROR
5993   IF NPAR=2 THEN
5994     Variable=Variable/Multiplier
5995     Uvariable=Uvariable/Multiplier
5996   END IF
5997   IF NPAR=4 THEN
5998     Variable=Variable/Multiplier
5999     Uvariable=Uvariable/Multiplier
6000   END IF
6001   SUBEND
6002 *****
6003 SUB Phase_delay
6004   Phase_delay:
6005   Init_com:
6006   OPTION BASE 1
6007   COM /Lab info/ REAL Rel_humidity, Temperature, Pressure, C
6008   COM /Sparms/ REAL Freq(801), S21(801), S12(801), S22(801)
6009   COM /Sparms/ Mag_s11_ids(40), Ang_s11_ids(40), Mag_s21_ids
6010   => (40)
6011   COM /Sparms/ Mag_s22_ids(40), Ang_s22_ids(40), Mag_s12_ids
6012   => (40)
6013   COM /Sparm/ Frequency(801), COMPLEX Sparm(3), INTEGER Num_poin
6014   => ts
6015   COM /S_array/ INTEGER Dcount, Sval_id, REAL S(801,5,2)
6016   Init_variables:
6017   END
6018 *****

```



```

6153 ELSE GOTO Enter_real
6154 END IF
6155 SELECT Test_real
6156 CASE <Low
6157 BEEP 1000,.3
6158 DISP " Number entered is TOO LOW. ";
6159 WAIT 2.1
6160 GOTO Enter_real
6161 CASE >High
6162 BEEP 1000,.3
6163 DISP " Number entered is TOO HIGH. ";
6164 WAIT 2.1
6165 GOTO Enter_real
6166 CASE ELSE
6167 Result=Test_real
6168 ! Number within limits
6169 END SELECT
6170 SUBEXIT
6171
6172
6173
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6177
6178 Trap_bad_number:
6179 SELECT ERRN
6180 CASE 15,32
6181 DISP CHR$(128); " what you ENTERED is not a number! Try again. ";
6182 DISP CHR$(128)
6183 Bad_number=1
6184 WAIT 1.7
6185 LINPUT "Please ENTER the number you wish for",Test$
6186 CASE ELSE
6187 DISP ERRN,ERRMS
6188 BEEP 850,.5
6189 Bad_number=1
6190 PAUSE
6191 END SELECT
6192 RETURN
6193
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6199
6200 SUB New_save_sparms(INTEGER Ssave)
6201 ! This subprogram saves the chosen S-parameters that are passed in common
6202 ! blocks. The B1's 0 through 3 of the pass parameter Ssave are used to
6203 ! select which S-parameters are to be saved.
6204 ! The S(I,J,K) array is (801,5,2) big. I=1:Ocount .. K=REAL/IMAG
6205 ! J=1:S11 .. J=2:S21 .. J=3:S12 .. J=4:S22 .. J=5:K=1:Frequency
6206 ! This array construct is the same as the NWA automated measurement program
6207 ! To save disk space when only interested in one S-parameter, other
6208 ! programs save only the frequency and the S-parameter. That is why the
6209 ! Basket array is a variable size that is determined by Num_sparms.
6210 ! If Basket contained all four S-parameters, its dimension would be
6211 ! (1,J)=(Ocount,9), where J=1:Frequency, 2/3:REAL/IMAG(S11), 4/5:REAL/IMAG(
6212 ! => S21), 6/7:REAL/IMAG(S12), 8/9:REAL/IMAG(S22)
6213 ! If Basket contained just S22 then (1,J)=(Ocount,3) where J=1:Frequency,
6214 ! J=2/3:REAL/IMAG(S22), which saves disk space.
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6087 DISP "Length of data_ids too long. You entered ",N;
6088 DISP " characters. Try again."
6089 WAIT 1.5
6090 IF NPAR=2 THEN
6091 IF IDS=<>Return_tests THEN
6092 OUTPUT 2 USING "#,K":Test$
6093 END IF
6094 GOTO Enter_id
6095 CASE =0
6096 IF NPAR=1 THEN
6097 IF Return_tests="ABORT" THEN
6098 IDS=Test$
6099 SUBEXIT
6100 END IF
6101 END IF
6102 DISP "you must ENTER SOMETHING or you'll ";
6103 DISP "never get out of this."
6104 BEEP
6105 WAIT 1.8
6106 GOTO Enter_id
6107 CASE ELSE
6108 I=Everything ok
6109 END SELECT
6110 IDS=Test$
6111 SUBEND
6112 !Enter_id
6113 ! *****
6114 !
6115 !
6116 ! SUB Save_file(f(*),INTEGER Datacount,ids)
6117 Save_file:
6118 OPTION BASE 1
6119 COM /Files/ Drives$,FileNames$
6120 ON ERROR CALL ErrorTrap
6121 DiskSpace=INT((3500*(Datacount*16))/256)*1
6122 CREATE BDAT FileNames&Drives$,DiskSpace,256
6123 CREATE ASCII FileNames&DiskDrives$,DiskSpace*2
6124 ASSIGN @Datapath TO FileNames&DiskDrives$
6125 OUTPUT @Datapath;"@"
6126 OUTPUT @Datapath;TRIMS(IDS)
6127 OUTPUT @Datapath;Datacount
6128 OUTPUT @Datapath;f(*)
6129 OUTPUT @Datapath;f(*)
6130 ASSIGN @Datapath TO *
6131 OFF ERROR
6132 SUBEND
6133 !Save_file
6134 ! *****
6135 !
6136 ! SUB Enter_real(Low,High,Prompt$,Result)
6137 OPTION BASE 1
6138 INTEGER Bad_number
6139 DIM Test$(160)
6140 REAL Test_real
6141 Enter_real:
6142 DISP "Please ENTER the desired value for";Prompt$:
6143 OUTPUT 2 USING "#,":Result
6144 LINPUT Test$
6145 Bad_number=0
6146 ON ERROR GOSUB Trap_bad_number
6147 Test_real=VAL(TRIMS(Test$))
6148 OFF ERROR
6149 IF Bad_number THEN
6150 IF Test$="" THEN
6151 SUBEXIT
6152

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6283 UNTIL !tst>=!tst_limit OR Force_converge
6284 Mrq_out:=OFF KEY
6285 Alambda=0
6286 CALL Mrgmin(X(*),Y(*),Sig(*),Ndata(*),A(*),Ma(*),Lista(*),Mfit(*),Cov
6287 => ar(*),Alpha(*),Nca(*),Chisq,Alambda,Sss,Debug)
6288 IF NOT Speed THEN
6289 GOSUB Plot_residuals
6290 END IF
6291 GOSUB Calc_sig
6292 GOSUB Calc_uncert
6293 GOSUB Print_results
6294 GOTO Subend
6295
6296
6297
6298
6299 Calc_uncert:
6300 CASE = 1
6301 Ua(1)=SQRT(Sig2*Covar(1,1)) IKx
6302 Ua(2)=SQRT(Sig2*Covar(2,2)) IKy
6303 Ua(3)=SQRT(Sig2*Covar(3,3)) IQ
6304 Ua(4)=SQRT(Sig2*Covar(4,4)) IF0
6305 Ua(5)=SQRT(Sig2*Covar(5,5)) !
6306 CASE = 2,3
6307 Ua(1)=SQRT(Sig2*Covar(1,1)) IUq0
6308 Ua(2)=SQRT(Sig2*Covar(2,2)) IUq1
6309 Ua(3)=SQRT(Sig2*Covar(3,3)) IUf0
6310 Ua(4)=SQRT(Sig2*Covar(4,4)) IUpHase
6311 Ua(5)=Ua(4)*360/2/PI ! in degrees
6312 Ua(3)=SQRT(Sig2*Covar(5,5)) IUmag (beta)
6313 CASE ELSE
6314 BEEP
6315 DISP "Sss=";Sss;" in Calc_uncert. PAUSED"
6316 PAUSE
6317 GOTO Calc_uncert
6318 END SELECT
6319 RETURN
6320
6321
6322 Calc_sig2:
6323 FOR I=1 TO Fitcount
6324 GOSUB Gen_cnum
6325 N=Start_index+I-1
6326 Sig2=Sig2+ABS(CPLX(S(N,1),S(N,NS,2))-Cnum)^2
6327 NEXT I
6328 Sig2=Sig2/(Fitcount-Mfit(2))
6329 RETURN
6330
6331
6332 Print_results:
6333 PRINTER IS 1
6334 PRINT Descriptions;" fit to ";Sparms;" values."
6335 PRINT USING "K,30,2(3X,K,MD,3DE)";Iteration #;" ,K,"Chi_squared=";Chis
6336 => q,"Alambda=";Alambda
6337 PRINT USING "30,K,30,K,30r;Dcount"; data points in file. Start_index
6338 => " ,Start_index," Stop_index=";Stop_index
6339 GOSUB Print_a
6340 GOSUB Print_u
6341 PRINT USING "K,D,4DE";"Chi_squared=";Chisq
6342 IF Printer ON THEN
6343 PRINTER IS Print_addr
6344 GOSUB Print_a
6345 GOSUB Print_u
6346 PRINT USING "K,D,4DE";"Chi_squared=";Chisq

```

```

6217 INTEGER Num_sparms,1
6218 Num_sparms:=BIT(SSave,0)+BIT(SSave,1)+BIT(SSave,2)+BIT(SSave,3)
6219 ALLOCATE Basket(1:Dcount,Num_sparms*2+1)
6220 New_save_sparms:=OFF KEY
6221 MAT Basket(1:Dcount,1)= S(1:Dcount,5,1) IF Frequency
6222 Column=2
6223 FOR I=1 TO 4
6224 IF BIT(SSave,I-1) THEN ISAVE S(11,S21,S12,S22
6225 MAT Basket(1:Dcount,Column)= S(1:Dcount,I,1)
6226 MAT Basket(1:Dcount,Column+1)= S(1:Dcount,I,2)
6227 Column=Column+2
6228 END IF
6229 NEXT I
6230 DISP CHR$(131);" Saving: ";CHR$(128);Filename$;Diskdrives
6231 DiskSpace=INT(Dcount*16.*(Num_sparms*2+1)/256)*20
6232 ON ERROR GOSUB Trap_error
6233 CREATE BOAT Filename$&Diskdrives$,DiskSpace,256
6234 CREATE ASCII Filename$&Diskdrives$,DiskSpace
6235 OFF ERROR
6236 ASSIGN abdatapath TO Filename$&Diskdrives$
6237 OUTPUT abdatapath;Dcount,SSave
6238 OUTPUT abdatapath;Basket(*)
6239 ASSIGN abdatapath TO *
6240 DEALLOCATE Basket(*)
6241 DISP
6242 SUBEXIT
6243 Trap_error: DISP ERRMS
6244 BEEP
6245 PAUSE
6246 RETURN
6247 Purge_old:PURGE Filename$&Diskdrives$
6248 BEEP
6249 PRINT "I just purged ";Filename$;Diskdrives$
6250 RETURN
6251 SUBEND
6252
6253
6254
6255 SUB Fit_sparms
6256 Fit Lorentz(Kx,Ky,Q,F0,Ukx,Uky,Uq,Uf0) Sss=1
6257 or Fit Lorentz(Q,Q,F0,Phase,Uq0,Uq1,Uf0,Uphase) Sss=2 or 3
6258 !Written by Eric J. Vanzura
6259 !Levenberg-Marquart algorithm CSUB taken from:
6260 !Numerical Recipes : the art of Scientific Computing
6261 !W. H. Press, B. P. Flannery, S. A. Teukolsky, W. T. Vetterling
6262 !Cambridge Press 1987
6263 Fit_sparms:=OFF KEY
6264 GOSUB Init
6265 GOSUB Init_est
6266 GOSUB Initial_plot
6267 REPEAT
6268 Do_mrq:
6269 GOSUB Printlatestvals
6270 Ochiq=Chisq
6271 MAT Old_e=A
6272 DISP "Performing Levenberg-Marquart Least-square-fit"
6273 OFF KEY
6274 CALL Mrgmin(X(*),Y(*),Sig(*),Ndata(*),A(*),Ma(*),Lista(*),Mfit(*)
6275 => ,Covar(*),Alpha(*),Nca(*),Chisq,Alambda,Sss,Debug)
6276 GOSUB Convergencekeys
6277 MAT Old_est=Sgen
6278 GOSUB Gen_estmate
6279 IF NOT Speed THEN
6280 GOSUB Plot_estimate
6281 END IF
6282 GOSUB Ask_converge

```

```

6347 PRINTER IS 1
6348 END IF
6349 Reset_keys=1
6350 RETURN
6351
6352
6353
6354 Initial_Plot:|
6355 ALLOCATE File(Fitcount*3,2),Old_est(fitcount,2),Orig_sparm(fitcount,2)
6356 =>
6357 ALLOCATE Sgen(fitcount,2)
6358 MAT Orig_sparm(*,1:2)= S(Start_index:Stop_index,ns,1:2)
6359 GOSUB Gen_estimate
6360 MAT Old_est= Sgen
6361 IF NOT Speed THEN GOSUB Plot_estimate
6362 RETURN
6363
6364
6365 Gen_cnum:|
6366
6367 SELECT Sss
6368 CASE 1 !S21
6369 Cnum=X(1)^2*CMPLX(A(1),-A(2))/CMPLX(X(1)^2-A(4)^2,-X(1)*A(4)/A(3)
6370 )
6371 CASE 2 !S11
6372 Degrees=360*A(4)/2/PI
6373 P=CMPLX(COS(Degrees),SIN(Degrees))
6374 Q1Zf=2*A(2)*X(1)-A(3))/A(3)
6375 Cnum=A(5)*CMPLX(A(2)/A(1)+Q1Zf*Q1Zf,A(1)+Q1Zf*Q1Zf*(1-A(2)/A(1)))**P/(1+Q1Zf
6376 )
6377 CASE 3 !S21
6378 Degrees=360*A(4)/2/PI
6379 P=CMPLX(COS(Degrees),SIN(Degrees))
6380 G0gl=1-A(1)/A(2)
6381 Dum=CMPLX(1,2)*X(1)-A(3))/A(2)
6382 Cnum=A(5)*G0gl*P/Dum
6383 CASE ELSE
6384 BEEP
6385 DISP "Sss=";Sss;" is not valid in Gen_cnum."
6386 PAUSE
6387 END SELECT
6388 RETURN
6389
6390
6391 Gen_estimates:|
6392 DISP "Generating latest estimate values"
6393 SELECT Sss
6394 CASE =1 !Hussenveig Lorentzian
6395 GOSUB Gen_cnum
6396 Sgen(1,1)=REAL(Cnum)
6397 Sgen(1,2)=IMAG(Cnum)
6398 NEXT 1
6399 FOR I=1 TO Fitcount
6400 GOSUB Gen_cnum
6401 Sgen(I,1)=REAL(Cnum)
6402 Sgen(I,2)=IMAG(Cnum)
6403 NEXT I
6404 FOR I=1 TO Fitcount
6405 GOSUB Gen_cnum
6406 Sgen(I,1)=REAL(Cnum)
6407 Sgen(I,2)=IMAG(Cnum)
6408 NEXT I
6409
6410
6411 NEXT I
6412 CASE ELSE
6413 DISP "Sss=";Sss;" is not valid in Gen_estimate."
6414 BEEP
6415 PAUSE
6416 GOTO Gen_estimate
6417 END SELECT
6418 RETURN
6419
6420
6421 Plot_estimate:|
6422 Pack_data(File(*),Old_est(*),Fitcount,Pen,Id$)
6423 Id$=Sparms
6424 Pen=2
6425 Pack_data(File(*),Orig_sparm(*),Fitcount,Pen,Id$)
6426 Id$="Old estimate"
6427 Pen=4
6428 Pack_data(File(*),Old_est(*),Fitcount,Pen,Id$)
6429 Id$="latest estimate"
6430 Pen=6
6431 Pack_data(File(*),Sgen(*),Fitcount,Pen,Id$)
6432 Tests="Real Part"
6433 Prompts="Imaginary Part"
6434 Id$="Results from file:"&Rdatafiles
6435 Init_graphics(Tests,Prompts,Id$)
6436 Ax$="X1"
6437 Autoscale(File(*),AX$)
6438 Refresh Labels
6439 Disp "S-parameters are: Red=measured data, Green=old estimate, Blue=L
6440 atest Estimate"
6441 Plot all(file(*))
6442 RETURN
6443
6444
6445
6446
6447
6448
6449
6450 Plot_residuals:|
6451 Pack_data(File(*))
6452 ALLOCATE Reres(fitcount,2),Imres(fitcount,2)
6453 FOR I=1 TO Fitcount
6454 Reres(I,1)=S(Start_index+1-1,5,1)/1.E+9
6455 Imres(I,1)=S(Start_index+1-1,5,1)/1.E+9
6456 Reres(I,2)=Sgen(I,1)-Orig_sparm(I,1)
6457 Imres(I,2)=Sgen(I,2)-Orig_sparm(I,2)
6458 NEXT I
6459 Id$="real residuals"
6460 Pen=2
6461 Pack_data(File(*),Reres(*),Fitcount,Pen,Id$)
6462 Id$="imag residuals"
6463 Pen=4
6464 Pack_data(File(*),Imres(*),Fitcount,Pen,Id$)
6465 Tests="Frequency (GHz)"
6466 Prompts="Residual"
6467 Id$="red=real green=imaginary"
6468 Init_graphics(Tests,Prompts,Id$)
6469 Ax$="X1"
6470 Autoscale(File(*),AX$)
6471 Refresh Labels
6472 Plot_all(File(*))
6473 RETURN
6474
6475
6476
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6479
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6481
6482

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6483 Init:=1
6484 DEG
6485 OPTION BASE 1
6486 INTEGER I,Prty,Test,Test_Limit,NS
6487 INTEGER Pen,Dome,Fitcount
6488 COMPLEX Chum,P,Dum
6489 REAL Debug
6490 DIM Test$(160),Prompt$(40),Id$(40),Rawdata$(160),PaperSize$(1)
6491 COM /History/ Status$(1),Time_orgn$(3),Date_orgn$(1)
6492 COM /History/ Time_chng$(3),Date_chng$(1),Description$(160)
6493 COM /Label$/ Label$(30),ID,INTEGER Lbl_Count,REAL Lbl_addr$(30,6)
6494 COM /Data_param/ INTEGER Data_Count,FileSize,CurveCount,Roster$(17,4)
6495 COM /Data_param/ REAL Sym_size,Symbol$(17),Curve_ID$(17),ID
6496 COM /Data_param/ REAL Min_data,Max_data
6497 COM /Data_param/ REAL Ymin_data,Ymax_data
6498 COM /Graphics/ INTEGER Speed
6499 COM /Background/ graphType$(12),Margins$(2),ID, PaperSize$(1)
6500 COM /Background/ REAL Pen_speed,INTEGER Background_pen,Auto_time
6501 COM /Background/ INTEGER Auto_file,REAL X_Cross,Y_Cross,X
6502 COM /Background/ Xprid_tick$(4),INTEGER Xmajor,Yminor
6503 COM /Background/ Xprid_tick$(4),INTEGER Ymajor
6504 COM /Background/ REAL Train_graph,Xmat_graph,Ymin_graph,Ymax_graph
6505 COM /Axes/ Labels/,String$(12),REAL Xprid_tick$(4)
6506 COM /Axes/ Labels/,REAL Xmin,Xmid,Xmax,Ymin,Ymid,Ymax
6507 COM /Windowspace/ REAL Xleft,Xright,Bottom,Top,paper_edges,UDUS
6508 COM /Windowspace/ REAL Xleft,Xright,Bottom,Top,paper_edges,UDUS
6509 COM /Log_scale/ REAL Xcycles,Xoeigin,Xcycles,Yoeigin
6510 COM /Plot_device/ Plot_Langs$(10),INTEGER Plot_addr
6511 COM /Hard_space/ REAL View_Lft,View_Rt,View_Btm,View_Top
6512 COM /Hard_space/ REAL ViewScale,Aspect_Ratio
6513 COM /Frame_onoff/ INTEGER Frame_Flag
6514 COM /Clear_space/ REAL Space_Lft,Space_Rt,Space_Btm,Space_Top
6515
6516 Init_com:=1
6517 COM /FitType/ INTEGER Sselect,Start_Index,Stop_Index
6518 COM /5 array/ INTEGER Dcount,Sval$(5,2)
6519 COM /Data/ INTEGER Mfit(2),List$(5,2),REAL A(5),Jn(5),INTEGER
=> Ndata(2)
6520 COM /Gsets/ Alanda,Chisq,INTEGER Ncat(2),REAL Alpha(5,5),Covar(5,5)
6521 COM /Output/ INTEGER Print_addr,Plotter,Printer_on,Plotter_on
6522 Init_vals:=1
6523 Speed=0
6524 K=0
6525 Iter=0
6526 Iter_Limit=3
6527 Plot_addr=3
6528 Plot_Langs="INTERNAL"
6529 M=3 IDimension of Covar,Alpha,Lista & A matrices
6530 Debug=0
6531
6532 IF NOT BIT(Sval$(2),Sselect) THEN
6533 BEEP
6534 DISP "The S(*) array for ",Sparms;" is not valid (doesn't have da
=> ter?) PAUSED"
6535 PRINT "The S(*) array for ",Sparms;" is not valid (doesn't have d
=> atar?) PAUSED"
6536 PAUSE
6537 END IF
6538 $$$=1 S21=Nussenzweig $$$=2 S11-Schulten $$$=3 S21-Schulten
6539 SELECT Sselect
6540 CASE 0 IS11
6541 $$$=2
6542 Sparms="S11"
6543 CASE 1 IS21
6544 $$$=1
6545 $$$=3
6546
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6618 Nca(2)=M
6619 Alambda=-2
6620 ALLOCATE REAL X(fitcount*2),Y(fitcount*2),Sig(fitcount*2)
6621 MAT SIG= (1)
6622 MAT X(fitcount)= S(Start_index:Stop_index,5,1)
6623 MAT Y(fitcount)= S(Start_index:Stop_index,5,1)
6624 MAT I(fitcount)= S(Start_index:Stop_index,Ns,1)
6625 MAT Y(fitcount+1:fitcount*2)= S(Start_index:Stop_index,Ns,2)
6626 RETURN Init_est
6627
6628
6629
6630
6631 Done_return:OFF KEY
6632 Dome=1
6633 RETURN
6634
6635
6636
6637 Force_converge:PRINT "Convergence will be forced."
6638 Force_converge=1
6639 RETURN
6640
6641
6642
6643 Preventconverge:PRINT "Convergence will be prevented."
6644 Itest=0
6645 RETURN
6646
6647
6648
6649 Graph_on: 1
6650 Speed=0
6651 ON KEY 4 LABEL "GRAPH OFF",Prty GOSUB Graph_off
6652 RETURN
6653
6654
6655
6656 Graph_off: 1
6657 Speed=1
6658 OFF KEY 4
6659 ON KEY 4 LABEL "GRAPH ON",Prty GOSUB Graph_on
6660 RETURN
6661
6662
6663
6664 Debug_on:PRINT "Debug flag is turned on."
6665
6666
6667 Debug_off:PRINT "Debug flag is turned off."
6668
6669 OFF KEY 6
6670 ON KEY 6 LABEL "DEBUG OFF",Prty GOSUB Debug_off
6671 RETURN
6672
6673
6674
6675 Debug_off:PRINT "Debug flag is turned off."
6676
6677 OFF KEY 6
6678 ON KEY 6 LABEL "DEBUG ON",Prty GOSUB Debug_on
6679 RETURN
6680
6681
6682
6683 Ask_converge=1
6684 SELECT Chisq
6685 CASE >Chisq

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6685 Itest=0
6686 PRINT USING "K,D,DE":New Chisq is greater than old Chisq by:";Ch
6687
6688 Itest=0
6689 PRINT USING "K,D,DE":New Chisq is equal to old Chisq"
6690
6691 CASE <Chisq
6692 PRINT USING "K,D,DE":New Chisq is LESS than old Chisq by:";Chisq
6693
6694 CASE <Chisq
6695 PRINT USING "K,D,DE":New Chisq is LESS than old Chisq by:";Chisq
6696
6697 IF ABS((Ochisq-Chisq)/Chisq)<.003 THEN Itest=Itest+1
6698 IF ABS((Ochisq-Chisq)/Chisq)<.00003 THEN Itest=Itest+1
6699 PRINT USING "K,D,DE":Itest
6700 PRINT USING "K,D,DE":ABS((Ochisq-Chisq)/Chisq)=";ABS((Ochisq-Chi
6701 sq)/Chisq)
6702 END SELECT
6703 RETURN
6704
6705
6706
6707 Convergencekeys:1
6708 PrtyVAL(SYS$(SYSTEM_PRIORITY))+1
6709 ON KEY 0 LABEL "FORCE CONVERGE",Prty GOSUB Force_converge
6710 ON KEY 2 LABEL "PREVENTCONVERGE",Prty GOSUB Preventconverge
6711 IF Speed THEN
6712 ON KEY 4 LABEL "GRAPH ON",Prty GOSUB Graph_on
6713 ELSE
6714 ON KEY 4 LABEL "GRAPH OFF",Prty GOSUB Graph_off
6715 END IF
6716 IF Debug THEN
6717 ON KEY 6 LABEL "DEBUG OFF",Prty GOSUB Debug_off
6718 ELSE
6719 ON KEY 6 LABEL "DEBUG ON",Prty GOSUB Debug_on
6720 END IF
6721 RETURN
6722
6723
6724
6725
6726
6727 Print a:1
6728
6729
6730
6731 Print a:1
6732
6733
6734
6735 IF A(4)<1.E+10 THEN
6736 PRINT USING A f(10g):"QL=";A(3) "FO=";A(4)/1.E+9 "FO+(Span/2
6737 )=";((A(4)+A(4)/2.E+9)/A(3))/1.E+9 "Span=";A(4)/A(3)/1000
6738 ELSE PRINT USING A f(10g):"QL=";A(3) "FO=";A(4)/1.E+9 "FO+(Span/2
6739 )=";((A(4)+A(4)/2.E+9)/A(3))/1.E+9 "Span=";A(4)/A(3)/1000
6740 END IF
6741 CASE =2,=3
6742 PRINT USING "2K,MD,4DE,2X),K,MD,9DE,2X,K,M3D,DD,2X,K,MD,3DE":;"Q0
6743 = " A(1) "Q1=" A(2) "FO=" A(3) "Degrees=";A(4)*360/2/PI "Mag=" A(5)
6744 CASE ELSE
6745 DISP "sss=";sss;" is not valid in print_a_ program paused"
6746 BEEP
6747 PAUSE
6748 END SELECT
6749 RETURN
6750 A f(10g):IMAGE K,60,X,K,DD,9D,X,K,DD,9D,X,K,DD,9D,X,K,3D,3D
6751 A f(10g):IMAGE K,60,X,K,DD,9D,X,K,DD,9D,X,K,DD,9D,X,K,3D,3D
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6745 | ////////////////////////////////////////////////////
6746 | ////////////////////////////////////////////////////
6747 | ////////////////////////////////////////////////////
6748 | Print ua:1
6749 | SELECT SSS
6750 | CASE =1
6751 | PRINT USING "2(K,MD,30E,4X),K,M5D,3X,K,3X,K,MD,30E","Ua(1),U
=> ky=" Ua(2), "Ua(3), "Ua(4)
6752 | CASE =2,=3
6753 | PRINT USING "2(K,MD,30E,2X),K,MD,30E,2X,K,M5D,30,2X,K,MD,30E","Ua
=> 0=" Ua(1), "Ua(1)=" Ua(2), "Ua(3), "Ua(4), "Ua(5)
6754 | CASE ELSE
6755 | DISP "SSS=" SSS; " is not valid in Print ua. program paused"
6756 | REP
6757 | PAUSE
6758 | END SELECT
6759 | RETURN
6760 | SUBEND
6761 | *****
6762 | *****
6763 | *****
6764 | *****
6765 | *****
6766 | *****
6767 | *****
6768 | *****
6769 | *****
6770 | *****
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6807 | *****
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6825 | *****

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i Original: 02 Dec 1987
i Revision: 02 Dec 1987
COM /Sys/ Sys: fcs(10)
If Sys: fcs(1,4)=$300" THEN i reset to $300 system keys
CONTROL KBD,15:0
CONTROL CRT,12:2
LOAD KEY
END IF
PAUSE
If Sys: fcs(1,4)=$300" THEN i set to $200 compatible keys
OUTPUT KBD USING "K,##"SCRATCH KEY X"
CONTROL KBD,15:1
CONTROL CRT,12:0
END IF
SUBEXIT
SUBEND
i *****
i *****
i *****

```

Appendix E

CAVITYPROG Program Listing


```

257 SELECT TRIMS(Sys_ids$)
258 CASE "PC300" "ED1KEYS:DOS,C"
260 LOAD KEY "ED1KEYS:DOS,C"
261 CASE "S300"
262 LOAD KEY "/BASIC/ED1KEYS:,1400,0"
263 END SELECT
264 STOP
265
266
267
268 Menu main:
269
270 CLEAR SCREEN
271 PRINT USING Image6;Sample_ids$
272 PRINT USING Image7;SampleLen*1.E+3;UsampleLen*1.E+3
273 PRINT USING Image8;Eps_re_guess
274 PRINT USING Image1;Degc,Udegc,Relh,Ureth
275 PRINT USING Image2;Mbar,Umbar
276 PRINT USING Image3;C,Uc
277 PRINT
278 PRINT USING Image4;Diameter*1.E+3,Udiameter*1.E+3
279 PRINT USING Image5;Length0*1.E+3,Ulength*1.E+3
280 DISP "Main program menu (select from options given below)"
281 RETURN
282
283
284
285
286
287
288
289
290
291
292
293
294
295 Image1:IMAGE "Temperature = ",20.20," +/- ",0.20," deg C Relative Humidity
=> = ",20.0," +/- ",0.0," %u
296
297
298 Image2:IMAGE "Barometric Pressure = ",30.0," +/- ",0.0," mbar"
=> = ",0.70652," +/- ",0.3052," meters/second
299
300 Image3:IMAGE "Speed of Light
=> = "
301
302 Image4:IMAGE "Resonator Diameter
Length (micr. zero) = ",20.50," +/- ",0.50," mm"
303
304 Image5:IMAGE "Sample
Length description: ",A,U," +/- ",0.50," mm"
305
306 Image6:IMAGE "Sample
Length = ",20.40," +/- ",0.40," mm"
307
308 Image7:IMAGE "
Epsilon Guess = ",30.40
309
310 Image8:IMAGE "
311
312
313
314
315
316
317
318 End program:OFF KEY
319 Prior_menu=1
320 RETURN
321
322
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496 Image3:IMAGE "Speed of Light" = "D.70ESZ," +/- "D.30ESZ," +/- "D.30ESZ," meters/second
498 Image4:IMAGE "Resonator Diameter" = "20.5D," +/- "D.5D," mm
500 Image5:IMAGE "Length (micr. zero)" = "30.5D," +/- "D.5D," mm
502 Image6:IMAGE "Sample Description: "K" = "20.4D," +/- "D.4D," mm
504 Image7:IMAGE "Initial guess for epsilon:" = "20.4D," +/- "D.4D," mm
506 Image8:IMAGE "Length" = "20.4D," +/- "D.4D," mm
508 Image9:IMAGE "Empty" = "20.4D," +/- "D.4D," mm
510 Image10:IMAGE "20X, "with Sample" = "2(52D.4D)," +/- "D.4D, 6X)," mm
512 Image11:IMAGE "Micrometer reading" = "2(2D.7D)," +/- "D.7D, 3X)," GHz
514 Image12:IMAGE "Resonance frequency" = "2(5D)," +/- "D.5D, 11X)," GHz
516 Image13:IMAGE "q value" = "2(5D.2D)," +/- "D.5D, 20)," +/- "D.5D, 20, 6X)," B, "kHz"
518 Image15:IMAGE "Epsilon" = "30.4D," +/- "D.4D, 5X)," Epsilon" = "20.5D," +/- "D.5D"
520 Image16:IMAGE "tan delta" = "D.6D," +/- "D.6D"
522
524
526
528
530
532
534
536 SUB Var_freq
538 Var_freq:OFF KEY
540 OPTION BASE 1
542 COM /Environment/ Degc, Mbar, Reth, C
544 COM /Environment/ Udegc, Umbar, Ureth, Uc
546 COM /sample/ Sample_ids[55], SampleLen, Eps_re, Eps_im, Tand
548 COM /sample/ UsampleLen, Ueps_re, Ueps_im, Utand, Eps_re_guess
550 COM /sample/ Errs[12], ErrTand[12]
552 COM /cavity dims/ Diameter, LengthD, Spectrum(40, 5), INTEGER Num_modes
554 COM /cavity vals/ Freq(2), Micr(2), Q(2), Bw(2)
556 COM /cavity vals/ Ufreq(2), Umicr(2), Uq(2), Ubw(2)
558 COM /System_status/ INTEGER Done, Prior_menu, Do_var_freq, Empty_cavity,
560 Prompt$(80)
562 COM /constants/ Bessel_root
564
566
568
570
572 PRty=VAL(SYSTEM_PRIORITY)
574 ON KEY 1 LABEL "Prior Menu", Prty CALL Prior_menu
576 ON KEY 2 LABEL "Sample Infor", Prty CALL Sample_info
578 ON KEY 3 LABEL "Resonance Freqs", Prty CALL Resonance_freqs
580 ON KEY 4 LABEL "Micr. Reading", Prty CALL Micr_reading
582 ON KEY 5 LABEL "q Values", Prty CALL q_values
584 ON KEY 6 LABEL "Band Widths", Prty CALL Band_widths
586 ON KEY 7 LABEL "Calc Epsilon", Prty CALL Calc_eps_freq
588 ON KEY 8 LABEL "Calc Errors", Prty CALL Calc_errs_freq
590 Done=0
592 Prior_menu=0
594 LOOP
596 IF Done THEN GOTO Var_freq
598 EXIT IF Prior_menu
600 END LOOP
602 OFF KEY
604 Prior_menu=0
606 Done=1
608 SUBEXIT
610
612 Calc_eps_freq:OFF KEY
614 Do_var_freq=1
616 CALL Calc_eps_re
618 CALL Calc_tand

```

```

620 CALL Calc_eps_im
622 CALL Calc_errs(1)
624 CALL Calc_err_anal
626 Done=1
628 RETURN
630
632
634
636 Calc_errs_freq:OFF KEY
638 Do_var_freq=1
640 CALL Calc_errs
642 Done=1
644 RETURN
646
648
650
652 Menu_var_freq:
654 CLEAR SCREEN
656 PRINT USING Image6; Sample_ids
658 PRINT USING Image9; Eps_re_guess, "Method:Variable Frequency"
660 PRINT USING Image7; SampleLen*1.E+3, UsampleLen*1.E+3
662 PRINT
664 PRINT USING Image1; Degc, Udegc, Reth, Ureth
666 PRINT USING Image2; Mbar, Umbar
668 PRINT USING Image3; C, Uc
670 PRINT
672 PRINT USING Image4; Diameter*1.E+3, Udiagram*1.E+3
674 PRINT USING Image5; LengthD*1.E+3, Ulength*1.E+3
676 PRINT
678 PRINT USING Image10ab
680 PRINT USING Image11b; Freq(1)/1.E+9, Ufreq(1)/1.E+9, Freq(2)/1.E+9, Ufreq(2)/1.E+9
682 PRINT USING Image11a; Micr(1)*1.E+3, Umicr(1)*1.E+3, Micr(2)*1.E+3, Umicr(2)*1.E+3
684 PRINT USING Image12ab; Q(1), Uq(1), Q(2), Uq(2)
686 PRINT USING Image13ab; Bw(1)/1.E+3, Ubw(1)/1.E+3, Bw(2)/1.E+3, Ubw(2)/1.E+3
688
690
692
694
696
698 Image1:IMAGE "Temperature" = "20.2D," +/- "D.2D," deg C
699 Image2:IMAGE "Relative Humidity" = "20.0," +/- "D.0," %
700 Image3:IMAGE "Barometric Pressure" = "30.0D," +/- "D.0D," mbar
702 Image4:IMAGE "Speed of Light" = "D.70ESZ," +/- "D.30ESZ," meters/second
704 Image5:IMAGE "Resonator Diameter" = "20.5D," +/- "D.5D," mm
706 Image6:IMAGE "Length (micr. zero)" = "30.5D," +/- "D.5D," mm
708 Image7:IMAGE "Sample Description: "K" = "20.4D," +/- "D.4D," mm
710 Image8:IMAGE "Initial guess for epsilon:" = "20.4D," +/- "D.4D," mm
712 Image9:IMAGE "Length" = "20.4D," +/- "D.4D," mm
714 Image10:IMAGE "20X, "with Sample" = "2(52D.4D)," +/- "D.4D, 6X)," mm
716 Image11:IMAGE "Micrometer reading" = "2(2D.7D)," +/- "D.7D, 3X)," GHz
718 Image12:IMAGE "Resonance frequency" = "2(5D)," +/- "D.5D, 11X)," GHz
720 Image13:IMAGE "q value" = "2(5D.2D)," +/- "D.5D, 20)," +/- "D.5D, 20, 6X)," B, "kHz"
722 Image15:IMAGE "Epsilon" = "30.4D," +/- "D.4D, 5X)," Epsilon" = "20.5D," +/- "D.5D"
724 Image16:IMAGE "tan delta" = "D.6D," +/- "D.6D"
726 SUBEND
728
730
732
734
736
738

```

SUB Calc_dimers
OPTION BASE 1

```

868 GOSUB Print_dims2
870 RETURN
872 I Print_dims2:1
874 PRINT "Filename$:";Filename$;
876 FIRST=1 TO Num_modes
880 FOR IF=1 TO Num_modes
882 IF Spectrum(1,IF) THEN
884 PRINT USING "#,K,DD";Modes;" Spectrum(1,1)
886 ELSE
888 PRINT USING "#,K,DD";" Spectrum(1,1)
890 END IF
892 NEXT I
894 PRINT USING "#,K,DD";" Spectrum(1,1)
896 END IF
898 NEXT I
900 PRINT
902 PRINT USING "#,2D,4D,K,D,4D,K,D,4D";" Diameter:";Diameter;" Diameter*1.E+3
=> *1.E+3 "+/-" ;Ldiameter*1.E+3, " Length:";Length*1.E+3, "+/-" ;Ulength*1.E+3
904 RETURN
906 I ///////////////////////////////////////////////////////////////////
908 Menu_calc_dimens:1
910 CLEAR SCREEN
912 PRINT USING Image1;Degc,Udegc,Relh,Ureth
914 PRINT USING Image2;Mbar,Ubar
916 PRINT USING Image3;C,Uc
918 PRINT USING Image4;Diameter*1.E+3,Udiameter*1.E+3
920 PRINT USING Image5;Length*1.E+3,Ulength*1.E+3
922 PRINT USING Image6;Micr(2)*1.E+3,Umicr(2)*1.E+3
924 PRINT USING Image7;Bessel_root
926 PRINT USING Image8;Xpos,Ypos
928 Men_calc_dimens:1
930 Xpos=0
932 Ypos=0
934 FOR I=1 TO Num_modes
936 Ypos=Ypos+1
938 PRINT TABXY(Xpos,Ypos);
940 IF Spectrum(1,5) THEN PRINT CHR$(129);
942 PRINT USING Image9c;Spectrum(1,1);Spectrum(1,2);1.E+9
944 IF I=9 THEN
946 PRINT CHR$(128);
948 Xpos=Xpos+35
950 Ypos=Ypos+8
952 PRINT TABXY(Xpos,Ypos);
954 PRINT USING Image8c
956 END IF
958 NEXT I
960 RETURN
962 I ///////////////////////////////////////////////////////////////////
964 Image1:IMAGE "Temperature = "2D.2D," +/- "D.D," deg C Relative Humidity
=> = "2D.D," +/- "D.D," %
966 Image2:IMAGE "Barometric Pressure = "3D.D," +/- "D.D," mbar"
970 Image3:IMAGE "Speed of Light = "D.7DESZ," +/- "D.3DESZ," meters/second
d"
972 Image4:IMAGE "Resonator Diameter Length (micr. zero) = "2D.5D," +/- "D.5D," mm"
974 Image5:IMAGE "Bessel function root: "2D.7D
976 Image6c:IMAGE "Axial mode Frequency (GHz)"
978 Image7c:IMAGE "X,2D,9X,2D,9D
980 Image8c:IMAGE "Micrometer reading = "2(S2D.4D," +/- "D.4D,6X),"mm"
982 Image11a:IMAGE "Micrometer reading = "2(S2D.4D," +/- "D.4D,6X),"mm"
984 I ///////////////////////////////////////////////////////////////////
986 I ///////////////////////////////////////////////////////////////////
990 Add_mode:OFF KEY
992 Num_modes=Num_modes+1

```

```

740 COM /Environment/ Degc,Mbar,Relh,C
742 COM /Environment/ Udegc,Ubar,Urelh,Uc
744 COM /Sample/ Sample_ids[55],SampleLen,Eps,re,Eps_im,Tand
746 COM /Sample/ UsampleLen,Udegs,re,Udegs_im,Utand,Eps_re,Udegs
748 COM /Sample/ Errors(12),ErrTand(12)
750 COM /Cavity_dims/ Diameter,Length,Spectrum(40,5),INTEGER Num_modes
752 COM /Cavity_dims/ Diameter,Ulength
754 COM /Cavity_dims/ Freq(2),Micr(2),Q(2),Bw(2)
756 COM /Cavity_vals/ Urreq(2),Umicr(2),Uq(2),Ubw(2)
758 COM /Constants/ Bessel_root
760 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
=> Prompts[180]
762 COM /Output/ INTEGER Print_addr,Plotter,Printer_on,Plotter_on
764 COM /Sys/ Sys_ids[10]
766 COM /Sys_ms/ Msi_ids[20]
768 COM /Files/ DiskDrives[20],Filename$[10]
770 INTEGER Xpos,Ypos,N
772 REAL CO
774 CO=2.99792456E+8 vacuum speed of light in meters/second
776 I
778 Diameter=temp*Diameter
780 Length=Temp*Length
782 Calc_dimens:OFF KEY
784 GOSUB Menu_calc_dimen
786 ns "CHR$(129);" this is the menu for resonator dimensions calculatio
=>
788 Prty=VAL(SYSTEM$(SYSTEM PRIORITY))+1
790 ON KEY 0 LABEL "Cancel Results",Prty GOSUB Cancel_results
792 ON KEY 1 LABEL "Keep Results",Prty CALL Prior_menu
794 ON KEY 2 LABEL "Add Del Mode",Prty GOSUB Add_mode
796 ON KEY 3 LABEL "Delete Mode"
798 ON KEY 4 LABEL "Chng Tog Mode",Prty GOSUB Change_mode
800 ON KEY 5 LABEL "Init Modes",Prty GOSUB Init_modes
802 ON KEY 6 LABEL "Root/Micr.",Prty GOSUB Change_root
804 ON KEY 7 LABEL "Disk Modes",Prty GOSUB Disk_modes
806 ON KEY 8 LABEL "Find Dimens",Prty GOSUB Find_dimens
808 ON KEY 9 LABEL "Done"
810 Done=0
812 Prior_menu=0
814 LOOP
816 IF Done THEN GOTO Calc_dimens
818 EXIT IF Prior_menu
820 END LOOP
822 OFF KEY
824 Prior_menu=0
826 Done=1
828 SUBEXIT
830 I ///////////////////////////////////////////////////////////////////
832 I ///////////////////////////////////////////////////////////////////
834 I ///////////////////////////////////////////////////////////////////
836 I ///////////////////////////////////////////////////////////////////
838 I ///////////////////////////////////////////////////////////////////
840 Cancel_results:OFF KEY
842 Diameter=Diameter,temp
844 Length=Length,temp
846 Prior_menu=1
848 RETURN
850 I ///////////////////////////////////////////////////////////////////
852 I ///////////////////////////////////////////////////////////////////
854 I ///////////////////////////////////////////////////////////////////
856 Print_dims:1
858 IF Printer ON THEN
860 GOSUB Print_dims2
862 PRINTER IS T
864 END IF

```

```

994 =>
996 INPUT "Enter the value of the axial mode number",Spectrum(Num_modes,1)
998 Spectrum(Num_modes,2)=Spectrum(Num_modes,2)*1.E+9
1000 INPUT "Enter the value of resonance Q",Spectrum(Num_modes,3)
1002 Spectrum(Num_modes,4)=-.05 *Resonance level
1004 INPUT "Will this mode be included in calculations (Y/N)",AS
1006 IF AS(1)="Y" OR AS(1)="y" THEN
1008 Spectrum(Num_modes,5)=1
1010 ELSE
1012 Spectrum(Num_modes,5)=0
1014 END IF
1016 Sort_modes(Spectrum(*), integer num_modes)
1018 MAT SORT Spectrum(*,1)
1020 Done=1
1022 RETURN
1024
1026 Change_mode:OFF KEY
1028 INPUT "Enter the value of the axial mode you wish to change",Edit_num
1030 J=1
1032 WHILE (Spectrum(J,1)<>Edit_num) AND (J<Num_modes)
1034 J=J+1
1036 END WHILE
1038 Test=Spectrum(J,1)
1040 OUTPUT 2 USING "k, #":Test
1042 INPUT "Enter the value of the axial mode number",Spectrum(J,1)
1044 Test=Spectrum(J,2)/1.E+9
1046 OUTPUT 2 USING "k, #":Test
1048 INPUT "Enter the value of the frequency (GHz)",Spectrum(J,2)
1050 Test=Spectrum(J,3)
1052 OUTPUT 2 USING "k, #":Test
1054 Spectrum(J,2)=Spectrum(J,2)*1.E+9
1056 Test=Spectrum(J,3)
1058 OUTPUT 2 USING "k, #":Test
1060 INPUT "Enter the resonance Q value",Spectrum(J,3)
1062 INPUT "Do you want to include this mode in calculations",AS
1064 IF AS(1)="Y" OR AS(1)="y" OR AS=" " THEN
1066 Spectrum(J,5)=1
1068 ELSE
1070 Spectrum(J,5)=0
1072 END IF
1074 MAT SORT Spectrum(*,1)
1076 CALL Sort_modes(Spectrum(*), integer num_modes)
1078 Done=1
1080 RETURN
1082
1084
1086 Toggle_mode:
1088 INPUT "Enter the mode number you wish to toggle",J
1090 FOR I=1 TO Num_modes
1092 IF Spectrum(I,1)=J THEN
1094 Spectrum(I,5)=0
1096 ELSE
1098 Spectrum(I,5)=1
1100 END IF
1102 END IF
1104 NEXT I
1106 Done=1
1108 RETURN
1110
1112
1114
1116
1118 Delete_mode:OFF KEY
1120 INPUT "Enter the value of the axial mode you wish to delete",Del_num
1122

```

```

1124
1126 FOR I=1 TO Num_modes
1128 IF Spectrum(I,1)=Del_num THEN
1130 Spectrum(I,1)=1000
1132 Spectrum(I,2)=0
1134 Spectrum(I,3)=0
1136 Spectrum(I,4)=0
1138 Spectrum(I,5)=0
1140 Hit=Hit+1
1142 END IF
1144 FOR K=J TO Num_modes
1146 Spectrum(K,1)=Spectrum(K+1,1)
1148 Spectrum(K,2)=Spectrum(K+1,2)
1150 NEXT K
1152 J=J-1
1154 Num_modes=Num_modes-1
1156 END IF
1158 NEXT J
1160 Num_modes=Num_modes-1
1162 MAT SORT Spectrum(*,1)
1164 Done=1
1166 RETURN
1168
1170 Init_modes:OFF KEY
1172 DISP "Do you really wish to initialize the spectrum modes?"
1174 Prty=VAL(SYSTEM("SYSTEM PRIORITY"))+1
1176 ON KEY 1 LABEL "yes" Prty GOSUB Continue
1178 ON KEY 2 LABEL "no", Prty CALL Prior_menu
1180 Done=0
1182 Prior_menu=0
1184 LOOP IF Done THEN Quit
1186 EXIT IF Prior_menu
1188 END LOOP
1190 OFF KEY
1192 Prior_menu=0
1194 Done=1
1196 RETURN
1198 Quit:
1200 RETURN
1202 Continue:MAT Spectrum= (0)
1204 FOR I=1 TO SIZE(Spectrum,1)
1206 Spectrum(I,1)=1000
1208 NEXT I
1210 Num_modes=0
1212 Done=1
1214 RETURN
1216
1218
1220 Change_root:OFF KEY
1222 DISP "Enter the value of the Bessel function root.:"
1224 INPUT Bessel_root
1226 PRINT TAB(10,7):
1228 PRINT USING "Image7:;Bessel_root"
1230 Done=1
1232 RETURN
1234
1236
1238
1240 Disk_modes:OFF KEY
1242 CALL Select_disk("File with frequency/Q data")
1244 IF Diskdrives="NO DISK" THEN GOTO Adm
1246 CALL Enterfilename("CAT", "File with frequency/Q data")
1248 IF filename="" THEN
1250 CAT Diskdrives
1252
1254

```

```

1256 CALL Enterfilename("ABORT", "File with frequency/0 data")
1257 END IF
1258 ALLOCATE A(40,2) (8(1,0))
1259 CALL Enter_fill(A(0), Num_modes, Ids)
1260 IF A(1,1) >= 8.23E+9 THEN
1261   IX=15
1262 ELSE
1263   IX=14
1264 END IF
1265 IC=0
1266 FOR I=1 TO Num_modes
1267   IF PROND(Spectrum(Ic,2),4)=PROND(A(1,1),4) THEN GOTO Skip_s
1268   pec_fill
1269   IC=IC+1
1270   Spectrum(Ic,1)=Ic+IX
1271   Spectrum(Ic,2)=A(1,1)
1272   Spectrum(Ic,3)=A(1,2)
1273   IF I-IX=27 THEN
1274     ELSE Spectrum(I,5)=0
1275     Spectrum(Ic,5)=I
1276     END IF
1277   END IF
1278   Num_modes=IC
1279   DEALLOCATE Ids,A(**)
1280   Done=1
1281   RETURN
1282   //////////////////////////////////////
1283   Find_dims:off key
1284   J=0
1285   FOR I=1 TO Num_modes
1286     IF Spectrum(I,5) THEN J=J+1
1287     NEXT I
1288     ALLOCATE Spec(1:J,2)
1289     FOR I=1 TO Num_modes
1290       IF Spectrum(I,5) THEN J=J+1
1291       IF Spectrum(I,5) THEN MAT Spec(J,*)=Spectrum(I,1:2)
1292     NEXT I
1293     CALL Find_dims(Spec(*),Length0,Diameter,Ulength,Udiameter,Micr(2),C,B
1294     => essel_root,J)
1295     DEALLOCATE Spec(*)
1296     PRINT TABXY(0,4)
1297     PRINT USING Image4,Diameter*1.E+3,Udiameter*1.E+3
1298     PRINT TABXY(0,5)
1299     PRINT USING Image5,Length0*1.E+3,Ulength*1.E+3
1300     PRINT "COVARIANCE MATRIX"
1301     PRINT Cvm(1,1),Cvm(1,2)
1302     PRINT Cvm(2,1),Cvm(2,2)
1303     Done=1
1304     RETURN
1305   //////////////////////////////////////
1306   Micr_reading:off key
1307   CALL Edit_data("the micrometer reading of the cavity (mm).",Micr(2),1
1308   => .E+3,Umicr(2))
1309   PRINT TABXY(0,6)
1310   PRINT USING Image11a,Micr(2)*1.E+3,Umicr(2)*1.E+3
1311   Done=1
1312
1368 RETURN
1369 SUBEND
1370 *****
1371 SUB Calc_c
1372 Equations from H. Liebe in "An Updated Model for Millimeter Wave
1373 Propagation in Moist Air
1374 RADIO SCIENCE, Vol. 20, Num 5, pp.1069-1089, Sept/Oct 1985
1375 *****
1376 Calc_c:off key
1377 COM /Environment/ Degc,Mbar,Relh,C
1378 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
1379 => Prompt$(80)
1380 GOSUB Menu_calc_c
1381 PRTY=VAL(SYSTEM$( "SYSTEM PRIORITY" ))+1
1382 ON KEY LABEL "Prior Menu",PrtY CALL Prior_menu
1383 ON KEY LABEL "Change Temp",PrtY CALL temp
1384 ON KEY LABEL "Change Rel Hum",PrtY CALL Rel_hum
1385 ON KEY LABEL "Change Pressure",PrtY CALL Pressure
1386 Prior_menu=0
1387 Done=0
1388 LOOP
1389 IF Done THEN GOTO Calc_c
1390 EXIT IF Prior_menu
1391 END LOOP
1392 Prior_menu=0
1393 Done=1
1394 SUBEXIT
1395
1396 Menu_calc_c:
1397 CLEAR SCREEN
1398 PRINT USING Image1,degc,Relh,Urelh
1399 PRINT USING Image2,Mbar,Umbar
1400 PRINT USING Image3,C,Uc
1401 DISP "Menu for speed of light calculation"
1402 RETURN
1403 Image:IMAGE "Temperature = ",2D,2D," +/- ",D,D," deg C Relative Humidit
1404 y = ",2D,D," +/- ",D,D," %"
1405 Image2:IMAGE "Barometric Pressure = ",3D,D," +/- ",D,D," mbar"
1406 Image3:IMAGE "Speed of Light = ",D,D,DESZ," +/- ",D,3DESZ," meters/seco
1407 nd"
1408 SUBEND
1409 *****
1410 SUB Calc_modes
1411 Calc_modes:off key
1412 OPTION BASE 1
1413 COM /Environment/ Degc,Mbar,Relh,C
1414 COM /Cavity_dims/ Diameter,Length0,Spectrum(40,5),INTEGER Num_modes
1415 COM /Cavity_dims/ Diameter,Length0,Ulength
1416 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
1417 => Prompt$(80)
1418 LOADSUB Modefreqs2 FROM "MODEFREQS2"
1419 CALL Modefreqs2(C,Diameter/2,Length0)
1420 DELSUB Modefreqs2
1421 Prior_menu=0
1422 Done=1
1423

```

```

1618 Image2:IMAGE "Barometric Pressure = ",D.70ESZ," +/- ",D.0," mbar"
1620 Images:IMAGE "Speed of Light = ",D.70ESZ," +/- ",D.30ESZ," meters/seco
rd"
1622 Image6:IMAGE "Resonator Diameter:
Length (micr. zero) = ",20.50," +/- ",D.50," mm"
1624 Image5:IMAGE "
Length (total) = ",30.50," +/- ",D.50," mm"
1626 Image5a:IMAGE "
Length (total) = ",30.50," +/- ",D.50," mm"
1628 Image6:IMAGE "Sample Description: ",K,"
1630 Image7:IMAGE "
Length = ",20.40," +/- ",D.40," mm"
1632 Image10ab:IMAGE 2BX,"Length",20X,"with Sample"
1634 Image11a:IMAGE "Micrometer reading = ",2(3D.40)," +/- ",D.40,DX),"mm"
1636 Image17:IMAGE "
Epsilon" = ",30.40," +/- ",D.40
1640 SUBEND
1642 ! *****
1644 !
1646 SUB Sample_descr
1648 Sample_descr:OFF KEY
1650 !
1652 OPTION BASE 1
1654 COM /Sample/ Sample_ids(5),SampleLen,Eps_re,Eps_im,land
1656 COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utand,Eps_re_guess
1658 COM /System_status/ Errrand(12),Errrand(12)
1660 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
=> Prompt$(80)
1662 LOADSUB Enter_id FROM "BAS.SUBS":1400,0,3"
1664 Enter_id(Sample_ids,"Sample")
1666 DELSUB Enter_id
1668 Done=1
1670 SUBEND
1672 ! *****
1674 !
1676 SUB Sample_thickns
1678 Sample_thickns:OFF KEY
1680 !
1682 OPTION BASE 1
1684 COM /Sample/ Sample_ids(5),SampleLen,Eps_re,Eps_im,land
1686 COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utand,Eps_re_guess
1688 COM /Sample/ Errrands(12),Errrand(12)
1690 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
=> Prompt$(80)
1692 !
1694 CALL Edit_data("the sample length (mm).",SampleLen,1.E+3,UsampleLen)
1696 Done=1
1700 SUBEND
1702 ! *****
1704 !
1706 SUB Resnce_freq
1708 Resnce_freq:OFF KEY
1710 !
1712 OPTION BASE 1
1714 COM /Cavity_vals/ Freq(2),Micr(2),q(2),Bw(2)
1716 COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
1718 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
=> Prompt$(80)
1720 !
1722 CALL Edit_data("the resonance frequency (GHz)",Freq(2),1.E-9)
1724 Done=1
1726 Ufreq(1)=Ufreq(2)
1728 IF Bw(1)>0 THEN
1730 q(1)=Freq(1)/Bw(1)
1732 Uq(1)=q(1)-Freq(1)/(Bw(1)+Ubw(1))
1734 !
1736 !
1738 !

```

```

1492 SUBEND
1494 ! *****
1496 !
1498 SUB Calc_freqs
1500 Calc_freqs:OFF KEY
1502 !
1504 OPTION BASE 1
1506 COM /Environment/ Degc,Mbar,Ureth,C
1508 COM /Environment/ Udegc,Umbar,Uureth,Uc
1510 COM /Sample/ Sample_ids(5),SampleLen,Eps_re,Eps_im,land
1512 COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utand,Eps_re_guess
1514 COM /Sample/ Errrands(12),Errrand(12)
1516 COM /Cavity_dims/ Diameter,Length,Ulength
1518 COM /Cavity_vals/ Ufreq(2),Micr(2),q(2),Bw(2)
1520 COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
1522 COM /Constants/ Bessel_root
1524 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
=> Prompt$(80)
1526 !
1528 GOSUB Menu_calc_freqs
1530 !
1532 Prty=VAL(SYSTEM$( "SYSTEM.PRIORITY" ))+1
1534 ON KEY 1 LABEL "Prior Menu",Prty CALL Prior_menu
1536 ON KEY 2 LABEL "Sample Thickns",Prty CALL Sample_thickns
1538 ON KEY 3 LABEL "Sample Epsilon",Prty CALL Epsilon
1540 ON KEY 4 LABEL "Cavity Diameter",Prty CALL Diameter
1542 ON KEY 5 LABEL "Cavity Length",Prty CALL Length
1544 ON KEY 6 LABEL "Cavity Reading",Prty CALL Micr_readings
1546 ON KEY 7 LABEL "Freq. Empty",Prty CALL Freq_empty
1548 ON KEY 8 LABEL "Freq. Loaded",Prty CALL Freq_loaded
1550 Prior_menu=0
1552 Done=0
1554 LOOP
1556 !
1558 IF Done THEN GOTO Calc_freqs
1560 EXIT IF Prior_menu
1562 END LOOP
1564 OFF KEY
1566 Prior_menu=0
1568 Done=1
1570 SUBEXIT
1572 ! *****
1574 !
1576 !
1578 !
1580 Menu_calc_freqs: !
1582 CLEAR SCREEN
1584 PRINT USING Image6: S_ids
1586 PRINT USING Image7: SampleLen*1.E+3,UsampleLen*1.E+3
1588 PRINT USING Image17: Eps_re,Ueps_re
1590 PRINT
1592 PRINT USING Image1: Degc,Udegc,Reh,Ureth
1594 PRINT USING Image2: Mbar,Umbar
1596 PRINT USING Image3: C,Uc
1598 PRINT USING Image4: Diameter*1.E+3,Udiameter*1.E+3
1600 PRINT USING Image5: Length*1.E+3,Ulength*1.E+3
1602 PRINT USING Image5a: (Length+Micr(1))*1.E+3,Ulength
1604 PRINT
1606 PRINT USING Image10ab
1608 PRINT USING Image11a: Micr(1)*1.E+3,Umicr(1)*1.E+3,Micr(2)*1.E+3,Umicr
(2)*1.E+3
1610 PRINT
1612 RETURN
1614 Image1:IMAGE "Temperature = ",20.20," +/- ",D.20," deg C
Relative Humidit
=> Y = ",20.D," +/- ",D.D," %

```

```

1740 END IF
1741 IF Bw(2)>0 THEN
1742   Q(2)=Freq(2)/Bw(2)
1743   Uq(2)=Q(2)-Freq(2)/(Bw(2)+Ubw(2))
1744   END IF
1745   Done=1
1746   SUBEND
1747 *****
1748 SUB Micr_readings
1749 Micr_readings:OFF KEY
1750 *****
1751 OPTION BASE 1
1752 COM /Cavity_vals/ Freq(2),Micr(2),Q(2),Bw(2)
1753 COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
1754 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
1755 Prompt$(80)
1756 *****
1757 Prompts="Do you want to change the EMPTY or LOADED cavity micrometer
1758 reading?"
1759 CALL Empty_or_Loaded
1760 IF Prior_menu THEN
1761   Prior_menu=0
1762   Done=1
1763   SUBEXIT
1764   END IF
1765 IF Empty_cavity THEN
1766   CALL Edit_data("the micrometer reading for the empty cavity (mm).")
1767 ELSE
1768   CALL Edit_data("the micrometer reading for the loaded cavity (mm)
1769   ")
1770   Micr(2),1.E-3,Umicr(2))
1771   END IF
1772   GOTO Micr_readings
1773   SUBEND
1774 *****
1775 SUB Q_values
1776 Q_values:OFF KEY
1777 *****
1778 OPTION BASE 1
1779 COM /Cavity_vals/ Freq(2),Micr(2),Q(2),Bw(2)
1780 COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
1781 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
1782 Prompt$(80)
1783 *****
1784 Prompts="Do you want to change the EMPTY or LOADED cavity Q-value?"
1785 CALL Empty_or_Loaded
1786 IF Prior_menu THEN
1787   Prior_menu=0
1788   Done=1
1789   SUBEXIT
1790   END IF
1791 IF Empty_cavity THEN
1792   CALL Edit_data("Q for the empty cavity.",Q(1),1,Uq(1))
1793 CALL Edit_data("Q for the empty cavity.",Q(1),1)
1794 Uq(1)=Q(1)/100
1795 Bw(1)=Bw(1)/Q(1)
1796 Ubw(1)=Bw(1)-Freq(1)/Q(1)
1797 ELSE
1798   CALL Edit_data("Q for the loaded cavity.",Q(2),1,Uq(2))
1799 CALL Edit_data("Q for the loaded cavity.",Q(2),1)
1800 Uq(2)=Q(2)/100
1801 *****
1802 SUB Resnce_freqs
1803 Resnce_freqs:OFF KEY
1804 *****
1805 Prompts="Do you want to change the EMPTY or LOADED cavity resonance f
1806 rency?"
1807 CALL Empty_or_Loaded
1808 IF Prior_menu THEN
1809   Prior_menu=0
1810   Done=1
1811   SUBEXIT
1812   END IF
1813 IF Empty_cavity THEN
1814   CALL Edit_data("the empty cavity resonance frequency (kHz).",freq
1815   (1),1.E-3)
1816 CALL Edit_data("the empty cavity resonance frequency (kHz).",freq
1817   (1),1.E-3)
1818 IF Bw(1)>0 THEN
1819   Q(1)=Freq(1)/Bw(1)
1820 Uq(1)=Q(1)-freq(1)/(Bw(1)+Ubw(1))
1821 ELSE
1822   CALL Edit_data("of the 3dB bandwidth for the empty cavity (kHz).")
1823   Bw(1),1.E-3,Ubw(1))
1824   Q(1)=Freq(1)/Bw(1)
1825 Uq(1)=ABS(Q(1)-freq(1))/(Bw(1)+Ubw(1))
1826 ELSE
1827   CALL Edit_data("of the 3dB bandwidth for the loaded cavity (kHz).")
1828   Bw(2),1.E-3,Ubw(2))
1829   Q(2)=Freq(2)/Bw(2)
1830 Uq(2)=ABS(Q(2)-freq(2))/(Bw(2)+Ubw(2))
1831 END IF
1832 GOTO Band_widths
1833 SUBEND
1834 *****
1835 SUB Resnce_freqs
1836 Resnce_freqs:OFF KEY
1837 *****
1838 Prompts="Do you want to change the EMPTY or LOADED cavity resonance f
1839 rency?"
1840 CALL Empty_or_Loaded
1841 IF Prior_menu THEN
1842   Prior_menu=0
1843   Done=1
1844   SUBEXIT
1845   END IF
1846 IF Empty_cavity THEN
1847   CALL Edit_data("the empty cavity resonance frequency (GHz).",freq
1848   (1),1.E-9)
1849 CALL Edit_data("the empty cavity resonance frequency (GHz).",freq
1850   (1),1.E-9)
1851 Ufreq(1),1.E-3)
1852 IF Bw(1)>0 THEN
1853   Q(1)=Freq(1)/Bw(1)
1854 *****

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2096      ))/(B1^2)
2098      =>
2100      Funct_b1=TAN(B1*SampleLen)/B1+TAN(B0*Y)/B0
2102      Funct_b1_p=(B1*SampleLen*(1/COS(B1*SampleLen)^2)-TAN(B1*SampleLen
2104      DelTab1=Funct_b1/Funct_b1_p
2106      IF ABS(DelTab1)<1.E-2 THEN DelTab1=DelTab1/10
2108      PRINT DelTab1
2110      UNTIL (I=100) OR (ABS(DelTab1)<1.E-4)
2112      Eps_re=(B1^2+K^2)/(B0^2+K^2)
2114      PRINT Eps_re
2116      PAUSE
2118      IF I=100 THEN Eps_re=0
2120      DEG
2122      SUBEND
2124      *****
2126      SUB Calc_tand
2130      Equations from E, Ni and U, Stumper "Permittivity Measurements Using a
2132      Frequency-tuned Microwave TE01 Cavity Resonator"
2134      IEEE PROCEEDINGS, Vol. 132, Pt. 1, No. 1, February 1985 p27-32.
2136      IEE
2138      Calc_tand:OFF KEY
2140      *****
2142      RAD
2144      OPTION BASE 1
2146      COM /Environment/ Degc,Mbar,Relh,C
2148      COM /Environment/ Udegc,Uambar,Urelh,Uc
2150      COM /Sample/ Sample_ids(55),SampleLen,Eps_re,Eps_im,Tand
2152      COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utand,Eps_re_guess
2154      COM /Sample/ Erreps(12),Errtand(12)
2156      COM /Cavity dims/ Diameter,length0,Spectrum(40,5),INTEGER Num_modes
2158      COM /Cavity dims/ Diameter,Ulength
2160      COM /Cavity_vals/ Freq(2),Micr(2),Q(2),Bw(2)
2162      COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
2164      COM /Constants/ Bessel_root
2166      COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2168      Prompts(80)
2170      IF Do_var_freq THEN !Variable frequency/ fixed length
2172      ELSE !L=Length0+Micr(1) !Variable length/ fixed frequency
2174      L0=Length0+Micr(2)
2176      END IF
2178      L=Length0+Micr(2)
2180      A=Diameter/2
2182      Fc=(Bessel_root*C)/(pi*Diameter)
2184      Q0=Fc/freq(2)
2186      B0=2*pi*Freq(2)*SQRT(1-Q0^2)/C
2188      B=2*pi*Freq(2)*SQRT(Eps_re-Q0^2)/C
2190      B=(B/80)^2
2192      X=B*Eps_re
2194      P=(SIN(X)^2)/(B*(COS(X)^2))
2196      C=(2*B*Eps_re*(Eps_re-Q0^2))/(1-Q0^2)*P)/(2*X-SIN(2*X))
2198      F=(40^2*(2*X*(B-1)-SIN(X)+(2*B*E-L-2*X)*P))/(2*X-SIN(2*X))
2200      E=Eps_re*(40^2+2*A*(1-Q0^2)/L0)
2202      D=(F+C)/E
2204      A=(2*X*Eps_re*(B-Eps_re)*SIN(2*X)+(2*B*E-L-2*X)*P)/(Eps_re*(2*X-SIN(2*
2206      => X)))
2208      Tand=(A*Bw(2)-0*Bw(1))/Freq(2)
2210      DEG
2212      SUBEND
2214      *****
2216      *****
2218      *****
2220      *****

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1970      Uq(1)=Q(1)-Freq(1)/(Bw(1)+Ubw(1))
1972      END IF
1974      ELSE
1976      CALL Edit_data("the loaded cavity resonance frequency (GHz).",Fre
1978      q(2),1.E-9)
1980      CALL Edit_data("the uncertainty in the loaded cavity res. freq. (
1982      => kHz).",Ufreq(2),1.E-3)
1984      IF Bw(2)>0 THEN
1986      Uq(2)=Freq(2)/Bw(2)
1988      Uq(2)=Uq(2)-Freq(2)/(Bw(2)+Ubw(2))
1990      END IF
1992      GOTO Resnce_freqs
1994      SUBEND
1996      *****
1998      *****
2000      SUB Micr_reading
2002      Micr_reading:OFF KEY
2004      *****
2006      OPTION BASE 1
2008      COM /Cavity_vals/ Freq(2),Micr(2),Q(2),Bw(2)
2010      COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
2012      COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2014      Prompts(80)
2016      *****
2018      CALL Edit_data("the micrometer reading (mm).",Micr(1),1.E+3,Umicr(1))
2020      Micr(2)=Umicr(1)
2022      Done=1
2024      SUBEND
2026      *****
2028      *****
2030      SUB Calc_eps_re
2032      *****
2034      *****
2036      Equations from F. Horner, T.A. Taylor, R. Dunsmair in "Resonance
2038      Methods of Dielectric Measurement at Centimetre Wavelengths"
2040      J. IEE, 1946, 93, Pt.111, pp.53-68
2042      Calc_eps_re:OFF KEY
2044      *****
2046      RAD
2048      OPTION BASE 1
2050      COM /Environment/ Degc,Mbar,Relh,C
2052      COM /Environment/ Udegc,Uambar,Urelh,Uc
2054      COM /Sample/ Sample_ids(55),SampleLen,Eps_re,Eps_im,Tand
2056      COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utand,Eps_re_guess
2058      COM /Sample/ Erreps(12),Errtand(12)
2060      COM /Cavity dims/ Diameter,length0,Spectrum(40,5),INTEGER Num_modes
2062      COM /Cavity dims/ Diameter,Ulength
2064      COM /Cavity_vals/ Freq(2),Micr(2),Q(2),Bw(2)
2066      COM /Cavity_vals/ Ufreq(2),Umicr(2),Uq(2),Ubw(2)
2068      COM /Constants/ Bessel_root
2070      *****
2072      Y=length0+Micr(2)-SampleLen
2074      Mu=(4*pi)*10^-7
2076      Eps_0=8.854188E-12
2078      Eps_0=Mu*(C/2)^2
2080      H=2*pi*Freq(2)
2082      K=Bessel_root/(Diameter/2)
2084      Eps_re=Eps_re_guess
2086      B0=SQRT(Mu^2*Mu*Eps_0*air-K^2)
2090      B1=SQRT(Mu^2*Mu*Eps_re*Eps_0-K^2)
2092      REPEAT
2094      I=I+1

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2728 => Prompts(80)
2729 ! CALL Edit_data("the relative epsilon of the sample.",Eps_re,1)
2730 Done=1
2731 SUBEND
2732 ! *****
2733 ! SUB Diameter
2734 Diameter:OFF KEY
2735 ! *****
2736 ! OPTION BASE 1
2737 COM /Cavity dims/ Diameter,Length0,Spectrum(40,5),INTEGER Num_modes
2738 COM /Cavity dims/ Diameter,Ulength
2739 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2740 => Prompts(80)
2741 ! CALL Edit_data("the diameter of the cavity (mm).",Diameter,1.E+3,Udia
2742 meter) Done=1
2743 SUBEND
2744 ! *****
2745 ! SUB Length
2746 Length:OFF KEY
2747 ! *****
2748 ! OPTION BASE 1
2749 COM /Cavity dims/ Diameter,Length0,Spectrum(40,5),INTEGER Num_modes
2750 COM /Cavity dims/ Diameter,Ulength
2751 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2752 => Prompts(80)
2753 ! CALL Edit_data("the length of the zeroed cavity (mm).",Length0,1.E+3,
2754 Ulength) Done=1
2755 SUBEND
2756 ! *****
2757 ! SUB Sort_modes(Array(*),N)
2758 IF N=1 THEN SUBEXIT
2759 FOR J=2 TO N
2760 Dummy1=Array(J,1)
2761 Dummy2=Array(J,2)
2762 FOR I=J-1 TO 1 STEP -1
2763 IF Array(I,1)<=Dummy1 THEN GOTO Mark
2764 Array(I+1,1)=Array(I,1)
2765 Array(I+1,2)=Array(I,2)
2766 NEXT I
2767 I=0
2768 Mark: I Array(I+1,1)=Dummy1
2769 Array(I+1,2)=Dummy2
2770 NEXT J
2771 SUBEND
2772 ! *****
2773 ! SUB Empty_or_loaded
2774 Empty_or_loaded:OFF KEY
2775 ! *****
2776 ! COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2777 => Prompts(80)
2778 ! DISP Prompts
2779
2848 Prty=VAL(SYSTEMS("SYSTEM_PRIORITY"))+1
2849 ON KEY 1 LABEL "PRIOR MENU" Prty CALL Prior_menu
2850 ON KEY 2 LABEL "EMPTY CAVITY" Prty GOSUB Emcyc
2851 ON KEY 3 LABEL "LOADED CAVITY" Prty GOSUB Loaded
2852 Done=0
2853 Prior_menu=0
2854 LOOP
2855 EXIT IF Prior_menu OR Done
2856 END LOOP
2857 SUBEXIT
2858 ! *****
2859 ! Empty_cavity=1
2860 Done=1
2861 RETURN
2862 ! *****
2863 ! Loaded:OFF KEY
2864 Empty_cavity=0
2865 Done=1
2866 RETURN
2867 SUBEND
2868 ! *****
2869 ! SUB Prior_menu
2870 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2871 => Prompts(80)
2872 Prior_menu=1
2873 SUBEND
2874 ! *****
2875 ! SUB Calc_errors(OPTIONAL INTEGER Dont_print)
2876 ! *****
2877 ! Calc_errors:OFF KEY
2878 ! *****
2879 ! OPTION BASE 1
2880 COM /Environment/ Degc,Mbar,Relh,C
2881 COM /Environment/ Udegc,Uubar,Ureth,Uc
2882 COM /Sample/ Sample_id$(55),SampleLen,Eps_re,Eps_im,Tand
2883 COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utand,Eps_re_guess
2884 COM /Sample/ Erreps(12),Errtand(12)
2885 COM /Cavity dims/ Diameter,Length0,Spectrum(40,5),INTEGER Num_modes
2886 COM /Cavity vals/ Freq(2),Micr(2),Uq(2),Ubw(2)
2887 COM /Cavity vals/ Freq(2),Umicr(2),Uq(2),Ubw(2)
2888 COM /Constants/ Bessel_root
2889 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
2890 => Prompts(80)
2891 Print_addr=9
2892 ! *****
2893 ! DIM Eps(12),Tandel(12)
2894 Diameter=Diameter*Udiameter
2895 CALL Calc_eps_re
2896 Eps(1)=Eps_re
2897 Diameter=Diameter-Udiameter
2898 Length0=Length0+Ulength
2899 CALL Calc_eps_im
2900 Eps(2)=Eps_re
2901 Length0=Length0-Ulength
2902

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2974 SampleLen=SampleLen+UsampleLen
2976 CALL Calc_eps_re
2978 Eps(3)=Eps_re
2980 SampleLen=SampleLen+UsampleLen
2982
2984 Degc=Degc+Udegc
2986 CALL Speed_of_light
2988 CALL Calc_eps_re
2990 Degc=Degc+Udegc
2992 Eps(4)=Eps_re
2994
2996 Relh=Relh+Urelh
2998 CALL Speed_of_light
3000 CALL Calc_eps_re
3002 Relh=Relh+Urelh
3004 Eps(5)=Eps_re
3006
3008 Mbar=Mbar+Umbar
3010 CALL Speed_of_light
3012 CALL Calc_eps_re
3014 Mbar=Mbar+Umbar
3016 Eps(6)=Eps_re
3018
3020 CALL Speed_of_light
3022 C=C+Uc
3024 C=C-Uc
3026 CALL Calc_eps_re
3028 Eps(7)=Eps_re
3030
3032 Freq(2)=Freq(2)+Ufreq(2)
3034 CALL Calc_eps_re
3036 Eps(8)=Eps_re
3038 Freq(2)=Freq(2)-Ufreq(2)
3040
3042 Micr(1)=Micr(1)+Umicr(1)
3044 CALL Calc_eps_re
3046 Eps(9)=Eps_re
3048 Micr(1)=Micr(1)-Umicr(1)
3050
3052 Micr(2)=Micr(2)+Umicr(2)
3054 CALL Calc_eps_re
3056 Eps(10)=Eps_re
3058 Micr(2)=Micr(2)-Umicr(2)
3060
3062 CALL Calc_eps_re
3064 Eps(11)=Eps_re
3066 Eps(12)=Eps_re
3068
3070 Diameter=Diameter+Udiameter
3072 CALL Calc_tand
3074 Diameter=Diameter-Udiameter
3076 Tand(1)=Tand
3078 Length=Length+Ulength
3080 CALL Calc_tand
3082 Length=Length-Ulength
3084 Tand(2)=Tand
3086
3088 SampleLen=SampleLen+UsampleLen
3090 CALL Calc_tand
3092 SampleLen=SampleLen+UsampleLen
3094 Tand(3)=Tand
3096
3098 Degc=Degc+Udegc
3100 CALL Speed_of_light
3102
3104 CALL Calc_tand
3106 Degc=Degc+Udegc
3108 Tand(4)=Tand
3110
3112 Relh=Relh+Urelh
3114 CALL Speed_of_light
3116 CALL Calc_tand
3118 Relh=Relh+Urelh
3120 Tand(5)=Tand
3122
3124 Mbar=Mbar+Umbar
3126 CALL Speed_of_light
3128 CALL Calc_tand
3130 Mbar=Mbar+Umbar
3132 Tand(6)=Tand
3134
3136 CALL Speed_of_light
3138 C=C+Uc
3140 CALL Calc_tand
3142 C=C-Uc
3144 Tand(7)=Tand
3146
3148 Freq(2)=Freq(2)+Ufreq(2)
3150 CALL Calc_tand
3152 Freq(2)=Freq(2)-Ufreq(2)
3154 Tand(8)=Tand
3156
3158 Micr(1)=Micr(1)+Umicr(1)
3160 CALL Calc_tand
3162 Micr(1)=Micr(1)-Umicr(1)
3164 Tand(9)=Tand
3166
3168 Micr(2)=Micr(2)+Umicr(2)
3170 CALL Calc_tand
3172 Micr(2)=Micr(2)-Umicr(2)
3174 Tand(10)=Tand
3176
3178 Bw(1)=Bw(1)+Ubw(1)
3180 CALL Calc_tand
3182 Bw(1)=Bw(1)-Ubw(1)
3184 Tand(11)=Tand
3186
3188 Bw(2)=Bw(2)+Ubw(2)
3190 CALL Calc_tand
3192 Bw(2)=Bw(2)-Ubw(2)
3194 Tand(12)=Tand
3196
3198 CALL Calc_tand
3200
3202 MAT Errreps=Eps
3204 MAT Errtand=Tand
3206 FOR J=1 TO 12
3208 Errreps(J)=Errreps(J)-Eps_re
3210 Errtand(J)=Errtand(J)-Tand
3212
3214 NEXT J
3216 Ueps_re=ABS(SORT(Errreps(1)-2*Errreps(2)-2*Errreps(3)-2*Errreps(7)-2*Errreps(8)
3218 Uatand=ABS(SORT(Errtand(1)-2*Errtand(2)-2*Errtand(3)-2*Errtand(7)-2*Errtand(8)
3220 Ueps_re=Ueps_re+Ueps_re
3222
3224 IF NPAR<1 THEN IDONT Print
3226 Calc_errs_menu: !
3228 CLEAR SCREEN
3230 PRINT "
3232 DELTA TAND"
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5362 References: J.D. Jackson "Classical Electrodynamics" 1962 page 255.
5364 M. Draper & H. Smith "Applied Regression Analysis, Second
5366 Edition" 1981 Chapter 2 and page 532.
5368 FOR I=1 TO Num_modes
5370 X(I,1)=1
5372 Y(I,1)=Spectrum(I,1)^2
5374 Y(I,1)=Spectrum(I,2)^2
5376 NEXT I
5378 MAT X%= TRN(X)
5380 MAT Y%= X*X
5382 MAT Xtx%= INV(X*X)
5384 MAT Yt%= TRN(Y)
5386 MAT Xty%= X*Y
5388 MAT Bt%= Xtx*Xty
5390 MAT Bt%= TRN(B)
5392 MAT Btxty%= Bt*Xty
5394 MAT S2%= IY-Btxty
5396 S2(1)=S2(1)/(N-2)
5400 MAT Cvm% Xtx*(S2(1))
5402 CALL T95percent(I, Num_modes-2)
5404 Length=C/(2*SQR(B(2,1))) Micr
5406 Diameter=Bessel_root*(SQR(B(1,1))) Micr
5408 Ulength=c*(1+SQR(Cvm(2,2)))/(2*B(2,1))*(1.5))
5410 Udiameater=Bessel_root*(1+SQR(Cvm(1,1)))/(2*B(1,1))*(1.5))
5412 DEALLOCATE X(*), Y(*), B(*), Bt(*), Xtx(*), Yty(*), Cvm(*)
5414 DEALLOCATE Xt(*), Yt(*), Xtx(*), S2(*), Yty(*), Btxty(*)
5416 SUBEND
5418 *****
5420 *****
5422 *****
5424 SUB Calc_err_anal
5426 Calc_err_anal:OFF KEY
5428
5430 OPTION BASE 1
5432 COM /Environment/ Degs, Mbar, Relh, C
5434 COM /Environment/ Udeg, Ubar, Urelh, Uc
5436 COM /Sample/ Sample_ids(151), SampleLen, Eps_re, Eps_im, Tand
5438 COM /Sample/ UsampleLen, Ueps_re, Ueps_im, Utand, Eps_re_guess
5440 COM /Sample/ Erreps(12), Errtand(12)
5442 COM /Cavity dims/ Diameter, Length, Spectrum(40,5), INTEGER Num_modes
5444 COM /Cavity dims/ Diameter, Ulength
5446 COM /Cavity vals/ Freq(2), Micr(2), Q(2), Bw(2)
5448 COM /Cavity vals/ Ufreq(2), Umicr(2), Uq(2), Ubw(2)
5450 COM /Constants/ Bessel_root
5452 COM /System_status/ INTEGER Done, Prior_menu, Do_var, freq, Empty_cavity,
=> Prompts(80)
5454 Equations: I
5456 F0=Freq(2)
5458 Fe=Freq(2)
5460 L0=Micr(1)
5462 Le=Micr(2)
5464 S=SampleLen
5466 D=Diameter
5468 Us=Bessel_root
5470 Eps=Eps_re
5472 YLe=S
5474 Fc=(U*Fc)/(PI*D)
5476 Xc=Fc/F0
5478 B0=(2*PI*F0*SQR(1-(X^2)))/C
5480 Be=(2*PI*Fe*SQR(Eps-(X^2)))/C
5482 Par_eps_b1=(X^2)
5484 Par_eps_q1=(-2*S*X)+(2*X)
5486 Par_b_be2*(1/(TAN(Be*S^2)))*B)/(TAN(B0*Y^2))
5488 Par_b_b0=2*(1-(TAN(Be*S^2)))/(1/(SIN(B0*Y^2)))*Y)/(TAN(B0*Y^2))
5490
5492 Par_b_b=2*(1-(TAN(Be*S^2))/(COS(Be*S^2)))*Be)/(TAN(B0*Y^2))
5494 Par_b_y=2*(1-(TAN(Be*S^2))/(SIN(B0*Y^2)))*B0)/(TAN(B0*Y^2))
5496 Par_be=fe*(2*PI*Fe*SQR(Eps_re-(X^2)))/C
5498 Par_b0=qc*(2*PI*Fe)/(C*SQR(Eps-(X^2)))
5500 Par_be=eps*(PI*Fe)/(C*SQR(Eps-(X^2)))
5502 Par_b0_f0=(2*PI*F0)/(C*SQR(1-(X^2)))/C
5504 Par_b0_qc=(2*PI*F0*SQR(1-(X^2)))/C
5506 Par_qc=1/F0
5508 Par_f0=f0-fc/(PI*D)
5510 Par_fc=c=U/(PI*D)
5512 Par_y=10=1
5514 Par_y=10=1
5516 Par_y=10=1
5518 Par_y=10=1
5520 Par_y=10=1
5522 Par_y=10=1
5524 Differential: I
5526 Del_s=UsampleLen
5528 Del_d=Udiameter
5530 Del_L=Umicr(2)
5532 Del_L0=Umicr(1)
5534 Del_f0=Ufreq(1)
5536 Del_fe=Ufreq(2)
5538 Del_c=Uc
5540 Del_y=(Par_y*Le*Del_le)+(Par_y*L0*Del_l0)+(Par_y*B*Del_s)
5542 Del_b0=(Par_b0_f0*Del_f0)+(Par_b0_c*Del_c)+(Par_b0_q*Del_q)
5544 Del_qc=(Par_qc*Del_fc)+(Par_qc*Del_f0)
5546 Del_be=(Par_be*Del_eps)+(Par_be*c*Del_c)+(Par_be*fe*Del_fe)+(Par_
=> be*Del_q)
5548
5550 L_y) Del_eps=(Par_eps_b*Del_b)+(Par_eps_q*Del_q)
5552 SUBEND
5554 *****
5556 *****
5558 *****
5560 *****
5562 *****
5564 *****
5566 *****
5568 *****
5570 *****
5572 *****
5574 *****
=> Prompts(80)
5576 DISP CHR$(129); " You may change Sample Information here. "; CHR$(128)
5578 Prty=VAL(SYSKEY$(SYSTEM PRIORITY));+1
5580 ON KEY 1 LABEL "Prior Menu", Prty CALL Prior_menu
5582 ON KEY 2 LABEL "Sample Desc", Prty CALL Sample_descr
5584 ON KEY 3 LABEL "Sample Thickness", Prty CALL Sample_thickns
5586 Done=0
5588 Prior_menu=0
5590 LOOP
5592 IF Done THEN GOTO Sample_info
5594 EXIT IF Prior_menu
5596 END LOOP
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5742 Ua(2)=SQRT(Sig2*Covar(2,2)) IUq1
5744 Ua(3)=SQRT(Sig2*Covar(3,3)) IUf0
5746 Ua(4)=SQRT(Sig2*Covar(4,4)) Uphase
5748 Ua(5)=Ua(4)*3.14159/2/PI i in degrees
5750 Ua(3)=SQRT(Sig2*Covar(5,5)) !Umeg (beta)
5752 CASE ELSE
5754 BEEP
5756 DISP "Sss=";Sss;" in Calc_uncert. PAUSED"
5758 PAUSE
5760 GOTO Calc_uncert
5762 END SELECT
5764 RETURN
5766 !
5768 !
5770 !
5772 Calc_sig2:
5774 FOR i=1 TO Fitcount
5776 GOSUB Gen_cnjum
5778 N=Start_index+1
5780 Sig2=Sig2+(ABS(CMPLX(S(N,1),S(N,2))-Cnum))^2
5782 NEXT i
5784 Sig2=Sig2/(Fitcount-Mfit(2))
5786 RETURN
5788 !
5790 !
5792 Print_results:
5794 PRINT "Printer IS 1
5796 PRINT "Descriptions: fit to "Sparms" values"
5798 PRINT USING "K,3D,2(3X,K,MD,3DE)";Iteration #,K,"Chi_squared=",Chisq
5800 q,"Alamda=",Alamda
5802 "Start_index=" Stop_index=" Stop_index
5804 GOSUB Print_a
5806 GOSUB Print_Ua
5808 PRINT USING "K,D,40E";"Chi_squared=",Chisq
5810 PRINT "Printer IS Print_addr
5812 GOSUB Print_a
5814 GOSUB Print_Ua
5816 PRINT USING "K,D,40E";"Chi_squared=",Chisq
5818 PRINT "Printer IS 1
5820 Reset_keys=1
5822 RETURN
5824 !
5826 !
5828 !
5830 Initial_plot:
5832 ALLOCATE File(Fitcount*3,2),Old_est(fitcount,2),Orig_sparm(Fitcount,2
5834 )
5836 CALL Orig_sparm(*,1:2)= S(Start_index:Stop_index,Ns,1:2)
5838 GOSUB Gen_estimate
5840 MAT Old_est=Sgen
5842 IF NOT Speed THEN GOSUB Plot_estimate
5844 RETURN
5846 !
5848 !
5850 Gen_cnjum:
5852 SELECT Sss
5854 CASE 1 IS21
5856 Cnum=-X(1)^2*CMPLX(A(1),-A(2))/CMPLX(X(1)^2-A(4)^2,X(1)^4/A(3
5858 )
5860 CASE 2 IS11
5862 Degrees=360*A(4)/2/PI
5864 P=CMPLX(COS(Degrees),SIN(Degrees))

```

```

5618 Sample_init_eps:OFF KEY
5620 OPTION BASE 1
5622 COM /Sample/ Sample_ids(55),SampleLen,Eps_re,Eps_im,Iand
5624 COM /Sample/ UsampleLen,Ueps_re,Ueps_im,Utamd,Eps_re_guess
5626 COM /System/ Errreps(12),Errrand(12)
5628 COM /System_status/ INTEGER Done,Prior_menu,Do_var_freq,Empty_cavity,
5630 Prompt$(80)
5632 !
5634 ! CALL Edit_data("the sample's initial guess for epsilon.",Eps_re_guess
5636 )
5638 SUBEND
5640 !
5642 !
5644 !
5646 SUB Fit_sparms
5648 ! Fit Lorentz(Kx,Ky,Q,F0,Ukx,Uky,Uq,Uf0) Sss=1
5650 ! or Fit Lorentz(Q,Q,F0,Phase,Uq,Uf0,Uphase) Sss=2 or 3
5652 ! Written by Eric J. VanZura
5654 ! Levenberg-Marquart algorithm CSUB taken from:
5656 ! Numerical Recipes: the art of Scientific Computing
5658 ! W. H. Press, B. P. Flannery, S. A. Teukolsky, W. T. Vetterling
5660 ! Cambridge Press 1987
5662 Fit_sparms:OFF KEY
5664 GOSUB Init
5666 GOSUB Init_est
5668 GOSUB initial_plot iPlot meas'd S-parms & initial estimate
5670 REPEAT
5672 Do_mrqt:
5674 GOSUB Printtestvals
5676 Ochisq=Chisq
5678 MAT Old_a=A
5680 DISP "performing Levenberg-Marquart Least-square-fit"
5682 CALL KEY
5684 OFF MRQIN(X*),Y(*),Sig(*),Mdata(*),A(*),Ma(*),Lista(*),Mfit(*)
5686 ,Covar(*),Alpha(*),Nca(*),Chisq,Alamda,Sss,Debug)
5688 GOSUB Convergencekeys
5690 MAT Old_est=Sgen
5692 GOSUB Gen_estimate
5694 IF Speed->1 THEN
5696 GOSUB Plot_estimate
5698 END IF
5700 GOSUB Ask_converge
5702 MRQ_out:OFF KEY
5704 UNTIL !tst=>!tst_limit OR Force_converge
5706 Alamda=0
5708 CALL MRQIN(X*),Y(*),Sig(*),Mdata(*),A(*),Ma(*),Lista(*),Mfit(*),Cov
5710 ar(*),Alpha(*),Nca(*),Chisq,Alamda,Sss,Debug)
5712 GOSUB Calc_sig2
5714 GOSUB Calc_uncert
5716 GOSUB Print_results
5718 GOTO Subend
5720 !
5722 Calc_uncert:
5724 SELECT Sss
5726 CASE 1
5728 Ua(1)=SQRT(Sig2*Covar(1,1)) IKx
5730 Ua(2)=SQRT(Sig2*Covar(2,2)) IKy
5732 Ua(3)=SQRT(Sig2*Covar(3,3)) I0
5734 Ua(4)=SQRT(Sig2*Covar(4,4)) IFO
5736 Ua(5)=SQRT(Sig2*Covar(5,5)) I
5738 CASE 2,3
5740 Ua(1)=SQRT(Sig2*Covar(1,1)) IUq0

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5866      Q12f=2*A(2)*X(1)-A(3))/A(3)
5868      Cnum=A(5)*CHPLX(A(2)/A(1)+Q12f*Q12f,Q12f*(1-A(2)/A(1)))**P/(1-Q12f
5870      *Q12f)
5872      CASE 3 IS21
5874      Degrees=360*A(4)/2/Pi
5876      P=CHPLX(COS(Degrees),SIN(Degrees))
5878      G0g1=1-A(1)/A(2)
5880      Dum=CHPLX(1,2)*X(1)-A(3))/A(2)
5882      Cnum=A(5)*G0g1**P/Dum
5884      CASE ELSE
5886      BEEP
5888      DISP "Ssss=";Ssss;" is not valid in Gen_cnum."
5890      PAUSE
5892      END SELECT
5894      RETURN
5896      !
5898      !
5900      Gen_estimate:
5902      DISP "Generating latest estimate values"
5904      MAT Sgen= (0)
5906      SELECT Sss
5908      CASE =1 !Nussenzweig Lorentzian
5910      FOR I=1 TO Fitcount
5912      GOSUB Gen_cnum
5914      Sgen(I,1)=REAL(Cnum)
5916      Sgen(I,2)=IMAG(Cnum)
5918      NEXT I
5920      CASE =2 !Schulzen s11
5922      FOR I=1 TO Fitcount
5924      GOSUB Gen_cnum
5926      Sgen(I,1)=REAL(Cnum)
5928      Sgen(I,2)=IMAG(Cnum)
5930      NEXT I
5932      CASE =3 !Schulzen s21
5934      FOR I=1 TO Fitcount
5936      GOSUB Gen_cnum
5938      Sgen(I,1)=REAL(Cnum)
5940      Sgen(I,2)=IMAG(Cnum)
5942      NEXT I
5944      DISP "Ssss=";Ssss;" is not valid in Gen_estimate."
5946      BEEP
5948      PAUSE
5950      GOTO Gen_estimate
5952      END SELECT
5954      RETURN
5956      !
5958      !
5960      !
5962      !
5964      Plot_estimate:
5966      Pack_data(File(*)
5968      Ids=Sparm$
5970      Pen=2
5972      Pack_data(File(*),Orig_sparm(*),Fitcount,Pen,Ids)
5974      Ids="Old estimate"
5976      Pen=4
5978      Pack_data(File(*),Old_est(*),Fitcount,Pen,Ids)
5980      Ids="Latest estimate"
5982      Pen=6
5984      Pack_data(File(*),Sgen(*),Fitcount,Pen,Ids)
5986      Pen=8
5988      Test$="Real Part"
5990      Prompt$="Imaginary Part"
5992
5994
5996      Ids="results from file:"&&Rawdatafiles
5998      Init_graphics(Test$,Prompts$,Ids)
6000      Axes="X"
6002      Autoscale(File(*),Axes)
6004      Refresh_labels
6006      Disp "S-parameters are: Red=measured data, Green=Old estimate, Blue=L
6008      atest Estimator"
6010      Plot all(file*)
6012      !
6014      !
6016      !
6018      Init:
6020      DEG
6022      OPTION BASE 1
6024      INTEGER I Prty, Itst, Itst Limit, Ns
6026      INTEGER Pen, Done, Fitcount
6028      COMPLEX Cnum, P, Dum
6030      REAL Debey
6032      DIM Test$(160), Prompts$(40), Ids$(40), Rawdatafiles$(10), Sparm$(3)
6034      COM /History/ Status$(1), Time_crgns$(8), Date_crgns$(11), Date_origns$(11)
6036      COM /Labels/ Labels$(30) 160), INTEGER Lbl_count, REAL Lbl_addr(30,6)
6038      COM /Data param/ INTEGER Data_count, Filesize, Curvcount, Roster(17,4)
6040      COM /Data param/ REAL Sym_size, Symbol$(17)$(2), Curve_ids$(17)$(40)
6042      COM /Data param/ REAL Xmin_data, Xmax_data
6044      COM /Data param/ REAL Ymin_data, Ymax_data
6046      COM /Graphics/ INTEGER Speed
6048      COM /Background/ Graphtypes$(12), Margins$(2)$(10), Papersizes$(11)
6050      COM /Background/ INTEGER Pen_speed, INTEGER Backgnd_pen, Auto time
6052      COM /Background/ INTEGER Auto_fill, REAL X_cross, Y_cross, X
6054      COM /Background/ Xgrid_ticks$(2), INTEGER Xmajor, Yminor
6056      COM /Background/ Ygrid_ticks$(4), INTEGER Xmajor, Yminor
6058      COM /Background/ REAL Xmin_graph, Xmax_graph, Ymin_graph, Ymax_graph
6060      COM /Axes_labels/ Print_xlabel$(3), Print_ylabel$(3)
6062      COM /Axes_labels/ Sig_digits$(2), REAL Xl_size, Yl_size
6064      COM /Windowspace/ REAL Xl_min, Xl_max, Yl_min, Yl_max; !max! graph edges UDU
6066      COM /Windowspace/ REAL Xl_left, Xl_right, Yl_bottom, Yl_top; !paper edges UDU$
6068      COM /Log scale/ REAL Xcycles, Xbegin, Ycycles, Ybegin
6070      COM /Plot device/ Plot_lang$(10), INTEGER Plot_addr
6072      COM /Hard space/ REAL View_scale, Aspect_ratio
6074      COM /Frame_onoff/ INTEGER Frame_flag
6076      COM /Clear_space/ REAL Space_lft, Space_rt, Space_btm, Space_top
6078      !
6080      !
6082      Init_com:
6084      COM /Fittype/ INTEGER Sselect, Start_index, Stop_index
6086      COM /S_array/ INTEGER Dcount, Svalid, REAL S(801,5,2)
6088      COM /Cdata/ INTEGER Mfit(2), lista(5,2), Ma(2), REAL A(5), Ua(5), INTEGER
6090      Ndata(2)
6092      COM /Cstats/ Alambda, Chisq, INTEGER Nca(2), REAL Alpha(5,5), Covar(5,5)
6094      COM /Output/ INTEGER Print_addr, Plotter, Printer_on, Plotter_on
6096      ASSIGN "alp-inter" TO Print_addr
6098      Init_vals:
6099      !
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6252 Mfit(1)=0 'number of unknowns to be fitted
6253 Mfit(2)=M
6254 IF SSS=1 THEN Mfit(2)=M-1
6255 MAT Lista= (0)
6256 LISTA(1,2)=1 'A array index of the unknowns to be fitted
6260 LISTA(2,2)=2
6262 LISTA(3,2)=3
6264 LISTA(4,2)=4
6266 LISTA(5,2)=5
6268 LISTA(6,2)=0
6270 IF SSS=1 THEN LISTA(5,2)=0
6272 MAT Covar= (0) 'Size of Alpha & Covar arrays >= Ma
6274 MAT ALpha= (0)
6276 Nca(1)=0
6278 Nca(2)=M
6280 ALambda=-2
6282 ALLOCATE Old a(M)
6284 MAT Sij= (1)
6286 ALLOCATE REAL X(Fitcount*2),Y(Fitcount*2),Sig(Fitcount*2)
6288 MAT X(:,Fitcount)= S(Start_index:Stop_index,5,1)
6290 MAT Y(:,Fitcount)= S(Start_index:Stop_index,5,1)
6292 MAT Sig(:,Fitcount)= S(Start_index:Stop_index,5,1)
6294 RETURN 'Init_est
6296
6298
6299
6300
6302 Done_return=OFF KEY
6304 Done=1
6306 RETURN
6308
6310
6312
6314
6316 Force_converge=PRINT "Convergence will be forced."
6318 Force_force_converge=1
6320 RETURN
6322
6324
6326
6328 Preventconverge=PRINT "Convergence will be prevented."
6330 list=0
6332 RETURN
6334
6336
6338
6340 Graph_on: i
6342 Speed=1
6344 OFF KEY 4
6346 ON KEY 4 LABEL "GRAPH OFF",Prty GOSUB Graph_off
6348 RETURN
6350
6352
6354
6356 Graph_off: i
6358 Speed=0
6360 OFF KEY 4
6362 ON KEY 4 LABEL "GRAPH ON",Prty GOSUB Graph_on
6364 RETURN
6366 Debug_on:PRINT "debug flag is turred on."
6368 Debug=1
6370 OFF KEY 6
6372 ON KEY 6 LABEL "DEBUG OFF",Prty GOSUB Debug_off
6374 RETURN
6376
6378
6380
6382 Debug_off:PRINT "Debug flag is turned off."

```

```

6124 ta?) PAUSED"
6126 PRINT "The S(*) array for ",Sparms," is not valid (doesn't have d
6128 atar?) PAUSED"
6130 PAUSE
6132 END IF
6134 SSS=1 S21=Nussenzeivig SSS=2 S11=Schulten SSS=3 S21=Schulten
6136 SELECT Sselect
6138 CASE 0 IS11
6140 Sparms="S11"
6142 CASE 1 IS21
6144 SSS=1
6146 SSS=3
6148 Sparms="S21"
6150 CASE 2 IS12
6152 SSS=1
6154 Sparms="S12"
6156 CASE 3 IS22
6158 SSS=2
6160 Sparms="S22"
6162 CASE ELSE
6164 BEEP
6166 DISP "Sselect is not valid in Fit_sparms"
6168 PAUSE
6170 END SELECT
6172 M=Sselect+1
6174 RETURN
6176
6178
6180
6182 Init_est: i
6184 Fitcount=Stop_index-Start_index+1
6186 Start_freq=S(Start_index,5,1)
6188 Stop_freq=S(Stop_index,5,1)
6190 SELECT Sss
6192 CASE =1 i Nussenzeivig Lorentzian
6194 ix=INT(Start_index+Fitcount/2)
6196 Angle=ARG(CMPLX(S(1x,Ms,1),S(1x,Ms,2)))
6198 Mag=ABS(CMPLX(S(1x,Ms,1),S(1x,Ms,2)))
6200 A(1)=2*10^(-5)*Mag*SIN(Angle)
6202 A(2)=2*10^(-5)*Mag*COS(Angle)
6204 A(3)=1.5*S(1x,5,1)/(Stop_freq-Start_freq)
6206 A(4)=S(1x,5,1)
6208 CASE =2,=3 i Schulten S11 or S21
6210 ix=INT(Start_index+Fitcount/2)
6212 A(1)=S(1x,5,1)/(Stop_freq-Start_freq)*1.8
6214 A(2)=S(1x,5,1)/(Stop_freq-Start_freq)*2.0
6216 A(1)=S(1x,5,1)/90000
6218 A(2)=S(1x,5,1)/79000
6220 A(3)=S(1x,5,1)
6222 A(4)=ARG(CMPLX(S(1x,Ms,1),S(1x,Ms,2)))*2*PI/160
6224 A(4)=0
6226 A(5)=ABS(CMPLX(S(1x,Ms,1),S(1x,Ms,2)))
6228 A(5)=.7
6230 CASE ELSE
6232 DISP "SSS=",SSS," is not valid in Init_est"
6234 REEP
6236 PAUSE
6238 GOTO Init_est
6240 END SELECT
6242 Mdata(1)=0
6244 Mdata(2)=Fitcount*2
6246 Ma(1)=0 'size of A array (unknowns)
6248 Ma(2)=M 'measure of fit
6250 Chisq=1

```

```

6506 => " ,A(3)/A(1), "Q1=" ,A(3)/A(2), "F0=" ,A(3), "degrees=" ,A(4)*360/2/PI, "Mag=" ,A(5
=> )
CASE ELSE
6508 DISP "Sss=" ,Sss, " is not valid in print_a. program paused"
6510 -BEOP
6512 PAUSE
6514 END SELECT
6516 RETURN
6518
6520 !
6522 !
6524 !
6526 Print_uj:
6528 -SELECT Sss
6530 CASE =1
6532 PRINT USING "2(K,MD,3DE,4X),K,M5D,3X,K,3X,MD,3DE",ujkx=" ,Ua(1), "U
=> ky=" ,Ua(2), "Uj=" ,Ua(3), "Uf0=" ,Ua(4)
6534 CASE =2, =3
6536 PRINT USING "2(K,MD,3DE,2X),K,MD,9DE,2X,K,MD,3D,3D,2X,K,MD,3DE", "UQ
=> O=" ,Ua(1), "UQ1=" ,Ua(2), "UQ2=" ,Ua(3), "Udegrees=" ,Ua(4), "Umag=" ,Ua(5)
6538 CASE ELSE
6540 DISP "Sss=" ,Sss, " is not valid in Print_uj. program paused"
6542 -BEOP
6544 PAUSE
6546 END SELECT
6548 RETURN
6550 !
6552 !
6554 !
6556 Subend:
6558 SUBEND ifit_sparms
6560 *****
6562 *****
6564 *****
6566 *****
6568 *****
6570 *****
6572 *****
6574 *****
6576 *****
6578 *****
6580 *****
6582 *****
6584 *****
6586 *****
6588 *****
6590 *****
6592 *****
6594 *****
6596 *****
6598 *****
6600 *****
6602 *****
6604 *****
6606 *****
6608 *****
6610 *****
6612 *****
6614 *****
6616 *****
6618 *****
6620 *****
6622 *****
6624 *****

```

```

6384 Debug=0
6386 OFF KEY 6
6388 ON KEY 6 LABEL "DEBUG ON",Prty GOSUB Debug_on
6390 RETURN
6392 !
6394 !
6396 !
6398 Ask_converge:
6400 SELECT Chisq
6402 CASE >Ochisq
6404 Itst=0
6406 PRINT USING "K,D,DE","New Chisq is greater than old Chisq by: ";Ch
=> isq-Ochisq
6408 CASE =Ochisq
6410 PRINT "Chisq is equal to old Chisq"
6412 CASE <Ochisq
6414 PRINT USING "K,D,DE","New Chisq is LESS than old Chisq by: ";Ochis
=> q-Chisq
6416 IF ABS((Ochisq-Chisq)/Chisq)<.003 THEN Itst=Itst+1
6418 IF ABS((Ochisq-Chisq)/Chisq)<.00003 THEN Itst=Itst+1
6420 PRINT USING "K,D,DE","Itst=" ,Itst
6422 PRINT USING "K,D,DE","ABS((Ochisq-Chisq)/Chisq)=" ,ABS((Ochisq-Chi
=> sq)/Chisq)
6424 END SELECT
6426 RETURN
6428 !
6430 !
6432 !
6434 Convergencekeys:
6436 PRYVAL(SYSTEM,PRIORITY)*1
6438 ON KEY 0 LABEL "FORCE CONVERG",Prty GOSUB force_converge
6440 ON KEY 2 LABEL "PREVENTCONVERG",Prty GOSUB Preventconverge
6442 IF Speed THEN
6444 ON KEY 4 LABEL "GRAPH OFF",Prty GOSUB Graph_off
6446 ON KEY 4 LABEL "GRAPH ON",Prty GOSUB Graph_on
6448 END IF
6450 IF Debug THEN
6452 ON KEY 6 LABEL "DEBUG OFF",Prty GOSUB Debug_off
6454 ON KEY 6 LABEL "DEBUG ON",Prty GOSUB Debug_on
6456 END IF
6458 ON KEY 6 LABEL "DEBUG ON",Prty GOSUB Debug_on
6460 END IF
6462 RETURN
6464 !
6466 !
6468 !
6470 Printlatestvals:
6472 PRINT USING "K,MD,3DE","Chisq=" ,Chisq
6474 Kx=1
6476 PRINT USING "K,3D,L,K,MD,3DE","Iteration #",K, "Alamda=" ,Alamca
6478 GOSUB Print_a
6480 RETURN
6482 !
6484 !
6486 !
6488 Print_a:
6490 -SELECT Sss
6492 CASE =1
6494 IF A(4)<1.E+10 THEN
6496 PRINT USING "2(K,MD,4DE,3X),K,M6D,3X,K,X,K,MD,9DE","Kx =" ,A(1),
=> "Ky =" ,A(2), "Q =" ,A(3), "F0 =" ,A(4)
6498 ELSE
6498 PRINT USING "2(K,MD,4DE,3X),K,M6D,3X,K,M2D,9DE","Kx =" ,A(1), "
=> Ky =" ,A(2), "Q =" ,A(3), "F0 =" ,A(4)
6500 CASE =2, =3
6502 END IF
6504 WAIT 1.8

```

```

6626      GOTO Enterfilename
6627      CASE "ABORT","EXIT"
6628          GOTO Abortline
6629      CASE ELSE
6630          DISP "Ac$=";Ac$;" in SUB Enterfilename"
6631          BEEP
6632          PAUSE
6633      END SELECT
6634      IF LEN(Tests$)>10 THEN
6635          BEEP
6636          DISP "ERROR in NAME ENTRY--up to 10 chars, you have "
6637          DISP LEN(Tests$);" "
6638          WAIT 1.8
6639          OUTPUT 2 USING "#,K,K";" #";Tests$
6640          GOTO Enterfilename
6641      END IF
6642      FOR I=1 TO LEN(FileName$)
6643          Ascii=num(NUM(FileName$(I)))
6644          SELECT Ascii,num
6645          CASE 65 TO 90,95,97 TO 122,48 TO 57
6646              !allowed characters
6647          CASE ELSE
6648              BEEP
6649              DISP "ERROR in NAME ENTRY--ILLEGAL CHARACTERS, TRY AGAIN."
6650              WAIT 1.8
6651              OUTPUT 2 USING "#,K,K";" #";FileName$
6652              GOTO Enterfilename
6653          END SELECT
6654      NEXT I
6655      SUBEXIT
6656      Abortline:FileName$=""
6657      SUBEND
6658      ! *****
6659      SUB Select disk(OPTIONAL Prompts)
6660          ! Original: 13 Nov 1984
6661          ! Revision: 02 Dec 1987
6662          COM /Files/ diskdrives(20),FileNames(10)
6663          DIM Disc$(30),Title$(40),Disp$(80)
6664          Local Prty=VAL(SYSTEM$(SYSTEM PRIORITY))+1
6665          Sys_id$=SYSTEM$(SYSTEM ID)
6666          OFF KEY
6667          ! Define the disk drives available for this system, reserve the
6668          ! first characters for the drive address and the characters after
6669          ! the - for a description of the drive.
6670          ! Example:
6671          ! Disc$(1)=":,700,0,0    HP 9133H HARD disk, volume 0."
6672          ! If NPAR>0 THEN
6673              Disp$=" SELECT DISK DRIVE for %gPrompts$" ... Abort will cancel"
6674              Disp$=Disp$(1,80)
6675          ELSE
6676              Disp$=" SELECT DISK DRIVE ... Abort will cancel. "
6677          END IF
6678          Title$=" Available disk drives for this system. "
6679          Prty=1 ! allow only one select
6680          IF Diskdrives$(1,1)<>" THEN Diskdrives$=""

```

```

6758      !
6759      !
6760      Msi_ids=SYSTEM$(Msi_id)
6761      Msi_id$=Msi_ids(POS(Msi_id$,";"),LEN(Msi_id$))
6762      Diskdrives=TRIMS(Diskdrives)
6763      Msi_ids=TRIMS(Msi_ids)
6764      IF LEN(Diskdrives)>0 AND LEN(Msi_ids)>0 THEN
6765          Disc$(1)=Diskdrives&RPTS(" ",17-LEN(Diskdrives))
6766          Disc$(1)=Disc$(1)&"- Last selected disk drive."
6767          Dd=1
6768          IF Diskdrives<Msi_ids THEN
6769              Disc$(2)=Msi_ids&RPTS(" ",17-LEN(Msi_ids))
6770              Disc$(2)=Disc$(2)&"- Start-up mass storage unit specifier."
6771          ELSE
6772              Disc$(1)=Disc$(1)&" Start-up MSUS."
6773          END IF
6774      ELSE
6775          IF LEN(Msi_ids)>0 THEN
6776              Disc$(1)=Msi_ids&RPTS(" ",17-LEN(Msi_ids))
6777              Disc$(1)=Disc$(1)&"- Start-up mass storage unit specifier."
6778          ELSE
6779              Dd=1
6780              Disc$(1)=Disc$(1)&" Start-up MSUS."
6781          END IF
6782      END IF
6783      ! ***** customize system drives here *****
6784      ! Follow format with - after unit specifier, description is
6785      ! optional but recommended.
6786      ! *****
6787      SELECT TRIMS(Sys_id$)
6788      CASE "$S300;20"
6789          Disc$(Dd+1)=":,1404
6790          Disc$(Dd+2)=":,1400,1
6791          Disc$(Dd+3)=":,1400,0
6792          Disc$(Dd+4)=":,1400,0,3
6793          Disc$(Dd+5)=":,700,0,0
6794          Disc$(Dd+6)=":,700,0,1
6795          ! Dd=Dd+3! add the number of drive specifiers above
6796      CASE "PC100"
6797          Disc$(Dd+1)=":,1500,0
6798          Disc$(Dd+2)=":,1500,1
6799          Disc$(Dd+3)=":,1500,2
6800          Disc$(Dd+4)=":,1500,3
6801          Disc$(Dd+5)=":,1500,4
6802          Disc$(Dd+6)=":,DOS,A
6803          Disc$(Dd+7)=":,DOS,B
6804          Disc$(Dd+8)=":,DOS,C
6805          ! Dd=Dd+8
6806          CASE ELSE
6807              BEEP
6808          DISP "you need to define your diskdrives in Select_disk for:.",Sys
6809          PAUSE
6810          END SELECT
6811      ! *****
6812      CALL Menu_scroll(Disp$,Title$,Disc$(*),Dd,Pt,Choose(*))
6813      IF Prty=0 THEN
6814          Diskdrives="NO DISK"
6815      ELSE
6816          Dd=POS(Disc$(Choose(Pt)),",")-1 ! find -

```

```

6886 IF Dd>5 THEN I valid msus
6888   Diskdrive$=TRIMS(Disc$(Choose(Pt.))(1,Dd))
6890 ELSE DISP " ERROR in reading MSUS from string, - chr not found. "
6892 BEEP
6894 CALL Pause_key_on
6896 Diskdrive$="NO DISK"
6898 IF
6900 END IF
6902 Diskselected:OFF KEY
6904 SUBEND
6906 SUBEXIT
6908 I *****
6910 I *****
6912 I *****
6914 I *****
6916 SUB Enter_file(T_f*),INTEGER Datacount,ids)
6920 OPTION BASE 1
6922 COM /Files/ Diskdrives$[20],filenames$[10]
6924 ON ERROR CALL Errortrap
6926 ASSIGN adatepath TO Filenames$Diskdrives$
6928 ENTER adatepath;Status$
6930 ENTER adatepath;ids
6932 ENTER adatepath;Datacount
6934 ENTER adatepath;Datacount
6936 ALLOCATE Temp(Datacount,2)
6938 ENTER adatepath;Temp(*)
6940 ASSIGN adatepath TO *
6942 IF SIZE(T_f,1)>=Datacount THEN
6944   FOR I=1 TO Datacount
6946     T_f(I,1)=Temp(I,1)
6948     T_f(I,2)=Temp(I,2)
6950   NEXT I
6952 ELSE DISP "SIZE OF FILE PASSED TO Enter_file=";SIZE(T_f,1);
6954   DISP " IS TOO SMALL. PROCESS ABORTED. "
6956 BEEP
6958 WAIT 1.8
6960 Datacount=0
6962 ICS=" "
6964 END IF
6966 DEALLOCATE Temp(*)
6970 OFF ERROR
6972 SUBEND
6974 I *****
6976 I *****
6978 I *****
6980 SUB Menu_scroll(0$,I$,Items$(*),INTEGER Item_cnt,To_select,Choose(*))
6982 Menu_scroll(: Original: 22 Jun 1987 Galen Koepke, NBS 723.04
6984   ! Revision: 02 Dec 1987
6986   ! A general purpose menu utility for scrolling items and
6988   ! selecting a given number of them.
6990   ! The items are arranged in screens of 15 items each and
6992   ! the user may access screens via softkeys. There may be
6994   ! up to 40 screens or 600 items to choose from.
6996   ! Items$(*) contains the item descriptions
6998   ! Item_cnt is the number of items in Items$(*)
7000   ! Choose(*) is dimensioned to the number of required choices
7002   ! and will be filled with the item numbers chosen.
7004   ! To_select is the number of required choices.
7006 I *****
7008 I *****
7010 OPTION BASE 1
7012 PRINT IS CRT
7014 DEG
7016 GOSUB Def_variables
7018
7020 GOSUB Define_screens
7022 IF Null_file THEN ! reset to zero
7024   Item_cnt=0
7026   Items$(1)=" "
7028   To_select=0 ! no valid selections
7030 END IF
7032 SUBEXIT
7034 I *****
7036 I *****
7038 I *****
7040 Def_variables:
7042 COM /Interrupts/ INTEGER Intr_prty
7044 COM /Bugs/ INTEGER Bug1, Bug2, Bug3, printer
7046 COM /Sys/ Sys_ids[10]
7048 I *****
7050 INTEGER Screen_cnt; Items_per_scn; First_item(40); Last_item(40)
7052 INTEGER Local_prty; J,K,First_line,Last_line,Active_screen,Pointer,Last_pt
7054 INTEGER Local_prty; Skips,Knobcount,PointerActive,K0,Null_file
7056 INTEGER Exit_Flag,temp
7058 DIM Markers[16],Tests[160]
7060 I *****
7062 I *****
7064 I *****
7066 I *****
7068 I *****
7070 I *****
7072 I *****
7074 Null_file=1
7076 Item_cnt=1
7078 To_select=0
7080 Items$(1)="**** Empty ****"
7082 I *****
7084 I *****
7086 I *****
7088 I *****
7090 I *****
7092 I *****
7094 I *****
7096 I *****
7098 I *****
7100 I *****
7102 I *****
7104 I *****
7106 Define_screens: ! Set up screens of 15 items each.
7108 I *****
7110 Items_per_scn=15 ! Maximum number of displayable items
7112 IF !N(Items_cnt/Items_per_scn)=Item_cnt/Items_per_scn THEN
7114   Screen_cnt=INT((Item_cnt/Items_per_scn))
7116 ELSE Screen_cnt=INT((Item_cnt/Items_per_scn))+1
7118 END IF
7120 FOR J=1 TO Screen_cnt ! set up each screen
7122   First_item(J)=J
7124   IF J+Items_per_scn-1<Item_cnt THEN
7126     Last_item(J)=J+Items_per_scn-1
7128   ELSE
7130     Last_item(J)=Item_cnt
7132   END IF
7134   J=J+Items_per_scn
7136 NEXT J
7138 RETURN
7140 I *****
7142 I *****
7144 I *****
7146 I *****
7018 GOSUB Define_screens
7020 GOSUB Make_selections
7022 IF Null_file THEN ! reset to zero
7024   Item_cnt=0
7026   Items$(1)=" "
7028   To_select=0 ! no valid selections
7030 END IF
7032 SUBEXIT
7034 I *****
7036 I *****
7038 I *****
7040 Def_variables:
7042 COM /Interrupts/ INTEGER Intr_prty
7044 COM /Bugs/ INTEGER Bug1, Bug2, Bug3, printer
7046 COM /Sys/ Sys_ids[10]
7048 I *****
7050 INTEGER Screen_cnt; Items_per_scn; First_item(40); Last_item(40)
7052 INTEGER Local_prty; J,K,First_line,Last_line,Active_screen,Pointer,Last_pt
7054 INTEGER Local_prty; Skips,Knobcount,PointerActive,K0,Null_file
7056 INTEGER Exit_Flag,temp
7058 DIM Markers[16],Tests[160]
7060 I *****
7062 I *****
7064 I *****
7066 I *****
7068 I *****
7070 I *****
7072 I *****
7074 Null_file=1
7076 Item_cnt=1
7078 To_select=0
7080 Items$(1)="**** Empty ****"
7082 I *****
7084 I *****
7086 I *****
7088 I *****
7090 I *****
7092 I *****
7094 I *****
7096 I *****
7098 I *****
7100 I *****
7102 I *****
7104 I *****
7106 Define_screens: ! Set up screens of 15 items each.
7108 I *****
7110 Items_per_scn=15 ! Maximum number of displayable items
7112 IF !N(Items_cnt/Items_per_scn)=Item_cnt/Items_per_scn THEN
7114   Screen_cnt=INT((Item_cnt/Items_per_scn))
7116 ELSE Screen_cnt=INT((Item_cnt/Items_per_scn))+1
7118 END IF
7120 FOR J=1 TO Screen_cnt ! set up each screen
7122   First_item(J)=J
7124   IF J+Items_per_scn-1<Item_cnt THEN
7126     Last_item(J)=J+Items_per_scn-1
7128   ELSE
7130     Last_item(J)=Item_cnt
7132   END IF
7134   J=J+Items_per_scn
7136 NEXT J
7138 RETURN
7140 I *****
7142 I *****
7144 I *****
7146 I *****

```

```

7150 Make_selections: MENU setup and use.
7152 Active_screen=1 ! first screen is active
7154 First_line=2 ! first printed line on screen = 2 or greater.
7156 GOSUB Write_screen ! activate screen at Active_screen
7158 ! and set First_line and Last_line for Pointer
7160 ! Write Marker$ to first non-selected line.
7162 ! Keys start at zero
7164 Exit_flag=0
7166 key_loop: ! allow ENTER key to exit when selections filled.
7168 ON KBD,Local_prty GOSUB Process_kbd
7170 ON KNOB,.01,Local_prty GOSUB Move_pointer
7172 IF Skips<To_select THEN
7174 DISP DS
7176 IF To_select>1 THEN
7178 Test$=" Select "&VAL$(Skips+1)&" of "&VAL$(To_select)
ELSE
7180 Test$=" Select"
7182 END IF
7184 ON KEY KO LABEL Test$,Local_prty GOSUB Select_item
7186 ELSE IF To_select>0 THEN
7188 DISP " Selection process complete ..."
ELSE
7190 DISP " Menu for information only ..."
7192 END IF
7194 ON KEY KO LABEL "Accept",Local_prty GOTO Exit_line
7200 END IF
7202 IF Active_screen=Screen_cnt THEN
7204 ON KEY KO+1 LABEL "Next Screen",Local_prty GOSUB Next_screen
ELSE OFF KEY KO+1
7210 END IF
7212 IF Active_screen>1 THEN
7214 ON KEY KO+2 LABEL "Last Screen",Local_prty GOSUB Last_screen
ELSE OFF KEY KO+2
7220 END IF
7222 IF Skips>0 THEN
7224 ON KEY KO+3 LABEL "Reset Select",Local_prty GOSUB Select_reset
ELSE OFF KEY KO+3
7230 END IF
7232 IF To_select>0 THEN
7234 ON KEY KO+4 LABEL "Abort ",Local_prty GOTO Escape_line
ELSE OFF KEY KO+4
7240 END IF
7242 IF Screen_cnt>2 THEN
7244 ON KEY KO+6 LABEL "Jump to Screen",Local_prty GOSUB Jump_to_scn
ELSE OFF KEY KO+6
7250 END IF
7252 IF Exit_flag THEN Exit_line
7254 GOTO Key_Loop
7256 MAT Choose=(0)
7258 Escape_line:Skips=0
7260 To_select=0
7262 Exit_line:OFF KEY
7264 OFF KNOB
7266 OFF KBD
7268 OUTPUT KBD:CHR$(255)&CHR$(75);
7270 PRINT CHR$(128);
7272 ! everything cleared, now go back to work.
7274 RETURN
7276 !
7278 !
7280 !

```

```

7282 Next_screen:
7284 OFF KBD
7286 OFF KNOB
7288 OFF KEY
7290 IF Active_screen=Screen_cnt THEN RETURN
7292 Active_screen=Active_screen+1
7294 GOSUB Write_screen
7296 RETURN
7300 !
7302 !
7304 !
7306 Last_screen:
7308 OFF KBD
7310 OFF KNOB
7312 OFF KEY
7314 IF Active_screen=1 THEN RETURN
7316 Active_screen=Active_screen-1
7318 GOSUB Write_screen
7320 RETURN
7322 !
7324 !
7326 !
7328 Jump_to_errors:DISP " Not a valid screen number ... try again."
7330 BEEP
7332 WAIT 1.8
7334 Jump_to_scn: !
7336 OFF KBD
7338 OFF KNOB
7340 OFF KEY
7342 DISP " ENTER the screen number desired (1 to ";screen_cnt,".";
7344 Test=TRIM$(Test$)
7346 IF LEN(Test$)=0 THEN Jump_to_return
7348 ON ERROR GOTO Jump_to_errors
7350 Temp=INT(VAL(Test$))
7352 OFF ERROR
7354 IF Temp<1 OR Temp>Screen_cnt THEN Jump_to_errors
7356 Active_screen=Temp
7358 GOSUB Write_screen
7360 Jump_to_return:
7362 DISP CHR$(12)
7364 Test$=""
7366 RETURN
7368 !
7370 !
7372 !
7374 !
7376 Select_item: !
7378 OFF KBD
7380 OFF KNOB
7382 OFF KEY
7384 IF NOT Pointinteractive THEN
7386 DISP "NO additional selections for this screen."
7388 BEEP
7390 WAIT 2
7392 DISP CHR$(12);
7394 RETURN
7396 END IF
7398 IF Skips=To_select THEN
7400 IF To_select=0 THEN
7402 DISP "this menu is for information only.";
7404 DISP " no selection allowed."
ELSE
7406 DISP "All selections have been filled.";
7408 DISP " Select reset' to repeat."
7410 END IF
7412 !

```



```

7678 PRINT CHR$(128);
7680 END IF
7682 PRINT TABXY(1,Pointer);Marker$,CHR$(128);
7684 END IF
7686 RETURN
7688 !
7690 !
7692 !
7694 !
7696 ! write screen: write the screen pointed to by Active_screen
7698 ! home and clear screen
7700 OUTPUT KBD;CHR$(255)&CHR$(84)&CHR$(255)&CHR$(75);
7702 Knobcount=KnobX-Knoby ! Clear knob and keyboard
7704 Knobcount=0
7706 Tests=KBD$
7708 Tests=""
7710 !
7712 ! PRINT TABXY(1,First_line-1);CHR$(132);" Item #1 screen #";
7714 ! PRINT USING "#,20,4A,2D,3A";Active_screen," of ",Screen_cnt;" | "
7716 ! J=0
7718 ! REPEAT
7720 ! IF J=Last_item(Active_screen)-First_item(Active_screen) THEN
7722 ! PRINT CHR$(132);
7724 ! PRINT TABXY(1,First_line+J);RPTS(" ",80)
7726 ! ELSE
7728 ! PRINT CHR$(128);
7730 ! END IF
7732 ! PRINT TABXY(5,First_line+J);
7734 ! PRINT USING "#D,A,#,First_item(Active_screen)+J,|"
7736 ! IF SKIPS>0 THEN ! make this line inverse video
7738 ! FOR I=1 TO SKIPS
7740 ! IF First_item(Active_screen)+J=Choose(I) THEN
7742 ! PRINT CHR$(129);
7744 ! END IF
7746 ! NEXT I
7748 ! END IF
7750 ! PRINT TABXY(10,First_line+J);Items$(First_item(Active_screen)+J)
7752 ! J=J+1
7754 ! UNTIL J>=(Last_item(Active_screen)-First_item(Active_screen)+1)
7756 ! Last_line=Last_item(Active_screen)-First_item(Active_screen)
7758 ! Last_line=Last_line+First_line
7760 ! !
7762 ! ! set marker to first non-selected item.
7764 ! !
7766 ! Pointeractive=0
7768 ! IF To_select>0 THEN Pointeractive=1
7770 ! IF SKIPS>0 AND Pointeractive=1 THEN ! find first non-selected item
7772 ! J=0
7774 ! LOOP
7776 ! Pointer=First_line+J
7778 ! FOR I=1 TO SKIPS
7780 ! IF First_item(Active_screen)+J=Choose(I) THEN Pointer=0
7782 ! NEXT I
7784 ! J=J+1
7786 ! IF First_line>J>Last_line THEN
7788 ! Pointeractive=0
7790 ! Pointer=First_line
7792 ! END IF
7794 ! EXIT IF Pointer<>0
7796 ! END LOOP
7800 ! ELSE
7802 ! Pointer=First_line
7804 ! END IF
7806 ! IF Pointeractive THEN
7808 ! IF Pointer=Last_line THEN

```

```

7810 PRINT CHR$(132);
7812 ELSE PRINT CHR$(128);
7814 END IF
7816 PRINT TABXY(1,Pointer);Marker$,CHR$(128);
7818 END IF
7820 RETURN
7822 SUBEND
7824 !
7826 !
7828 ! *****
7830 !
7832 ! SUB Errortrap
7834 ! Errortrap: ! Original: 13 Nov 1984
7836 ! Revision: 02 Dec 1987
7838 ! Trap most errors here
7840 ! COM /Files/ Diskdrives$(20),FileNames$(10)
7842 ! DIM Files$(20),Tests$(160),Whats$(20),ACS$(5)
7844 ! BEEP 400,.6
7846 ! SELECT ERRN
7848 ! CASE 5:
7850 ! DISP "DUPLICATE FILE NAME: ";FileNames$;
7852 ! DISP "...PURGE old one? (Y/N)";
7854 ! INPUT Whats$
7856 ! Whats$=TRIMS(Whats$)
7858 ! SELECT Whats$(1,1)
7860 ! CASE "y","Y":
7862 ! PURGE FileNames$&Diskdrives$
7864 ! CASE ELSE
7866 ! ACS="VALID"
7868 ! CALL Enterfilename(ACS)
7870 ! END SELECT
7872 ! CASE 52,53
7874 ! DISP "Improper FILE NAME --- ENTER NEW FILE NAME";
7876 ! OUTPUT 2 USING "#,K,K,;" #";FileNames$
7878 ! INPUT FileNames$
7880 ! FileNames$=TRIMS(FileNames$)
7882 ! CASE 54
7884 ! DISP "FILE: ";FileNames$;" is not on this disk, please insert";
7886 ! DISP " correct disk"
7888 ! CALL Pause_key_on
7890 ! CASE 6:
7892 ! DISP "this disk is full, PLEASE insert clean disk"
7894 ! CALL Pause_key_on
7896 ! CASE 56
7898 ! DISP "DATA INPUT disk must be in drive!! ";
7900 ! DISP "...CONTINUE when ready."
7902 ! CALL Pause_key_on
7904 ! CASE 72,73,76
7906 ! DISP Diskdrives$;
7908 ! DISP " is not available, type correct";
7910 ! DISP " unit specifier (ie. ;:707,0)";
7912 ! OUTPUT 2 USING "#,;"Diskdrives$
7914 ! INPUT Diskdrives$
7916 ! CASE 60
7918 ! DISP "CHECK DISK drive coord!"
7920 ! CALL Pause_key_on
7922 ! CASE ELSE
7924 ! DISP ERRMS;" 'CONTINUE' when fixed"
7926 ! CALL Pause_key_on
7928 ! END SELECT
7930 ! DISP CHR$(12)
7932 ! SUBEXIT
7934 ! SUBEND
7936 !
7938 ! *****
7940 !

```



```

7942 SUB Pause_key_on
7943 Pause_key_on: i Make sure that CONTINUE key exists.
7944 i Original: 02 Dec 1987
7945 i Revision: 02 Dec 1987
7946
7947 COM /Sys/ Sys_id$[10]
7948 IF Sys_id$[1,4]="S300" THEN i reset to S300 system keys
7949 CONTROL KBD,15;0
7950 CONTROL CRT,12;2
7951 LOAD KEY
7952 END IF
7953 PAUSE
7954 IF Sys_id$[1,4]="S300" THEN i set to S200 compatible keys
7955 OUTPUT KBD USING "K,#"SCRATCH KEY X"
7956 CONTROL KBD,15;1
7957 CONTROL CRT,12;0
7958 END IF
7959 SUBEXIT
7960 SUBEND
7961
7962 *****
7963 *****
7964 *****
7965 *****
7966 *****
7967 *****
7968 *****
7969 *****
7970 *****
7971 *****
7972 *****
7973 *****
7974 *****
7975 *****
7976 *****
7977 *****
7978 *****
7979 *****
7980 *****
7981 *****
7982 *****

```

(4) The -3 -dB correction is required because this is a direct measure of $\mathcal{L}(f)$ of two oscillators, assuming that the oscillators are of a similar type and that the noise contribution is the same for each oscillator. If one oscillator is sufficiently superior to the other, this correction is not required.

Other defined spectral densities can be calculated and plotted as desired. The plotted or stored value of the spectral density of phase fluctuations in decibels relative to one square radian (dBc rad²/Hz) is calculated as

$$S_{\delta\phi}(f) = \mathcal{L}(f) + 3. \quad (70)$$

The spectral density of phase fluctuations, in radians squared per hertz, is calculated as

$$S_{\delta\phi}(f) = 10 \exp(S_{\delta\phi}(f)/10), \quad (71)$$

The spectral density of frequency fluctuations, in hertz squared per hertz, is

$$S_{\delta\nu}(f) = f^2 S_{\delta\phi}(f). \quad (72)$$

where $S_{\delta\phi}(F)$ is in decibels with respect to 1 radian.

5. System Noise Floor Verification

A plot of the system noise floor (sensitivity) is obtained by repeating the automated measurement procedures with the system modified as shown in Fig. 12. Accurate measurements can be obtained using the configuration shown in Fig. 12a. The reference source supplies 10 dBm to one side of the mixer and 0 dBm to the other mixer input through equal path lengths; phase quadrature is maintained with the phase shifter.

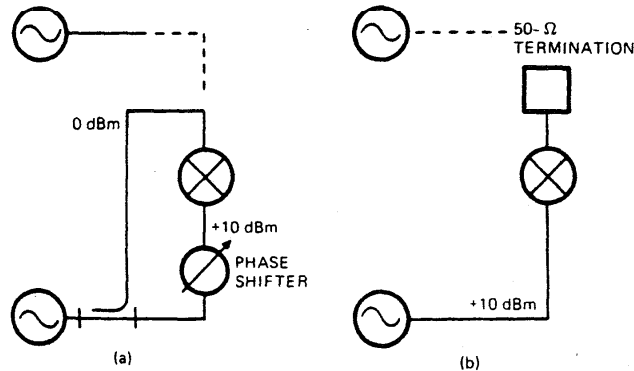


FIG. 12 System configurations for measuring the system noise floor (sensitivity): (a) configuration used for accurate measurements; (b) alternate configuration sometimes used.

The configuration shown in Fig. 12b is sometimes used and does not greatly degrade the noise floor because the reference signal of 10 dBm is larger than the signal frequency. See Sections IV.B and IV.C.4 for additional discussions related to system sensitivity and recommended system evaluation.

Proper selection of drive and output termination of the double-balanced mixer can result in improvement by 15 to 25 dB in the performance of phase-noise measurements, as discussed by Walls *et al.* (1976). The beat frequency between the two oscillators can be a sine wave, as previously mentioned, with proper low drive levels. This requires a proper terminating impedance for the mixer. With high drive levels, the mixer output waveform will be clipped. The slope of the clipped waveform at the zero crossings, illustrated by Walls *et al.* (1976), is twice the slope of the sine wave and therefore improves the noise floor sensitivity by 6 dB, i.e., the output signal, proportional to the phase fluctuations, increases with drive level. This condition of clipping requires characterization over the Fourier frequency range, as previously mentioned for the Hewlett-Packard 3047 phase noise measurement system. An amplifier can be used to increase the mixer drive levels for devices that have insufficient output power to drive the double-balanced mixers.

Lower noise floors can be achieved using high-level mixers when available drive levels are sufficient. A step-up transformer can be used to increase the mixer drive voltage because the signal and noise power increase in the same ratio, and the spectral density of phase of the device under test is unchanged, but the noise floor of the measurement system is reduced.

Walls *et al.* (1976) used a correlation technique that consisted primarily of two phase-noise measurement systems. At TRW the technique is used as shown in Fig. 13. The cross spectrum is obtained with the fast Fourier transform (FFT) analyzer that performs the product of the Fourier transform of one signal and the complex conjugate of the Fourier transform of

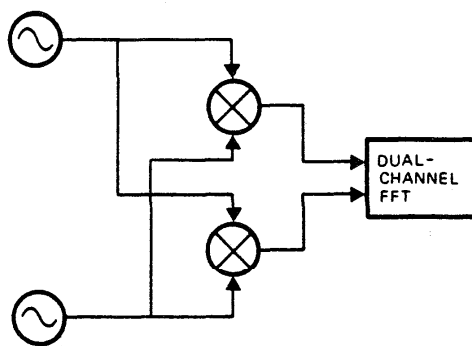


FIG. 13 Cross-spectrum measurement using the two-oscillator technique.

the second signal. This cross spectrum, which is a phase-sensitive characteristic, gives a phase and amplitude sensitivity measure directly. A signal-to-noise enhancement greater than 20 dB can be achieved.

If the double-balanced phase noise measurement system does not provide a noise floor sufficient for measuring a high-quality source, frequency multiplier chains can be used if their inherent noise is 10–20 dB below the measurement system noise. In frequency multiplication the noise increases according to

$$10 \log(\text{final frequency}/\text{original frequency}). \quad (73) *$$

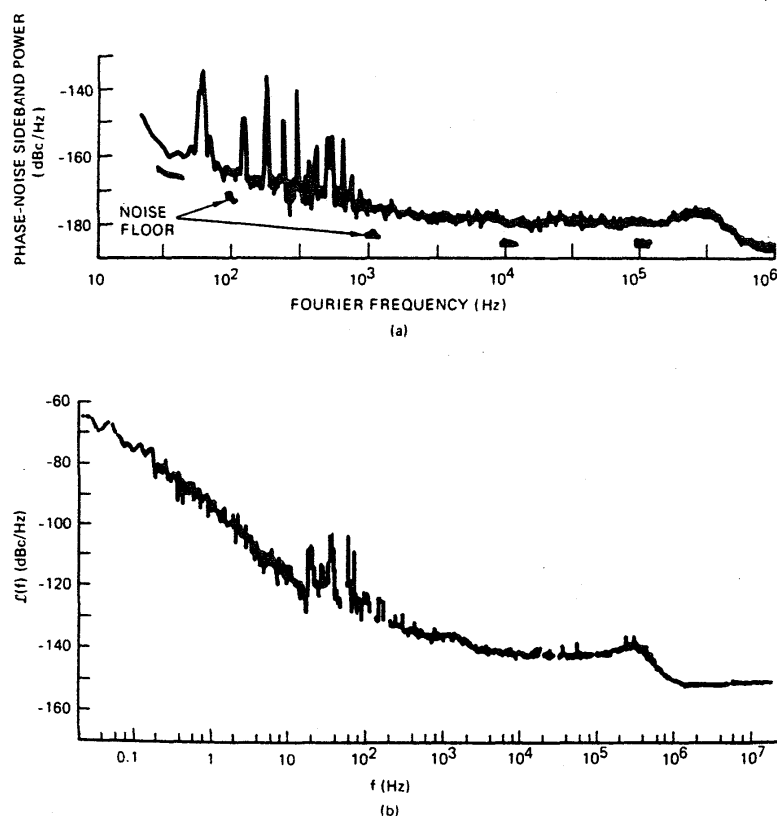


FIG. 14 Data plots of the automated phase-noise measurement system: (a) a high quality 5-MHz quartz oscillator; (b) combined noise of two H.P. 8662A synthesizers (subtract 3 dB for a single unit).

* See Appendix Note # 32

The following equation is used to correct for noise-floor contribution P_{nf} , in dBc/Hz, if desired or necessary:

$$\mathcal{L}(f)(\text{corrected}) = -\mathcal{L}(f) + 10 \log \left[\frac{P_{\mathcal{L}(f)} - P_{nf}}{P_{\mathcal{L}(f)}} \right] \quad (74)$$

The correction for noise-floor contribution can also be obtained by using the measurement of $S_{\delta v}(f)$ of Eq. (57). Measurement of $S_{\delta v}(f)$ of the oscillator plus floor is obtained, then $S_{\delta v}(f)$ is obtained for the noise floor only. Then,

$$S_{\delta v}(f) \Big|_{\text{cor}} = S_{\delta v}(f) \Big|_{(\text{osc} + \text{nf})} - S_{\delta v}(f) \Big|_{\text{nf}} \quad (75)$$

Figure 14a shows a phase noise plot of a very high-quality (5-MHz) quartz oscillator, measured by the two-oscillator technique. The sharp peaks below 1000 Hz represent the 60-Hz line frequency of the power supply and its harmonics and are not part of the oscillator phase noise. Figure 14b shows measurements to 0.02 Hz of the carrier at a frequency of 20 MHz.

IV. Single-Oscillator Phase-Noise Measurement Systems and Techniques

*

The phase-noise measurements of a single-oscillator are based on the measurement of *frequency fluctuations* using discriminator techniques. The practical discriminator acts as a filter with finite bandwidth that suppresses the carrier and the sidebands on both sides of the carrier. The ideal carrier-suppression filter would provide infinite attenuation of the carrier and zero attenuation of all other frequencies. The effective Q of the practical discriminator determines how much the signals are attenuated.

Frequency discrimination at very high frequencies (VHF) has been obtained using slope detectors and ratio detectors, by use of lumped circuit elements of inductance and capacitance. At ultrahigh frequencies (UHF) between the VHF and microwave regions, measurements can be performed by beating, or heterodyning, the UHF signal with a local oscillator to obtain a VHF signal that is analyzed with a discriminator in the VHF frequency range. Those techniques provide a means for rejecting residual amplitude-modulated (AM) noise on the signal under test. The VHF discriminators usually employ a limiter or ratio detector.

Ashley *et al.* (1968) and Ondria (1968) have discussed the microwave cavity discriminator that rejects AM noise, suppresses the carrier so that the input level can be increased, and provides a high discriminated output to improve the signal-to-noise floor ratio. The delay line used as an FM discriminator has been discussed by Tykulsky (1966), Halford (1975), and

* See Appendix Note # 6

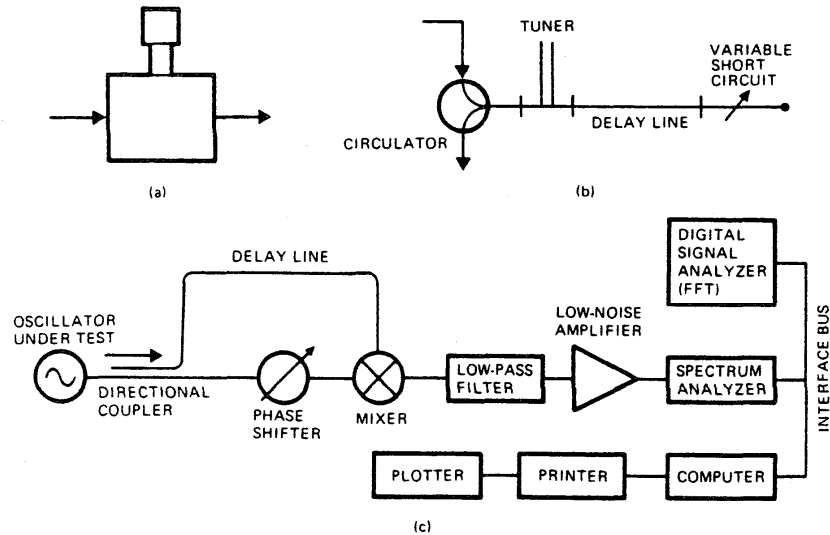


FIG. 15 Single-oscillator phase-noise measurement techniques: (a) cavity discriminator; (b) reflective-type delay-line discriminator; (c) one-way delay line.

Ashley *et al.* (1968). Ashley *et al.* (1968) proposed the reflective-type delay-line discriminator shown in Fig. 15b. The cavity can also be used to replace delay line. The one-way delay line shown in Fig. 15c is implemented in the TRW measurement systems. The theory and applications set forth in this section are based on a system of this particular type.

A. THE DELAY LINE AS AN FM DISCRIMINATOR

1. The Single-Oscillator Measurement System

The single-oscillator signal is split into two channels in the system shown in Fig. 15. One channel is called the nondelay or reference channel. It is also referred to as the local-oscillator (LO) channel because the signal in this channel drives the mixer at the prescribed impedance level (the usual LO drive). The signal in the second channel arrives at the mixer through a delay line. The two signals are adjusted for phase quadrature with the phase shifter, and the output of the mixer is a fluctuating voltage, analogous to the frequency fluctuations of the source, centered on approximately zero dc volts.

The delay line yields a phase shift by the time the signal arrives at the balanced mixer. The phase shift depends on the instantaneous frequency of

the signal. The presence of frequency modulation (FM) on the signal gives rise to differential phase modulation (PM) at the output of the differential delay and its associated (nondelay) reference line. This relationship is linear if the delay τ_d is nondispersive. This is the property that allows the delay line to be used as an FM discriminator. In general, the conversion factors are a function of the delay (τ_d) and the Fourier frequency f but not of the carrier frequency.

The differential phase shift of the nominal frequency ν_0 caused by the delay line is

$$\Delta\phi = 2\pi\nu_0\tau_d. \quad (76)$$

where τ_d is the time delay.

The phase fluctuations at the mixer are related to the frequency fluctuations (at the rate f) by

$$\delta\phi = 2\pi\tau_d\delta\nu(f). \quad (77)$$

The spectral density relationships are

$$S_{\delta\phi}(f)\Big|_{\text{mixer}} = (2\pi\tau_d)^2 S_{\delta\nu}(f)\Big|_{\text{osc}} \quad (78)$$

and

$$S_{\delta\nu}(f) = f^2 S_{\delta\phi}(f). \quad (79)$$

Then,

$$S_{\delta\phi}(f)\Big|_{\text{dlm}} = (2\pi f\tau_d)^2 S_{\delta\phi}(f)\Big|_{\text{osc}} \quad (80)$$

where the subscript dlm indicates delay-line method. From Eq. (56), the spectral density of phase for the two-oscillator technique, in radians squared per hertz, is

$$S_{\delta\phi}(f) = 4 \frac{S_{\delta\nu}(f)}{(V_{\text{ptp}})^2} = \frac{S_{\delta\nu}(f)}{2(V_{\text{rms}})^2} = \left[\frac{(\delta v_{\text{rms}})^2}{2(V_{\text{rms}})B} \right] \quad (81)$$

because

$$(V_{\text{ptp}})^2 = 8(v_{\text{rms}})^2 = 4(V_p)^2 = 4[2(v_{\text{rms}})^2]$$

and

$$\mathcal{L}(f) = 2(S_{\delta\nu}(f)/(V_{\text{ptp}})^2) = (\delta v_{\text{rms}})^2/4(V_{\text{rms}})^2 B \quad (82)$$

per hertz.

The sensitivity (noise floor) of the two-oscillator measurement system includes the thermal and shot noise of the mixer and the noise of the base-band preamplifier (referred to its input). This noise floor is measured with the oscillator under test inoperative. The measurement system sensitivity of the two-oscillator system, on a per hertz density basis (dBc/Hz) is

$$\mathcal{L}(f)_{\text{nf}} = 10 \log[2(\delta v_n)^2/(V_{\text{ptp}})^2], \quad (83)$$

where δv_n is the rms noise voltage measured in a one-hertz bandwidth.

The two-oscillator system therefore yields the output noise from both oscillators. If the reference oscillator is superior in performance, as assumed in the previous discussions, then one obtains a direct measure of the noise characteristics of the oscillator under test. If the reference and test oscillators are the same type, a useful approximation is to assume that the measured noise power is twice that associated with one noisy oscillator. This approximation is in error by no more than 3 dB for the noisier oscillator. Substituting in Eq. (80) and using the relationships in Eq. (56), we have, per hertz,

$$\mathcal{L}(f) \Big|_{\text{dim}} = 2[(\delta v_{\text{rms}})^2/(V_{\text{ptp}})^2](2\pi f\tau_d)^2 \quad (84)$$

Examination of this equation reveals the following.

(1) The term in the brackets represents the two-oscillator response. *Note that this term represents the noise floor of the two-oscillator method.* Therefore, adoption of the delay-line method results in a higher noise by the factor $(2\pi f\tau_d)^2$ when compared with the two-oscillator measurement method. The sensitivity (noise floor) for delay lines with different values of time delay are illustrated in Fig. 17.

(2) Equation (84) also indicates that the measured value of $\mathcal{L}(f)$ is periodic in $\omega = 2\pi f$. This is shown in Fig. 21. The first null in the responses is at the Fourier frequency $f = 1/\tau_d$. The periodicity indicates that the calibration range of the discriminator is limited and that valid measurements occur only in the indicated range, as verified by the discriminator slope shown in Fig. 16. (See Fig. 23.)

(3) The maximum value of $(2\pi f\tau_d)^2$ can be greater than unity (it is 4 at $f = 1/2\tau_d$). This 6-dB advantage is utilized in the noise-floor measurement. However, it is beyond the valid calibration range of the delay-line system. The 6-dB advantage is offset by the line attenuation at microwave frequencies, as discussed by Halford (1975).

The delay-line discriminator system has been analyzed in terms of a power-limited system (a particular idealized system in which the choice of power oscillator voltage, the attenuator of the delay line, and the conversion loss

of the mixer are limited by the capability of the mixer) by Tykulsky (1966), Halford (1975), and Ashley *et al.* (1977). For this particular case, Eq. (83) indicates that an increase in the length of the delay line (to increase τ_d for decorrelation of Fourier frequencies closer to the carrier) results in an increase in attenuation of the line, which causes a corresponding decrease in V_{ptp} . The optimum length occurs where τ_d is such that the decrease in V_{ptp} is approximately compensated by the increase in $(2\pi f\tau_d)$, i.e., where

$$\frac{d}{d\tau_d} \frac{2\pi f\tau_d}{V_{\text{ptp}}} = 0. \quad (85)$$

This condition occurs where the attenuation of the delay line is 1 Np (8.686 dB). However, when the system is not power limited, the attenuation of the delay line is not limited, because the input power to the delay line can be adjusted to maintain V_{ptp} at the desired value. The optimum delay-line length is determined at a particular selectable frequency. However, since the attenuation varies slowly (approximately proportional to the square root of frequency), this characteristic allows near-optimum operation over a considerable frequency range without appreciable degradation in the measurements.

A practical view of the time delay (τ_d) and Fourier-frequency functional relationship can be obtained by reviewing the basic concepts of the dual-channel time-delay measurement system discussed by Lance (1964). If the differential delay between the two channels is zero, there is no phase difference at the detector output when a swept-frequency cw signal is applied to the system. Figure 16 shows the detected output interference display when a swept-frequency cw signal (zero to 4 MHz) is applied to a system that has

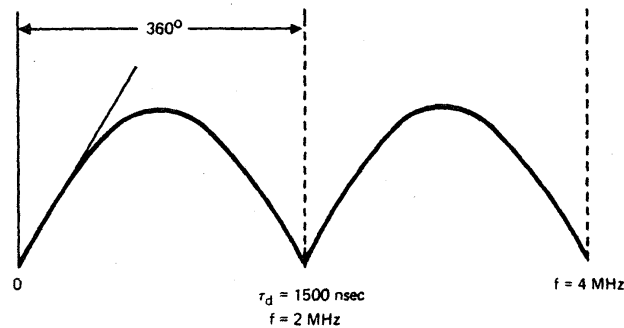


FIG. 16 Swept-frequency interference display at the output of a dual-channel system with a differential delay of 500 nsec.

a differential delay of 500 nsec between the two channels. The signal amplitudes are assumed to be almost equal, thus producing the familiar voltage-standing-wave pattern or interference display. Because this is a two-channel system, there is a null every 360° , as shown.

2. System Sensitivity (Noise Floor) When Using the Differential Delay-Line Technique

Halford (1975) has shown that the sensitivity (noise floor) of the single-oscillator differential delay-line technique is reduced relative to the two-oscillator techniques. The sensitivity is modified by the factor

$$S_d = 2(1 - \cos 2\pi f\tau_d). \quad (86)$$

For $\omega\tau_d = 2\pi f\tau_d \ll 1$ a good approximation is

$$S_d^2 = 2(1 - \cos 2\pi f\tau_d) = (\omega\tau_d)^2 [1 - \frac{1}{12}(\omega\tau_d)^2] = (2\pi f\tau_d)^2 = \theta^2, \quad (87)$$

where θ is the phase delay of the differential delay line evaluated at the frequency f . Figure 17 shows the relative sensitivity (noise floor) of the two-oscillator technique and the single-oscillator technique with different delay-line lengths. The f^{-2} slope is noted at Fourier frequencies beyond

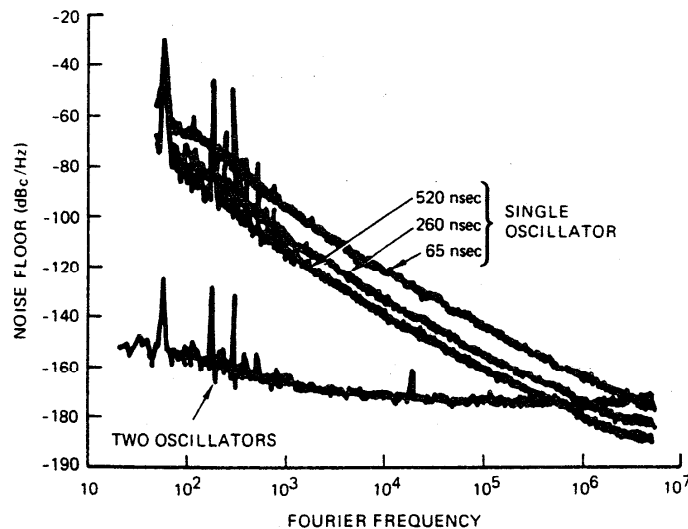


FIG. 17 Relative sensitivity (noise floor) of single-oscillator and two-oscillator phase-noise measurement systems.

about 1 kHz. For Fourier frequencies closer to the carrier, the slope is f^{-3} , i.e., the sum of the f^{-2} slope of Eq. (87) and the f^{-1} flicker noise.

Phase-locked sources have phase-noise characteristics that cannot be measured at close-in Fourier frequencies using this basic system. The relative sensitivity of the system can be improved by using a dual (two-channel) delay-line system and performing cross-spectrum analysis, which will be presented in this chapter.

Labaar (1982) developed the delay-line rf bridge configuration shown in Fig. 18. At microwave frequencies where a high-gain amplifier is available, suppression of the carrier by the rf bridge allows amplification of the noise going into the mixer. A relative sensitivity improvement of 35 dB has been obtained without difficulty. The limitations of the technique depend on the available rf power and the carrier suppression by the bridge. Naturally, if the rf input to the bridge is high one must use the technique with adequate precautions to prevent mixer damage that can occur by an accidental bridge unbalance. Labaar (1982) indicated the added advantage of using the rf bridge carrier-suppression technique when attempting to measure phase noise close to the carrier when AM noise is present. Figure 19 shows the

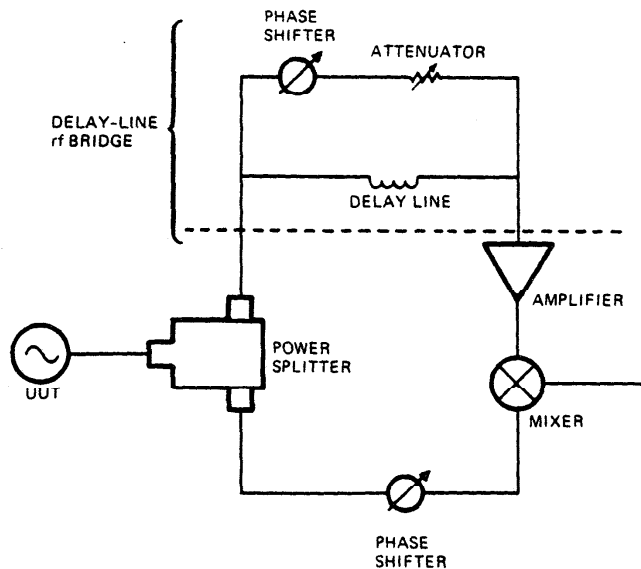


Fig. 18 Carrier suppression using an rf bridge to increase relative sensitivity. (Courtesy Instrument Society of America.)

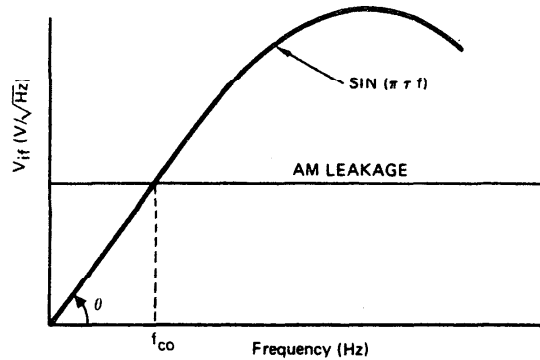


FIG. 19 Phase detector output (AM-PM crossover); τ , delay time.

mixer output for phase (PM) and amplitude (AM) noise in the single-oscillator delay-line FM discriminator system. It is noted that the phase noise and AM noise intersect and that the AM will therefore limit the measurement accuracy near the carrier. Even though AM noise is much lower than phase noise in most sources, and even though the AM is normally suppressed about 20 dB, there is still AM at the mixer output. This output is AM leakage and is caused by the finite isolation between the mixer ports. The two-oscillator technique does not experience this problem to this extent because the phase noise and AM noise maintain their relative relationships at the mixer output independent of the offset frequency from the carrier.

B. CALIBRATION AND MEASUREMENTS USING THE DELAY LINE AS AN FM DISCRIMINATOR

The block diagram of a practical single-oscillator phase noise measurement system is shown in Fig. 20. The signals in the delay-line channel of the system experience the one-way delay of the line. With adequate source power, the system is not limited to the optimum 1 Np (8.686 dB) previously discussed for a power-limited system. Measurements are performed using the following operational procedures.

- (1) Measure the tracking spectrum analyzer IF bandwidths as set forth in Section III.C.1.
- (2) Establish the system power levels (Section IV.B.1).
- (3) Establish the discriminator calibration factor (Section IV.B.2).
- (4) Measure and plot the oscillator characteristics in the automatic system used (Section IV.B.3).
- (5) Measure the system noise floor (sensitivity) (Section IV.B.4).

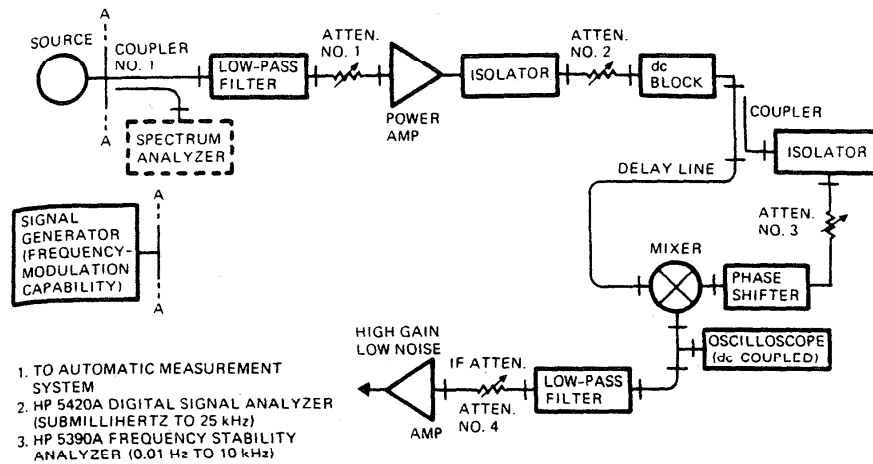


Fig. 20 Single-oscillator phase noise measurement system using the delay line as an FM discriminator. (From Lance *et al.*, 1977a.)

1. System Power Levels

The system power levels are set using attenuators, as shown in Fig. 20. Because the characteristic impedance of attenuator No. 4 is 50 Ω , mismatch errors will occur if the mixer output impedance is not 50 Ω . As previously discussed, the mixer drive levels are set so that the mixer output signal, as observed during calibration, is sinusoidal. This has been accomplished in TRW systems with a reference (LO) signal level of 10 dBm and a mixer input level of about 0 dBm from the delay line.

A power amplifier can be used to increase the source signal to the measurement system. This amplifier must not contribute appreciable additional noise to the signal.

2. Discriminator Calibration

The discriminator characteristics are measured as a function of frequency and voltage. The hertz-per-volt sensitivity of the discriminator is defined as the *calibration factor* (CF). The calibration process involves measuring the effects of intentional modulation of the source (carrier) frequency. A known modulation index must be obtained to calculate the calibration factors of the discriminator. The modulation index is obtained by using amplitude modulation to establish the carrier-to-sideband ratio when there is considerable instability of the source or when the source cannot be frequency modulated.

It is convenient to consider the system equations and calibration techniques in terms of frequency modulation of stable sources. If the source to be measured cannot be frequency modulated, it must be replaced, during the calibration process, with a modulatable source. The calibration process will be described using a modulatable source and a 20-kHz modulation frequency. However, other modulation frequencies can be used. The calibration factor of this type discriminator has been found to be constant over the usable Fourier frequency range, within the resolution of the measuring technique. The calibration factor of the discriminator is established after the system power levels have been set with the unit under test as the source.

The discriminator calibration procedures are as follows.

- (1) Set attenuator No. 4 (Fig. 20) to 50 dB.
- (2) Replace the oscillator under test with a signal generator or oscillator that can be frequency modulated. *The power output and operating frequency of the generator must be set to the same precise frequency and amplitude values that the oscillator under test will present to the system during the measurement process.*
- (3) Select a modulation frequency of 20 kHz and increase the modulation until the carrier is reduced to the first Bessel null, as indicated on the spectrum analyzer connected to coupler No. 1. This establishes a modulation index ($m = 2.405$).
- (4) Adjust the phase shifter for zero volts dc at the output of the mixer, as indicated on the oscilloscope connected as shown in Fig. 20. *This establishes the quadrature condition for the two inputs to the mixer.* This quadrature condition is continuously monitored and is adjusted if necessary.
- (5) Tune the tracking spectrum analyzer to the modulation frequency of 20 kHz. The power reading at this frequency is recorded in the program and is corrected for the 50-dB setting of attenuator No. 4, which will be set to zero decibel indication during the automated measurements.

$$P(\text{dBm}) = (-\text{dBm power reading}) + 50 \text{ dB} \quad (88)$$

This power level is converted to the equivalent rms voltage that the spectrum analyzer would have read if the total signal had been applied:

$$V_{\text{rms}} = \sqrt{10^{P/10}/1000 + R}. \quad (89)$$

- (6) The discriminator calibration factor can now be calculated because this power in dBm can be converted to the corresponding rms voltage using the following equation:

$$V_{\text{rms}} = \sqrt{(10^{P/10}/1000) \times R}, \quad (90)$$

where $R = 50 \Omega$ in this system.

(7) The discriminator calibration factor is calculated in hertz per volt as

$$CF = mf_m/\sqrt{2}V_{rms} = 2.405f_m/\sqrt{2}V_{rms}. \quad (91)$$

The modulation index m for the first Bessel null as used in this technique is 2.405. The modulation frequency is f_m .

3. Measurement and Data Plotting

After the discriminator is calibrated, the modulated signal source is replaced with the frequency source to be measured. Quadrature of the signals into the mixer is reestablished, attenuator No. 4 (Fig. 20) is set to 0 dB, and the measurement process can begin.

The measurements, calculations, and data plotting are completely automated. The calculator program selects the Fourier frequency, performs autoranging, and sets the bandwidth, and measurements of Fourier frequency power are performed by the tracking spectrum analyzer. Each Fourier frequency noise-power reading P_n (dBm) is converted to the corresponding rms voltage by

$$v_{1rms} = \sqrt{10^{(P_n + 2.5)/10}/1000 \times R}. \quad (92)$$

The rms frequency fluctuations are calculated as

$$\delta v_{rms} = v_{1rms} \times CF. \quad (93)$$

The spectral density of frequency fluctuations in hertz squared per hertz is calculated as

$$S_{\delta v}(f) = (\delta v_{rms})^2/B, \quad (94)$$

where B is the measured IF noise-power bandwidth of the spectrum analyzer. The spectral density of phase fluctuations in radians squared per hertz is calculated as

$$S_{\delta\phi}(f) = S_{\delta v}(f)/f^2. \quad (95)$$

The NBS-designated spectral density in decibels (carrier) per hertz is calculated as

$$\mathcal{L}(f)_{dB} = 10 \log \frac{1}{2} S_{\delta\phi}(f). \quad (96)$$

Spectral density is plotted in real time in our program. However, the data can be stored and the desired spectral density can be plotted in other forms. Integrated phase noise can be obtained as desired.

4. Noise Floor Measurements

The relative sensitivity (noise floor) of the single-oscillator measurement system is measured as shown in Fig. 12a for the two-oscillator technique. The delay line must be removed and equal channel lengths constructed, as in Fig. (12a). The same power levels used in the original calibration and measurements are reestablished, and the noise floor is measured at specific Fourier frequencies, using the same calibration-measurement technique, or by repeating the automated measurement sequence.

A correction for the noise floor requires a measurement of the rms voltage of the oscillator ($v_{1,rms}$) and a measurement of the noise floor rms voltage ($v_{2,rms}$). These voltages are used in the following equation to obtain the corrected value:

$$v_{rms} = \sqrt{(v_{1,rms})^2 - (v_{2,rms})^2}. \quad (97)$$

The value v_{rms} is then used in the calculation of frequency fluctuations. If adequate memory is available, each value of $v_{1,rms}$ can be stored and used after the other set of measurements are performed at the same Fourier frequencies.

The following technique was developed by Labaar (1982). Carrier suppression is obtained using the rf bridge illustrated in Fig. 18. One can easily improve sensitivity more than 40 dB. At 2.0 and 3.0 GHz 70-dB carrier suppression was realized. In general, the improvement in sensitivity will depend on the availability of an amplifier or adequate input power.

Figure 21 shows the different noise floors in a delay-line bridge discriminator. It is good measurement discipline to always determine these noise floors; also, the measurements, displayed in Fig. 21, give a quick understanding of the physical process involved. The first trace is obtained by terminating the input of the baseband spectrum analyzer. The measured output noise power is then a direct measure of the spectrum analyzer's noise figure (NF). The input noise is thermal noise and is usually indicated by "KTB," which is short for "the thermal noise power at absolute temperature of T degrees K(elvin) per one hertz bandwidth (B). This KTB number is, at 18°C, about -174 dBm/Hz.

Figure 21a shows that trace number 1 for frequencies above about 1 kHz is level with a value of about -150 dBm = $-(174-24)$ dBm, which means that the spectrum analyzer has an NF of 24 dB. At 20 Hz the NF has gone up to about 48 dB. To improve the NF, a low-noise (NF, 2dB), low-frequency (10 Hz-10 MHz) amplifier is inserted as a preamplifier. Terminating its input now results in trace number 2. At the high frequency end, the measured power goes up by about 12-13 dB, and the amplifiers gain is 34 dB. This means that the NF is improved by $34 - 12-13 = 21-22$ dB, which is an

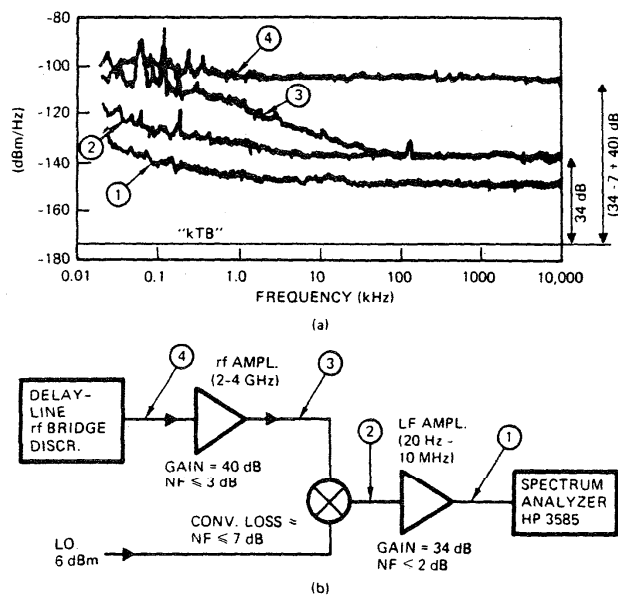


FIG. 21 (a) Noise contribution analysis; (b) phase noise test setup using a delay-line rf bridge discriminator (rf = 2.8 GHz; O, termination points. NF, noise figure; LF, low frequency. (From Seal and Lance, 1981.)

$NF \approx 2-3$ dB as expected, i.e., the first stage noise predominates. The low-frequency end at 20 Hz gives an NF of 26 dB, which overall is quite an improvement.

In trace number 3 the mixer is included with its rf (signal) port terminated. It is clear from this trace that certainly up to 100 kHz, the noise generated by the mixer diodes being "pumped" by the LO signal dominates. This case represents the "classic" delay-line discriminator. The last trace (number 4) includes the low-noise, high-gain rf amplifier that can be used because the carrier is suppressed in the delay-line rf bridge discriminator, in contrast to the classic delay-line discriminator case. This trace shows that from 1 kHz on up the measured output power is flat, representing a 2-3-dB NF.

At about 20-40 Hz, trace numbers 3 and 4 begin nearing their cross-over floor. In this particular case, which is discussed in full by Labaar (1982), the measurement systems noise floor (resolution) has been improved by 40 dB.

Figure 22 shows plots of phase noise as measured at two frequencies using delay lines of different lengths. The delay line used measure at 600 MHz

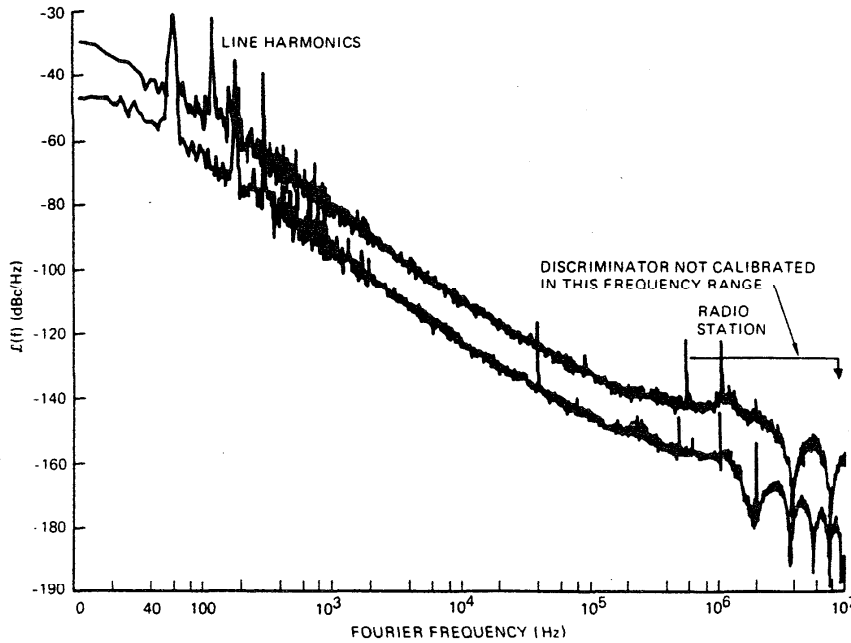


FIG. 22 Phase noise of 600-MHz oscillator multiplied to 2.4 GHz (From Lance *et al.*, 1977a.)

was about 500 nsec long, as noted by the first null, i.e., the reciprocal of the Fourier frequency of 2 MHz is the approximate differential time delay. Note that a shorter delay line (approximately 250 nsec differential) is used to measure the higher frequency because the delay-line discriminator calibration is valid only to a Fourier frequency at approximately 35% of the Fourier frequency at which the first null occurs, if a linear transfer function is assumed.

The actual transfer function of a delay-line discriminator (classic and rf bridge types) is sinusoidal, as shown in Fig. 23a. The baseband spectrum analyzer measures power in a finite bandwidth, and as a consequence it is possible to measure through a transfer-function null if the noise power does not change substantially over a spectrum-analyzer bandwidth. The following power relations then hold:

$$P_{\text{meas}}(\omega) = 1/\Delta\omega \int_{\omega - \Delta\omega/2}^{\omega + \Delta\omega/2} P(\omega') d\omega' \approx \frac{P(\omega)}{\Delta\omega} \int_{\omega - \Delta\omega/2}^{\omega + \Delta\omega/2} d\omega' = P(\omega). \quad (98)$$

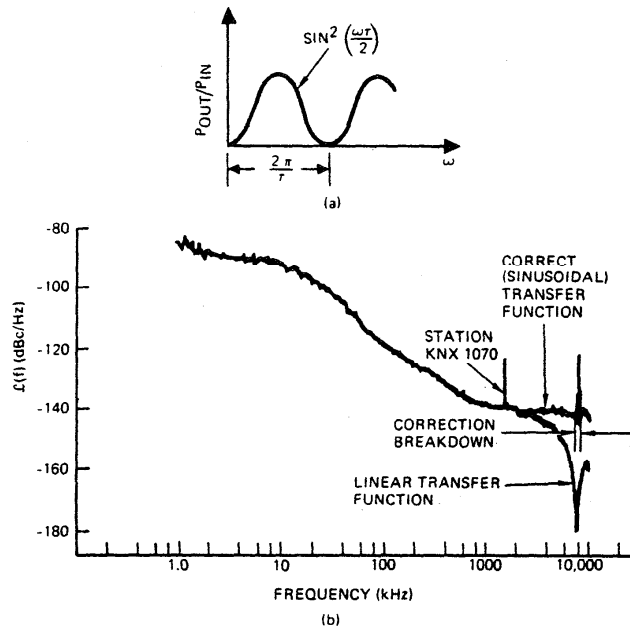


FIG. 23 Transfer functions for a delay-line rf bridge discriminator: (a) actual; (b) approximate (linear) and "correct" (sinusoidal). Phase noise: H.P. 8672A at 2.4 GHz.

Figure 23b shows the results using a linear approximation and the "correct" transfer function for a delay-line rf bridge discriminator. The correct transfer function breaks down close to the null because the signal level drops below the system's noise floor, as explained by Labaar (1982).

Using the sinusoidal transfer function in the calculator software gives correct results barring frequency intervals of 5 to 10 spectrum analyzer's bandwidths ($10 \times 30 = 300$ kHz) centered at the transfer function nulls. These particular data were selected to illustrate the characteristics of the system. Recall that one can easily make the noise floor 40 dB lower using the rf bridge shown in Fig. 18.

C. DUAL DELAY-LINE DISCRIMINATOR

1. Phase Noise Measurements

The dual delay-line discriminator is shown in Fig. 24. This system was suggested by Halford (1975) as a technique for lowering the noise floor of the delay-line phase noise measurement system. The system consists of

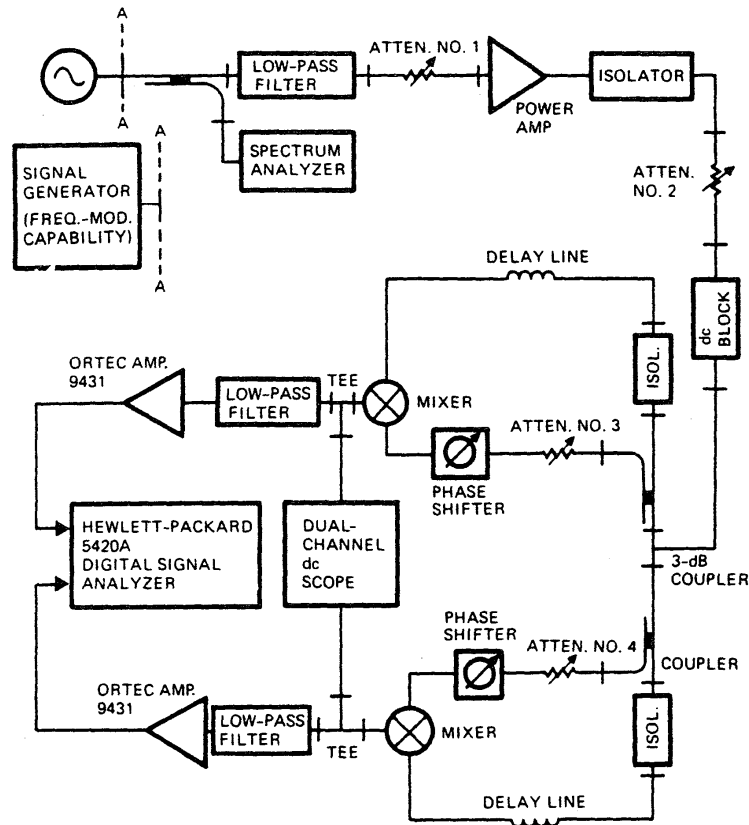


FIG. 24 A dual delay-line phase noise measurement system. (Courtesy Instrument Society of America.)

two differential delay-line systems. The single-oscillator signal is applied to both systems and cross-spectrum analysis is performed on the signal output from the two delay-line systems. Signal processing is performed with the Hewlett-Packard 5420A digital signal analyzer. The *cross spectrum* is obtained by taking the product of the Fourier transform of one signal and the complex conjugate of the Fourier transform of a second signal. It is a phase-sensitive characteristic resulting in a complex product that serves as a measurement of the relative phase of two signals. Cross spectrum gives a phase- and amplitude-sensitive measurement directly. By performing the product $S_y(f) \cdot S_x(f)^*$, a certain signal-to-noise enhancement is achieved.

The low-noise amplifiers preceding the digital signal analyzer are used when performing measurements at Fourier frequencies from 1 Hz to 25 kHz. The amplifiers are not used when performing measurements below the Fourier frequency of 1 Hz.

2. Calibrating the Dual Delay-Line System

Each delay line in the system is calibrated separately following the same basic procedure set forth in Section IV.B. The Hewlett-Packard 5420 measures the one-sided spectral density of frequency fluctuations in hertz squared per hertz. The spectral density of phase fluctuation in radians squared per hertz can be calculated as

$$S_{\delta\phi}(f) = S_{\delta\nu}(f)/f^2, \quad (99)$$

and

$$\mathcal{L}(f) = S_{\delta\nu}(f)/2f^2, \quad (100)$$

per hertz. The Hewlett-Packard 5420 measurement of $S_{\delta\nu}(f)$ in Hz^2/Hz must, therefore, be corrected by $1/2f^2$. However, the f^2 correction must be entered in terms of radian frequency ($\omega = 2\pi f$). This conversion is accomplished by

$$\mathcal{L}(f) = S_{\delta\nu}(f)(1/2f^2)(4\pi^2/4\pi^2) = [2\pi^2 S_{\delta\nu}(f)]/(\omega)^2 \quad (101)$$

per hertz since Eq. (100) can be stated in the following terms:

$$[2\pi^2 S_{\delta\nu}(f)]/4\pi^2 f^2.$$

Signal-to-noise enhancement greater than 20 dB has been obtained using the dual-channel delay-line system.

D. MILLIMETER-WAVE PHASE-NOISE MEASUREMENTS

1. Spectral Density of Phase Fluctuations

The delay line used as an FM discriminator is based, in principle, on a *nondispersive* delay line. However, a waveguide can be used as the delay line because the Fourier frequency range of interest is a small percentage of the operating bandwidth (seldom over 100 MHz), and the dispersion can be considered negligible.

The calibration and measurement are performed as set forth in Section IV.B. The modulation index m is usually established using the carrier-to-sideband ratio that uses amplitude modulation because millimeter sources are either unstable or cannot be modulated. The two approaches to measurements at millimeter frequencies are shown in Figs. 25 and 26. Figure 25 shows the *direct measurement* using a waveguide delay line. This system offers improved sensitivity if adequate input power is available. The rf bridge and delay-line portion of the system differs from Fig. 18 because pre- and

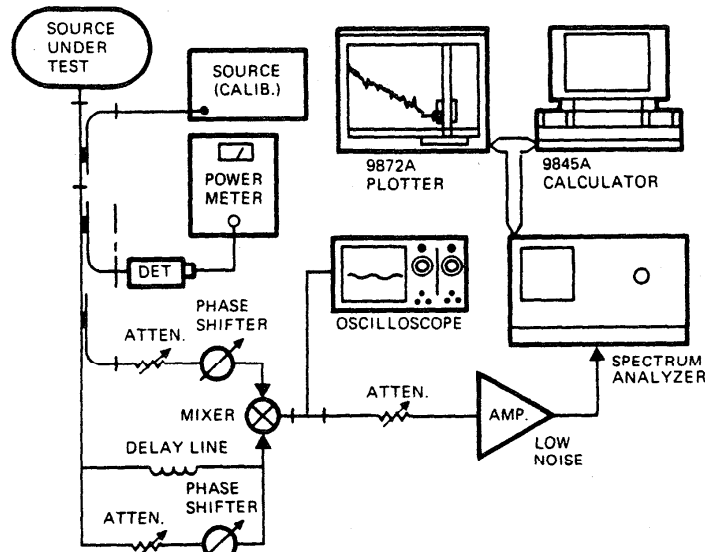


FIG. 25 Millimeter-wave phase noise measurements using a waveguide delay line. (From Seal and Lance, 1981.)

post-bridge amplifiers with appropriate gain are not available, so the sensitivity can equal the amount of carrier suppression.

Figure 26 shows the use of a harmonic mixer to downconvert to the convenient lower frequency where post-bridge amplifiers are available. The relatively low sensitivity to frequency drift that is characteristic of delay-line discriminators becomes an advantage here. A separate calibration generator is required, as shown in Fig. 25, and a power meter is used to assure proper power levels during the calibration process.

2. SPECTRAL DENSITY OF AMPLITUDE FLUCTUATIONS

AM noise measurements require equal electrical length in the two channels that supply the signals to the mixer. The delay line must be replaced with the necessary length of transmission line to establish the equal-length condition when the systems shown in Figs. 25 and 26 are used. The AM noise measurement system is calibrated and the noise measurements are performed directly in units of power for a direct measurement of $m(f)$ in dBc/Hz. $m(f)$ is the spectral density of one modulation sideband divided by the total signal power at a Fourier frequency difference f from the signal's average frequency ν_0 . The system calibration establishes the detection characteristics in terms of total power output at the IF port of the mixer (detector).

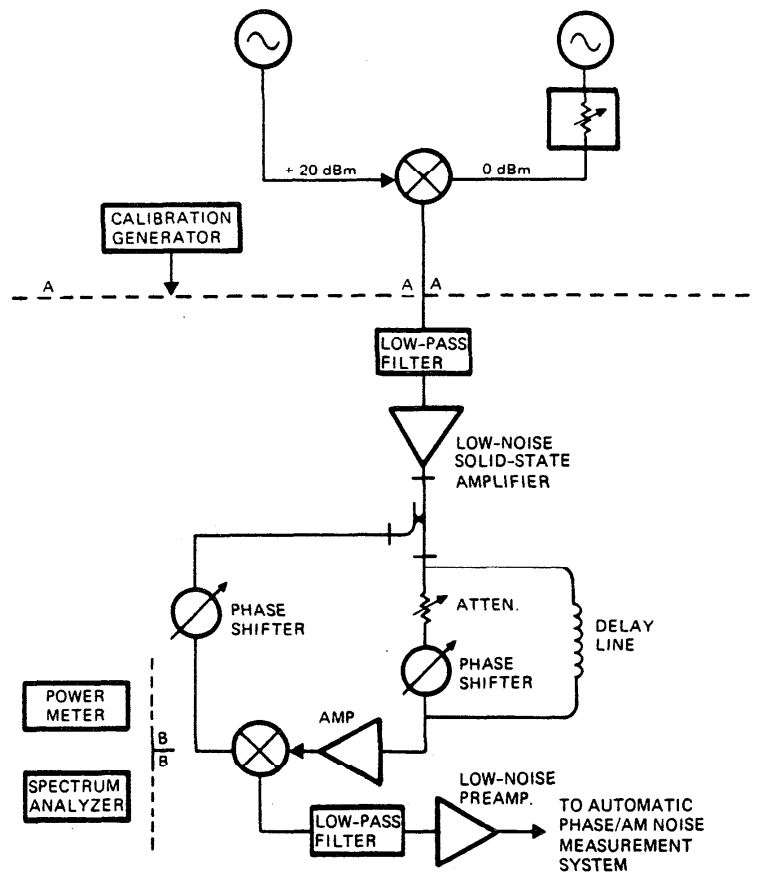


FIG. 26 A millimeter-wave hybrid phase noise measurement system that produces IF frequency and uses a delay-line discriminator at IF frequency. (From Seal and Lance, 1981.)

The AM noise measurements are performed according to the following.

(1) A known AM modulation (carrier-sideband ratio) must be established to calibrate this detector in terms of total power output at the IF port. The modulation must be low enough so that the sidebands are at least 20 dB below the carrier. This is to keep the total added power due to the modulation small enough to cause an insignificant change in the detector characteristics.

(2) The rf power levels are adjusted for levels of approximately 10 dBm at the reference port and 0 dBm at the test port of the mixer.

(3) Approximately 40 dB is set in the precision IF attenuator. The system is adjusted for an *out-of-phase quadrature condition*.

(4) The modulation frequency and power level are measured by the automatic baseband spectrum analyzer. The total carrier-power reference level is measured power, plus the carrier-sideband modulation ratio, plus the IF attenuator setting.

(5) The AM modulation is removed, the IF attenuator set to 0 dB, and the system re-checked to verify the out-of-phase quadrature (maximum dc output from the mixer IF port). Noise (V_n) is measured at the selected Fourier frequencies. A direct calculation of $m(f)$ in dBc/Hz is

$$m(f) = [(\text{modulation power (dBm)} + \text{carrier-sideband ratio (dB)} \\ + \text{IF attenuation (dB)} - \text{noise power (dBm)} + 2.5 \text{ dB} \\ - 10 \log(\text{BW})]. \quad (102)$$

Figure 27 illustrates the measurements of AM and phase noise of two GUNN oscillators that were offset in frequency by 1 GHz. The measurements were performed using the coaxial delay-line system.

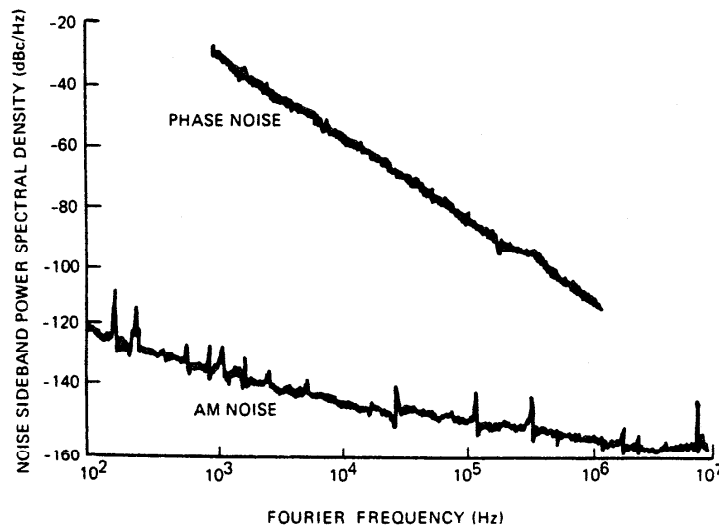


FIG. 27 Phase noise and AM noise of 40- and 41-GHz Gunn oscillators. (From Seal and Lance, 1981.)

V. Conclusion

The fundamentals and techniques for measurement of phase noise have been set forth for two basic systems. The two-oscillator technique provides the capability for measuring high-performance cw sources. The system sensitivity is superior to the single-oscillator technique for measuring phase noise very close to the carrier.

High-stability sources such as those used in frequency standards applications can be measured without using phase-locked loops. However, most microwave sources exhibit frequency instability that requires phase-locked loops to maintain the necessary quadrature conditions. The characteristics of the phase-locked loops must be evaluated to obtain the source phase noise characteristics. Also, in principle, one must have three sources at the same frequency to characterize a given source. If three sources are not available, one must assume that either one source is superior in performance or that they have equal phase noise contributions.

The single-oscillator technique employing the delay line as an FM discriminator has adequate sensitivity for measuring most microwave sources. The economic advantages of using this system include the fact that only one source is required, phase-locked loops are not required, system configuration is relatively inexpensive, and the system is inherently insensitive to oscillator frequency drift.

The single-oscillator technique using the delay-line discriminator can be adapted to measure the phase noise of pulsed sources. Pulsed sources have been measured at 94 GHz by F. Labaar at TRW, Redondo Beach, California.

ACKNOWLEDGMENTS

Our initial preparations for developing a phase noise measurement capability were the result of discussions with Dr. Jorg Rauc of TRW. Our first phase noise development effort was assisting in the evaluation of phase noise measurement systems designed and developed by Bill Hook of TRW (Hook, 1973). The efforts of Don Leavitt of TRW were vital in initiating the measurement program.

We are very grateful to Dr. Donald Halford of the National Bureau of Standards in Boulder, Colorado, for his interest, consultations, and valuable suggestions during the development of the phase noise measurement systems at TRW.

We appreciate the measurement cross-checks performed by C. Reynolds, J. Oliverio, and H. Cole of the Hewlett-Packard Company, Dr. J. Robert Ashley of the University of Colorado, Colorado Springs, and G. Rast of the U.S. Army Missile Command, Redstone Arsenal, Huntsville, Alabama.

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PERFORMANCE OF AN AUTOMATED HIGH ACCURACY
PHASE MEASUREMENT SYSTEM

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Summary

A fully automated measurement system has been developed that combines many properties previously realized with separate techniques. This system is an extension of the dual mixer time difference technique, and maintains its important features: zero dead time, absolute phase difference measurement, very high precision, the ability to measure oscillators of equal frequency and the ability to make measurements at the time of the operator's choice. For one set of design parameters, the theoretical resolution is 0.2 ps, the measurement noise is 2 ps rms and measurements may be made within 0.1 s of any selected time. The dual mixer technique has been extended by adding scalars which remove the cycle ambiguity experienced in previous realizations. In this respect, the system functions like a divider plus clock, storing the epoch of each device under test in hardware.

The automation is based on the ANSI/IEEE-583 (CAMAC) interface standard. Each measurement channel consists of a mixer, zero-crossing detector, scalar and time interval counter. Four channels fit in a double width CAMAC module which in turn is installed in a standard CAMAC crate. Controllers are available to interface with a wide variety of computers as well as any IEEE-488 compatible device. Two systems have been in operation for several months. One operates 24 hours a day, taking data from 15 clocks for the NBS time scale, and the other is used for short duration laboratory experiments.

Review of the Dual Mixer
Time Difference Technique

It is advantageous to measure time directly rather than time fluctuations, frequency or frequency fluctuations. These measurements constitute a hierarchy in which the subsequently listed quantities may always be calculated from the previous ones. However, the reverse is not true when there are gaps in the measurements. In the past, frequency was usually not derived from time measurements for short sample times because time interval measurements could not be performed with adequate precision. The dual

mixer technique, illustrated in Figure 1, made it possible to realize the precision of the beat frequency technique in time interval measurements.

The signals from two oscillators (clocks) are applied to two ports of a pair of double balanced mixers. Another signal synthesized from one of the oscillators is applied to the remaining two ports of the mixer pair. The input signals may be represented in the usual fashion

$$V_1(t) = V_{10} \sin [2\pi\nu_{10}t + \phi_1(t)],$$

$$V_2(t) = V_{20} \sin [2\pi\nu_{20}t + \phi_2(t)] \text{ and}$$

$$V_s(t) = V_{s0} \cos [2\pi\nu_{s0}t + \phi_s(t)]$$

where $\nu_s = \nu_{10}(1-1/R)$ and R is a constant usually called the heterodyne factor.

The low passed outputs of the two mixers are

$$V_{B1} = V_{B10} \sin [\phi_1(t) - \phi_s(t)] \text{ and}$$

$$V_{B2} = V_{B20} \sin [\phi_2(t) - \phi_s(t)] \text{ where}$$

$$\phi(t) = 2\pi\nu_0 t + \phi(t).$$

The time interval counter starts at time t_M when V_{B1} crosses zero in the positive direction and stops at time t_N , the time of the very next positive zero crossing of V_{B2} . Thus

$$\phi_1(t_M) - \phi_s(t_M) = 2M\pi \text{ and}$$

$$\phi_2(t_N) - \phi_s(t_N) = 2N\pi \text{ where}$$

N and M are integers.

Subtracting the two equations in order to compare the phases of oscillators 1 and 2, one obtains

$$\phi_2(t_N) - \phi_1(t_M) = \phi_s(t_N) - \phi_s(t_M) + 2(N-M)\pi.$$

The phase of an oscillator at time t_N may be written in terms of its phase at t_M and its

average frequency over the interval $t_M < t_N$.

$$\phi(t_N) = \phi(t_M) + 2\pi[\bar{v}(t_M; t_N)](t_N - t_M) \text{ and}$$

when we apply this equation to both ϕ_2 and ϕ_5 we find

$$\phi_2(t_M) - \phi_1(t_M) = 2(N-M)\pi - 2\pi[\bar{v}_{B2}(t_M; t_N)](t_N - t_M)$$

where $v_{B2} = v_2 - v_5$.

Since M and N are not measurable with the equipment in Figure 1, the dual mixer technique has heretofore only been used to measure the phase difference between two oscillators modulo 2π . We denote the period of the time interval counter time base by τ_c and the number of counts recorded in a measurement by P. Then the phase difference between the two oscillators is given by

$$[\phi_2(t_M) - \phi_1(t_M)] \text{ mod } 2\pi = -2\pi[\bar{v}_{B2}(t_M; t_N)]\tau_c P$$

Figure 2 illustrates the output of the measurement system over a period of time. If a measurement begins and ends without the time interval counter making a transition between zero and its maximum value, e.g., $t_a < t_M < t_N < t_b$, then the phase difference can be calculated from the data. If $t_a < t_M < t_b < t_N < t_c$, then the data must be corrected by 2π to calculate the phase difference. Experience has shown that there are many measurement situations for which the number of transitions of the time interval counter which occur between t_M and t_N cannot be known. For this reason, a modification has been developed which removes the ambiguity by measuring M and N.

Extended Dual Mixer Time Difference Measurement Technique

In order to configure the system to acquire complete phase information, two scalars are added to count the zero crossings of each mixer. Figure 3 is the block diagram of a two channel system. It is constructed from identical circuit modules and therefore contains an unused time interval counter. However, this design permits very straightforward and inexpensive extension to the comparison of an arbitrarily large number of oscillators with no need for switching any signals.

The counter outputs are combined to form the phase difference between oscillators.

$$\phi_2(t_M) - \phi_1(t_M) = 2(N_0 - M_0)\pi + 2(N-M)\pi - 2\pi[\bar{v}_{B2}(t_M; t_N)]\tau_c P$$

The first term is a constant which represents the choice of the time origin and can be ignored. The last two terms and their sum are plotted in Figure 4.

The average beat frequency $\bar{v}_{B2}(t_M; t_N)$ cannot be known exactly. However, it may be estimated with sufficient precision from the previous pair of measurements designated ' and ". The average frequency is approximately

$$\bar{v}_{B2}(t_M; t_N) \cong (N'' - N') / [R(M'' - M') / v_{10} + \tau_c (P'' - P')]$$

provided that it changes sufficiently slowly compared to the interval $t_M < t_N$. A typical value for this error will be given in the following section.

Hardware Implementation

All measurement channels consist of a mixer, zero-crossing detector, scaler and time interval counter. Four such circuits can be built in a double width CAMAC module. The system is easily expanded to compare many oscillators and a complete system for making phase comparisons among four clocks is shown in Figure 5. We have chosen parameters which are reasonable for comparing state-of-the-art atomic standards. Thus, the synthesizer is offset 10Hz below oscillator # 1 and $R = 5 \times 10^6$. The outputs from both mixers are approximately 10Hz. The noise bandwidth is 100 Hz. The time interval counter is twice the frequency of oscillator #1 or approximately 10 MHz. The quantization error is $1/2R = 10^{-6}$ cycle or 0.2ps which is a factor of ten smaller than the measurement noise. As stated earlier, an error will result from frequency changes which violate the constancy assumption used to estimate v_{B2} . A change in v_2 by 10^{-10} during the interval between two measurements will result in a time deviation error of 10ps. Thus, one must make more closely spaced measurements for oscillators which have large dynamic frequency changes than for more stable devices. Two other sources of inaccuracy are the sensitivities to the amplitude and phase of the common oscillator. Figure 6 shows the measured value of $x = \phi/2\pi v_0$ as a function of the amplitude of the input signal and the phase of the synthesizer.

The new measurement system has many desirable features and properties:

- (1) It has very high resolution, limited by the internal counters to 0.2 ps and by noise to approximately 2 ps.
- (2) It has much lower noise than divider based measurement systems. However compromises made to achieve low cost, low power, small size and automatic operation degrade the performance compared to state-of-the-art systems for comparing 2 oscillators.
- (3) The operation is fully automatic.
- (4) NBS has developed a detailed operating manual for the equipment and software.

- (5) All oscillators in the range of 5 MHz \pm 5 Hz may be compared. Other carrier frequencies such as 1 MHz, 5.115 MHz, 10 MHz and 10.23 MHz are also usable. However, different carrier frequencies may not be mixed on the same system. The system has been successfully tested with an oscillator offset 4.6 Hz from nominal 5MHz. Measurements were made at intervals of 2 hours between which the system had to accumulate approximately $2 \times 10^6 \pi$. The system has also been tested with an oscillator offset 4×10^5 , and no errors were detected during a period of 40 days.
- (6) All sampling times in the range of 1 second to 16 days with a resolution of 0.1 second are possible. Measurements may be made on command or in a preprogrammed sequence.
- (7) Measurements are synchronized precisely, i.e. at the picosecond level, with the reference clock. They may therefore be synchronized with important user system events, such as the switching times of a FSK or PSK system.
- (8) All oscillators are compared synchronously and all measurements are performed within a maximum interval of 0.1 second. As a result, the phase of any oscillator needs to be interpolated to the chosen measurement time for an interval of 0.1 second maximum. This capability, which is not present in either single heterodyne measurement systems or switched measurement systems eliminates a source of "measurement" error which is generally much larger than the noise induced errors. For example, interpolation of the phase of a high performance Cs clock ($\sigma_y \sim 10^{-11}/\tau^2$) over a period of 3 hours would produce approximately 1.5 ns phase uncertainty. To maintain 4 ps accuracy requires measurements simultaneous to 0.1s.
- (9) There are no phase errors due to the switching of rf signals since there is no switching anywhere in the analog measurement system.
- (10) No appreciable phase errors are introduced when it is necessary to change the reference clock since, as shown in Figure 6, the peak error due to changes in synthesizer phase is 20 ps.
- (11) The measurement system is capable of measuring its own phase noise when the same signal is applied to two input ports. Figure 7 shows the phase deviations between two such channels over a period of 75,000 seconds and Figure 8 is the corresponding Allan variance plot. Figure 9 shows the phase deviations between 2 input channels over a period of 40 days.
- (12) Since the IEEE-583 (CAMAC) interface standard has been followed for all the custom

hardware, the system may be easily interfaced to almost any instrument controller. NBS has already tested the system using a large minicomputer, a small minicomputer and a desk top calculator. Interfaces between IEEE-583 and IEEE-488 controllers are available and have been used successfully.

- (13) The system is capable of comparing a very large number of oscillators at a reasonable cost per device.

There are also disadvantages to this measurement system. The most important are:

- (1) The complexity of the hardware is greater than for some systems. It is possible that this will reduce reliability.
- (2) A high level of redundancy is difficult to achieve. The system design stresses size, power, convenience and cost, resulting in an increase in the number of possible single point failure mechanisms compared to some other techniques. For example, a CAMAC power supply failure will result in a loss of data for all devices being measured.
- (3) A substantial commitment is required in both specialized hardware and software.
- (4) If an oscillator under test experiences a phase jump which exceeds 1 cycle, the measurement system records a jump with incorrect absolute magnitude. As a result, it may not be applicable to signals which are frequency modulated with discontinuous phase steps larger than 2π .

Conclusions

We have demonstrated a new phase measurement system with very desirable properties: All oscillators in the range of 5MHz \pm 5Hz may be measured directly. The sampling times are only restricted by the requirement that they exceed one second. The noise floor is $\sigma_y(2,\tau) = 3 \times 10^{-12}/\tau$ in short term and the time deviations are less than 100 ps. All circuitry is designed as modules which allows expansion at modest cost. Compatibility with a variety of computers is insured through the use of the IEEE-583 interface and adapters are available to permit use with an IEEE-488 controller. The system makes it feasible to make completely automated phase measurements at predetermined times on large numbers of atomic clocks. It's own noise is one-hundred times less than the state-of-the-art in clock performance. It will be used in the near future to make all measurement needed to compute NBS atomic time, but it will also be very valuable for any laboratory which uses three or more atomic clocks.

References

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and frequency stability of precision oscillators," NBS Tech. Note 669 (1975).

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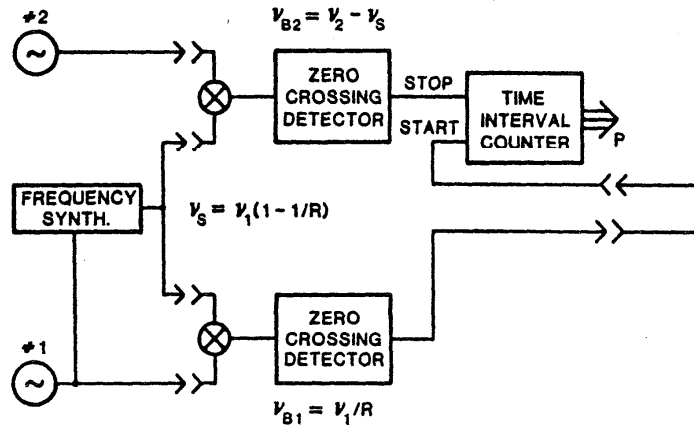


Figure 1. Dual Mixer Time Difference Measurement System

$$[\phi_2(t_H) - \phi_1(t_H)] \bmod 2\pi = -2\pi[\bar{v}_{B2}(t_H; t_H)]r_c P$$

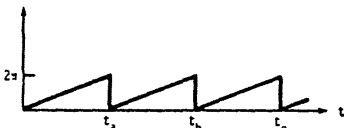


Figure 2. Dual Mixer Data

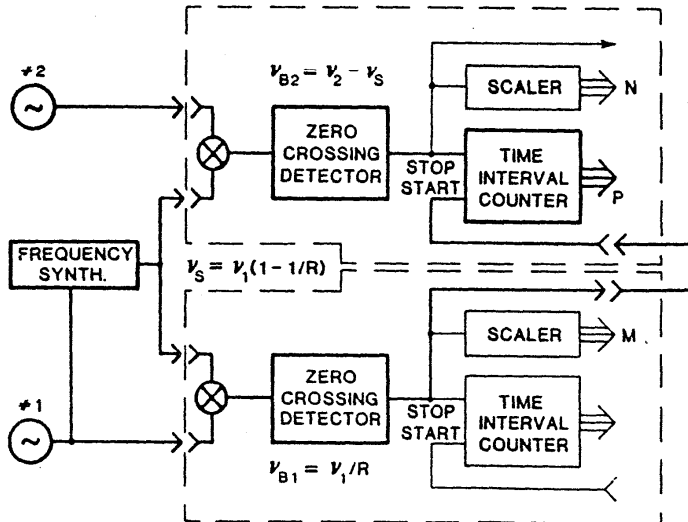


Figure 3. Extended Dual Mixer Time Difference Measurement System

$$\phi_2(t_N) - \phi_1(t_N) = 2(N_0 - N_0)\pi + 2(N-M)\pi - 2\pi[\bar{v}_{B2}(t_N; t_N)]\tau_c P$$

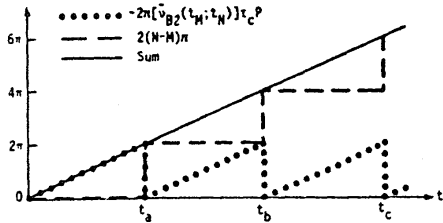


Figure 4. Extended Dual Mixer Data

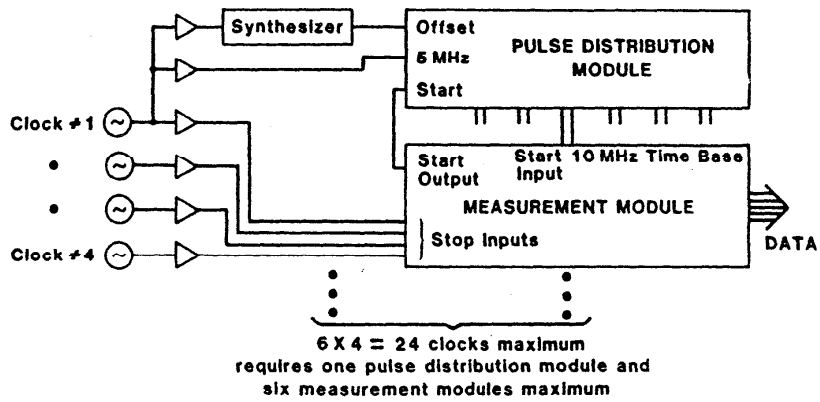


Figure 5. System Block Diagram

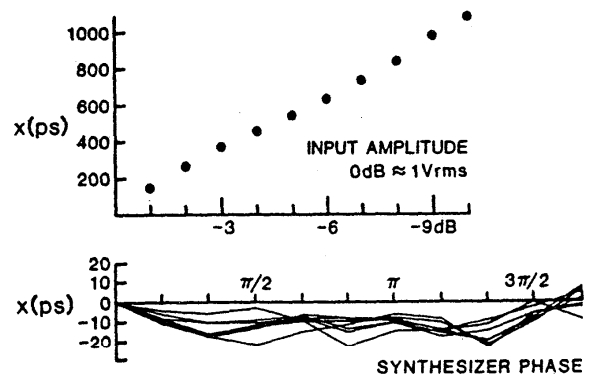
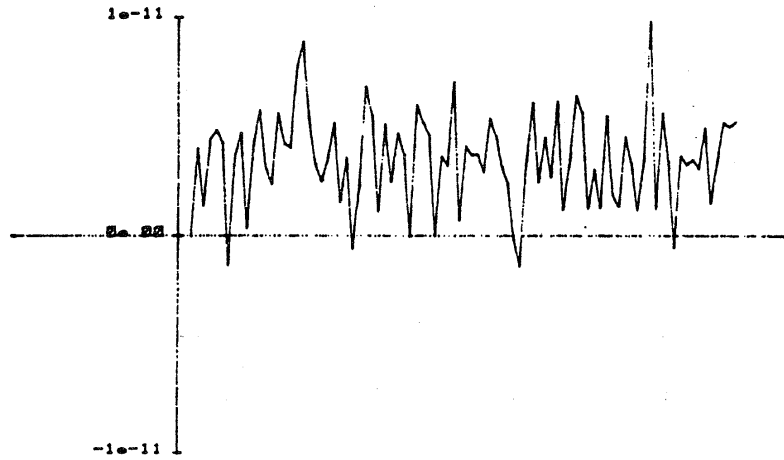


Figure 6. Measured Time Difference vs. Input Amplitude and Synthesizer Phase

PHASE PLOT Clock No. 4-8 nbe4a - nbe4b
 Let Sqr Slope of $-1.959200e-17/S$ Removed E)W 05 may 82



T= 74700 SECONDS

Figure 7. Raw Phase Data for Two Channels Driven from the Same Source

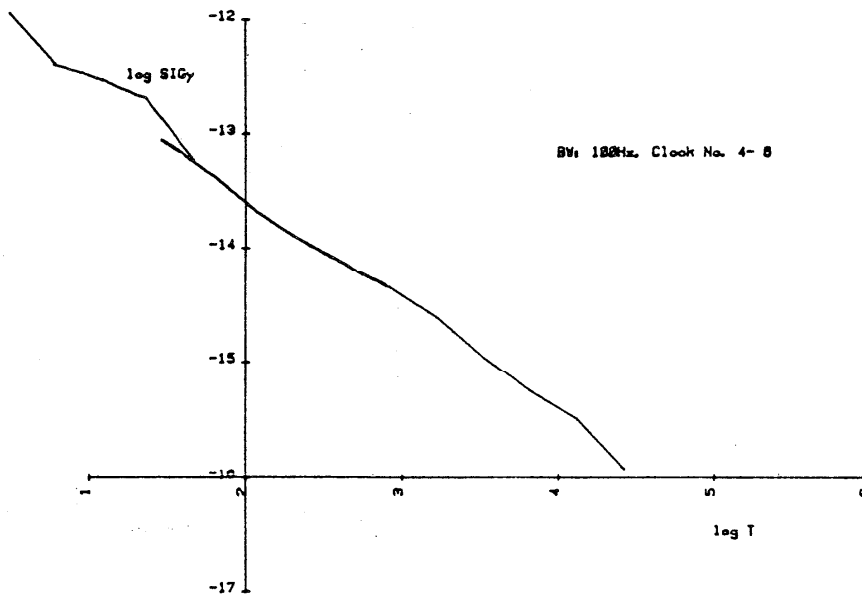


Figure 8. Noise Floor of Measurement System

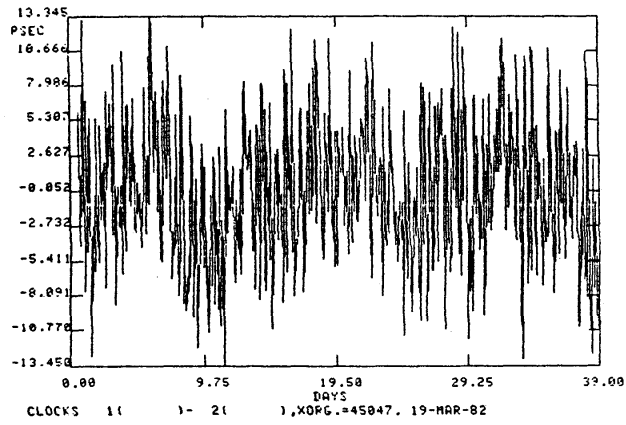


Figure 9. Raw Phase Data for Two Channels Driven from the Same Source

BIASES AND VARIANCES OF SEVERAL FFT SPECTRAL ESTIMATORS
AS A FUNCTION OF NOISE TYPE AND NUMBER OF SAMPLES

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Abstract

We theoretically and experimentally investigate the biases and the variances of Fast Fourier transform (FFT) spectral estimates with different windows (data tapers) when used to analyze power-law noise types f^0 , f^{-2} , f^{-3} and f^{-4} . There is a wide body of literature for white noise but virtually no investigation of biases and variances of spectral estimates for power-law noise spectra commonly seen in oscillators, amplifiers, mixers, etc. Biases (errors) in some cases exceed 30 dB. The experimental techniques introduced here permit one to analyze the performance of virtually any window for any power-law noise. This makes it possible to determine the level of a particular noise type to a specified statistical accuracy for a particular window.

I. Introduction

Fast Fourier transform (FFT) spectrum analyzers are very commonly used to estimate the spectral density of noise. These instruments often have several different windows (data tapers) available for analyzing different types of spectra. For example, in some applications spectral resolution is important; in others, the precise amplitude of a widely resolved line is important; and in still other applications, noise analysis is important. These diverse applications require different types of windows.

We theoretically and experimentally investigate the biases and variances of FFT spectral estimates with different windows when used to analyze a number of common power-law noise types. There is a wide body of literature for white noise but virtually no investigation of these effects for the types of power-law noise spectra commonly seen in oscillators, amplifiers, mixers, etc. Specifically, we present theoretical results for the biases associated with two common windows — the uniform and Hanning windows — when applied to power-law spectra varying as f^0 , f^{-2} and f^{-4} . We then introduce experimental techniques for accurately determining the biases of any window and use them to evaluate the biases of three different windows for power-law spectra varying as f^0 , f^{-2} , f^{-3} and f^{-4} . As an example we find with f^{-4} noise that the uniform window can have errors ranging from a few dB to over 30 dB, depending on the length of span of the f^{-4} noise.

We have also theoretically investigated the variances of FFT spectral estimates with the uniform and Hanning windows (confidence of the estimates) as a function of the power-law noise type and as a function of the amount of data. We introduce experimental techniques that make it relatively easy to independently determine the variance of the spectral estimate for virtually any window on any FFT spectrum analyzer. The variance that is realized on a particular instrument depends not only on the window but on the specific implementation in both hardware and software. We find that the variance of the spectral density estimates for white noise, f^0 , is very similar for three specific windows available on one instrument and almost

identical to that obtained by standard statistical analysis. The variances for spectral density estimates of f^{-4} noise are only 4% higher than that of f^0 noise for two of the windows studied. The third window — the uniform window — does not yield usable results for either f^{-3} or f^{-4} noise.

Based on this work it is now possible to determine the minimum number of samples necessary to determine the level of a particular noise type to a specified statistical accuracy as a function of the window. To our knowledge this was previously possible only for white noise — although the traditional results are generally valid for noise that varied as $f^{-\beta}$, where β was equal to or less than 4.

II. Spectrum Analyzer Basics

The spectrum analyzer which was used in the experimental work reported here is fairly typical of a number of such instruments currently available from various manufacturers. The basic measurement process generally consists of taking a string of $N_s = 1024$ digital samples of the input wave form, which we represent here by X_1, X_2, \dots, X_{N_s} . The basic measurement period was 4 ms. This yields a sampling time $\Delta t = 3.90625 \mu s$. Associated with the FFT of a time series with N_s data points, there are usually $(N_s/2) + 1 = 513$ frequencies

$$f_j = \frac{j}{N_s \Delta t}, \quad j = 0, 1, \dots, N_s/2.$$

The fundamental frequency f_1 is 250 Hz, and the Nyquist frequency $f_{N_s/2}$ is 128 kHz. Since the spectrum analyzer uses an anti-aliasing filter which significantly distorts the high frequency portion of the spectrum, the instrument only displays the measured spectrum for the lowest 400 nonzero frequencies, namely, $f_1 = 250$ Hz, $f_2 = 500$ Hz, \dots , $f_{400} = 100$ kHz.

The exact details of how the spectrum analyzer estimates the spectrum for X_1, \dots, X_{N_s} are unfortunately not provided in the documentation supplied by the manufacturer, so the following must be regarded only as a reasonable guess on our part as to its operation (see [1] for a good discussion on the basic ideas behind a spectrum analyzer; two good general references for spectral analysis are [2] and [4]). The sample mean,

$$\bar{X} \equiv \frac{1}{N_s} \sum_{t=1}^{N_s} X_t,$$

is subtracted from each of the samples, and each of these "de-measured" samples is multiplied by a window h_t (sometimes called a data taper) to produce

$$X_t^{(h)} = h_t (X_t - \bar{X}).$$

The spectral estimate,

$$\hat{S}_1(f_j) = \Delta t \left| \sum_{t=1}^{N_s} X_t^{(h)} e^{-i2\pi f_j t \Delta t} \right|^2, \quad j = 0.1, \dots, N_s/2.$$

is then computed using an FFT algorithm.

The subscript "1" on $\hat{S}_1(f_j)$ indicates that this is the spectral estimate formed from the first block of N_s samples. A similar spectral estimate $\hat{S}_2(f_j)$ is then formed from the second block of contiguous data $X_{N_s+1}, X_{N_s+2}, \dots, X_{2N_s}$. In all, there are N_b different spectral estimates from N_b contiguous blocks, and the spectrum analyzer averages these together to form

$$\hat{S}(f_j) = \frac{1}{N_b} \sum_{k=1}^{N_b} \hat{S}_k(f_j). \quad (1)$$

It is the statistical properties of $\hat{S}(f_j)$ with which we are concerned in this paper.

Unfortunately some important aspects of the windows are not provided in the documentation for the instrument. One important detail is the manner in which the window is normalized. There are two common normalizations:

$$\sum_{i=1}^{N_s} [h_i(X_i - \bar{X})]^2 = \frac{1}{N_s} \sum_{i=1}^{N_s} (X_i - \bar{X})^2$$

and

$$\sum_{i=1}^{N_s} h_i^2 = 1. \quad (2)$$

The first of these is common in engineering applications because it ensures that the power in the windowed samples $X_i^{(A)}$ is the same as in the original demeaned samples; the second is equivalent to the first in expectation and is computationally more convenient, but it can result in small discrepancies in power levels. Either normalization affects only the level of the spectral estimate and not its shape.

There are three windows built into the spectrum analyzer used here. The first is the uniform (rectangular, default) window $h_i^{(U)} = 1/\sqrt{N_s}$. The second is the Hanning data window, for which there are several slightly different definitions in the literature. In lieu of specific details, we assume the following symmetric definition:

$$h_i^{(H)} = C^{(H)} \left(1 - \cos \frac{2\pi(t-0.5)}{N_s} \right), \quad 1 \leq t \leq \frac{N_s}{2}, \\ = h_{N_s-t+1}^{(H)}, \quad \frac{N_s}{2} + 1 \leq t \leq N_s;$$

where $C^{(H)}$ is a constant which forces the normalization in Equation (2). The third window is a proprietary "flattened peak" window, about which little specific information is available (it is evidently designed to accurately measure the heights of peaks in a spectrum).

III. Expected Value and Bias of Spectral Estimates

III.A. Theoretical Analysis

We need to assume a noise model for the X_t 's in order to determine the statistical properties of $\hat{S}(f_j)$ in Equation (1). We consider three different models, each of which is represented in terms of a Gaussian white noise process ϵ_t with mean zero and variance σ_t^2 . The second-order properties of each model are given by a spectral density function $S(\cdot)$ defined over the interval $[-1/(2\Delta t), 1/(2\Delta t)]$ in cycles/ Δt . The first model is a discrete parameter, white noise process (f^0 noise):

$$X_t = \epsilon_t \quad \text{and} \quad S(f) = \sigma_t^2 \Delta t.$$

The second model is a discrete-parameter, random-walk process (nominally f^{-2} noise):

$$X_t = \sum_{s=1}^t \epsilon_s \quad \text{and} \quad S(f) = \frac{\sigma_t^2 \Delta t}{4 \sin^2(\pi f \Delta t)}.$$

The third model is a discrete-parameter, random-run process (nominally f^{-4} noise):

$$X_t = \sum_{r=1}^t \sum_{s=1}^r \epsilon_s \quad \text{and} \quad S(f) = \frac{\sigma_t^2 \Delta t}{16 \sin^4(\pi f \Delta t)}.$$

Continuous parameter versions of these three models have been used extensively in the literature as models for noise commonly seen in oscillators.

For each of the three models we have derived expressions for $E\{\hat{S}(f)\}$, the expected value of $\hat{S}(f)$. These expressions depend on the window h_t , the number of samples N_s in each block and — in the case of a random-run process — the number of blocks N_b . The details behind these calculations will be reported elsewhere [3]; here we merely summarize our conclusions for the three models in combination with the uniform and Hanning windows and $N_s = 1024$.

First, for a white noise process,

$$E\{\hat{S}(f_j)\} = S(f_j), \quad j = 1, 2, \dots, 512,$$

when the uniform window is used. For the Hanning window, the above equality also holds to a very good approximation for $2 \leq j \leq 511$ and to within 0.8 dB for $j = 1$ and 512 (the latter is of no practical importance since the highest frequency index given by the spectrum analyzer is $j = 400$). These theoretical calculations agree with our experimental data except at f_1 (see Table 1).

Second, for a random-walk process,

$$E\{\hat{S}(f_j)\} = 2S(f_j), \quad j = 1, 2, \dots, 512,$$

when the uniform window is used, i.e., the expected value is twice what it should be at all frequencies. This theoretical result has been verified by Monte Carlo simulations, but it does not agree with our experimental data, which shows no significant level shift in the estimated spectrum. The source of this discrepancy is currently under investigation, but it may be due to either (a) factors in the experimental data which effectively make it band-limited, random-walk noise, i.e., its spectral shape is markedly different from f^{-2} for, say, $0 < f < f_1$ or (b) an incorrect guess on our part as to how the spectral estimate is normalized by the spectrum analyzer. For the Hanning window, we found that

$$E\{\hat{S}(f_j)\} = \begin{cases} 1.08S(f_j) & j = 1; \\ 1.48S(f_j) & j = 2; \\ 1.15S(f_j) & j = 3; \\ 1.07S(f_j) & j = 4; \\ 1.04S(f_j) & j = 5; \\ S(f_j) & 6 \leq j \leq 511 \text{ to within } 3\%. \end{cases}$$

i.e., $\hat{S}(f_j)$ is essentially an unbiased spectral estimate except for the lowest few frequencies. This theoretical result has been verified by Monte Carlo simulations and also agrees in general with our experimental data.

Third, for a random-run process,

$$E\{\hat{S}(f_j)\} = C N_s f^{-2}, \quad 1 \leq j \leq 400.$$

to a good approximation when the uniform window is used, where C_{N_b} is a constant which depends on the number of blocks N_b and increases as N_b increases. Thus the shape of $E\{\hat{S}(f_j)\}$ follows that of a random-walk process (f^{-2}) rather than that of a random-run process (f^{-4}). This shape has been verified experimentally (see the next subsection), but the dependence of the level on N_b has not. The increase in level of $E\{\hat{S}(f_j)\}$ as N_b increases is due to the fact that the expected value of the sample variance of a block of N_s samples increases with time — by contrast, it is constant with time for the white noise and random-walk cases. For the Hanning window, we found that

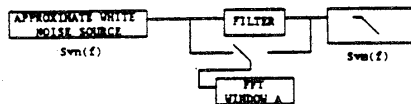
$$E\{\hat{S}(f_j)\} = C'_{N_b} S(f_j), \quad 4 \leq j \leq 400,$$

to a good approximation, where again C'_{N_b} is a constant — different from C_{N_b} — which depends on the number of blocks N_b and increases as N_b increases. For frequencies less than f_4 the theoretical results indicate significant (greater than 4%) distortion in the shape, but these do not agree in detail with the experimental values reported in Table 1. For $f_j \geq f_4$ the shape has been verified experimentally, but the dependence of the level on N_b has not. The discrepancy in level between the theoretical and experimental results is yet to be resolved, but it is probably due to a mismatch between the assumed random-run model and the true spectrum for the data (possibly band-limited random-run).

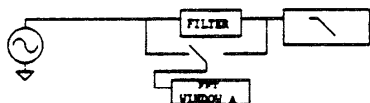
1) Verify Spectral Density Function & Voltage Reference



2) Measure Spectrum Relative to Noise Source Requires Multiple Scans



3) Measure Filter Transfer Function h(f)



4) Calculate Biases

$$B(f) = S_w(f) - h^2(f) S_{wn}(f)$$

Figure 1. Outline of measurement procedure for determining the biases in spectral estimators.

III.B. Experimental Determination

The following procedure can be used to experimentally determine the bias in the spectral estimate of any noise spectrum using any window in a particular instrument. The basic concept is to implement a filter that, when applied to white noise, mimics the approximate noise spectra of interest and then measures the level of the white noise and the filter transfer function in a way which has high precision and accuracy as illustrated in Figure 1. First, the level of a known white noise is measured over a convenient range. The higher the frequency span the faster that this is accomplished. Obviously, the chosen range must be one over which the noise source is accurate. To obtain a

precision of order 0.2 dB generally requires 1000 samples. This measurement verifies that the spectral density function and the internal reference voltage of the FFT are accurately calibrated and working properly. Virtually all of the windows accurately determine the value of white noise if the first few channels are ignored as explained above. Figure 2 shows the measurement of a noise source, which has been independently determined to have a noise spectral density of 99.8 dBV/Hz by the three windows. (Appendix A shows the circuit diagram for this noise source which has an accuracy of better than 0.2 dB for frequencies from 20 to 20 kHz.)

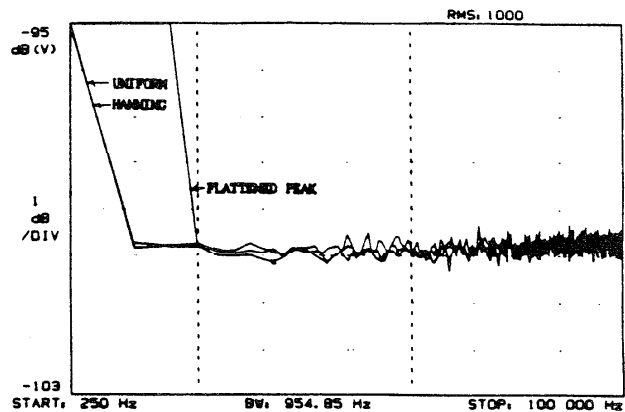
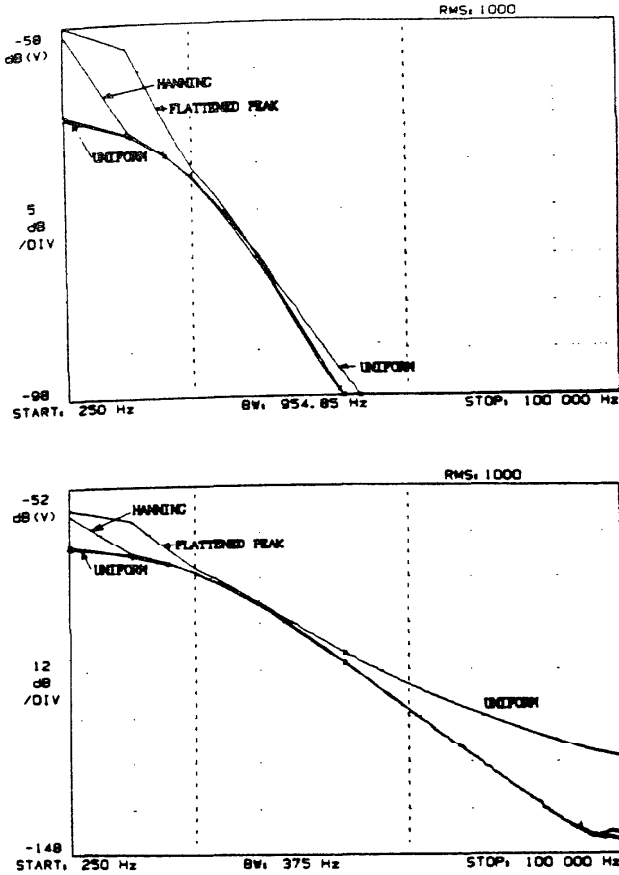


Figure 2. Spectral estimation of a white noise standard using the uniform, Hanning and the proprietary "flattened peak" windows.

Second, an approximately flat spectrum is measured over the frequency range of interest. It is not important if there are small variations in the level that change slowly over the frequency span. Third, the transfer function of the filter is determined for the frequencies of interest using a very narrow spectral source (typically an audio oscillator is sufficient). The very narrow source is accurately measured by the window since there is no problem with either high frequency or low frequency noise biasing the estimate. The use of a window with a flattened peak response is helpful but not necessary if the frequency source is sufficiently stable. This transfer function is then applied to the measured white noise spectrum in step two above. This yields a very accurate value for the "true" spectral density of the white noise source as measured through the filter. This "true value" is then compared to that obtained by the FFT analyzer. The difference between that measured in steps two and three and that measured directly with the FFT is the bias in the spectral estimate for that particular window and noise type. The accuracy of this approach comes from the fact that the calibration has been broken up into steps that can individually be determined with high precision and very small bias. The primary assumption is that the FFT analyzer is linear. Even this assumption can be checked by using precision attenuators. If the known white noise in step one does not extend to the frequencies of interest, then there is an additional assumption that the FFT is flat with frequency. This assumption is nearly always good except perhaps near the last few channels where the effect of the antialiasing filter might cause small inaccuracies.

Figures 3a and 3b show the "true" spectral estimate and the estimates as measured on a particular instrument using the uniform, Hanning, and the instrument's proprietary "flattened peak" windows for noise that varies as f^{-4} over much of the



Figures 3a (top) and 3b (bottom). Difference between the true spectrum (top) which varies approximately as f^{-4} and that estimated by the uniform, Hanning and proprietary "flattened peak" windows (bottom).

range from 1 kHz to 100 kHz. The scan is 0 to 100 kHz, and 1000 samples were taken for all curves. Note the considerable difference between the spectral estimates for channels 1 to 3 for the Hanning and proprietary "flattened peak" windows. These results confirm the theoretical calculations above showing that, for the Hanning window, the first 3 channels should be ignored. For the "flattened peak" window, the first 14 channels should be ignored. For both f^{-3} and f^{-4} noise, the uniform window does not yield usable spectral estimates over any portion of the scan. Note in this example that at frequencies above 80 kHz there is a small step in the spectral estimates. This is due to digitizing errors of the signal due to quantization. If the digitizer had more bits, these errors would not occur. This problem of dynamic range is common whenever the spectrum of interest covers many decades. The usual solution is to use filters to divide the spectrum into various frequency range segments which are suitable for the dynamic range of the FFT.

Table 1 summarizes the measured experimental biases in the spectral estimates of a particular instrument with three different windows for power-law noise types varying from f^0 to f^{-4} . This covers most of the random types of noise found in oscillators and signal processing equipment. We do not advocate using the biases reported in this table to correct data — they

Table 1. Approximate Biases in FFT Spectral Estimates

noise type f^0			
channel #	uniform	Hanning	flattened peak
1	19.6 dB	19.6 dB	20.1 dB
2	small	small	16.7 dB
3	↓	↓	7.2 dB
4			small
5			↓
noise type f^{-4}			
channel #	uniform	Hanning	flattened peak
1	unusable	8.6 dB	10.0 dB
2		0.4 dB	9.1 dB
3		0.4 dB	4.0 dB
4		small	1.2 dB
5		↓	1.1 dB
6			1.1 dB
7			1.0 dB
8			0.8 dB
9			0.6 dB
10			0.6 dB
11			0.5 dB
12			0.4 dB
13			0.4 dB
14			small
15			↓

only indicate which channels should not be relied upon for data analysis.

IV. Variances of Spectral Estimates

IV.A. Theoretical Analysis

We have derived expressions for $\text{var}\{\hat{S}(f)\}$ — the variance of $\hat{S}(f)$ — for each of the three models considered in Section III.A. These expressions depend primarily on the number of blocks N_b . Again, the details behind these calculations will be reported elsewhere [3].

First, for a white noise process, the uniform window yields

$$\text{var}\{\hat{S}(f_j)\} = S^2(f_j)/N_b, \quad 1 \leq j \leq 511.$$

while the Hanning window yields

$$\text{var}\{\hat{S}(f_j)\} = \begin{cases} 0.69S^2(f_j)/N_b, & j = 1; \\ S^2(f_j)/N_b, & 2 \leq j \leq 510; \\ 1.03S^2(f_j)/N_b, & j = 511. \end{cases}$$

These results are consistent with our experimental results and with standard statistical theory.

Second, for a random-walk process, the uniform window yields

$$\text{var}\{\hat{S}(f_j)\} = 5S^2(f_j)/N_b, \quad 1 \leq j \leq 511.$$

while the Hanning window yields

$$\text{var}\{\hat{S}(f_j)\} = \begin{cases} 1.30S^2(f_j)/N_b, & j = 1; \\ 2.20S^2(f_j)/N_b, & j = 2; \\ 1.31S^2(f_j)/N_b, & j = 3; \\ 1.15S^2(f_j)/N_b, & j = 4; \\ 1.09S^2(f_j)/N_b, & j = 5; \\ 1.06S^2(f_j)/N_b, & j = 6; \\ 1.04S^2(f_j)/N_b, & j = 7; \\ S^2(f_j)/N_b, & 8 \leq j \leq 511 \text{ to within } 3\%. \end{cases}$$

Except for the few lowest frequencies, the results for the Hanning window agree with our experimental results and with standard statistical theory; however, the factor of five in the variance for the uniform window disagrees with our experiments and with standard theory (although it has been verified by Monte Carlo techniques). The cause of this discrepancy is under investigation, but we think it is due to the band-limited nature of the experimental data.

Third, for a random-run process, the variance computations are not useful since the variance is dominated by the fact that the expected value of the sample variance for each block of samples increases with time. The agreement which we found between standard statistical theory and our experimental results on the $1/N_b$ rate of decrease of variance is undoubtedly due to the band-limited nature of the experimental data. We will attempt to verify these conclusions in the future using Monte Carlo techniques.

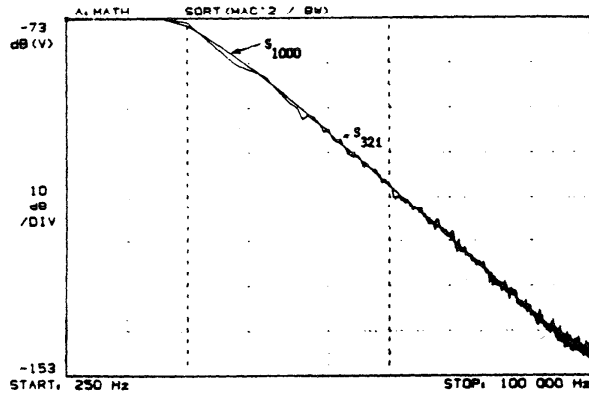


Figure 4. Comparison of the spectral estimate of f^{-4} power-law noise with 1000 samples with that obtained with 32 samples. The text explains how these two curves are used to obtain the fractional RMS confidence of the spectral estimate for 32 samples.

IV.B. Experimental Determination

The following procedure can be used to experimentally determine the variance of the spectral estimates of virtually any type of noise spectrum with any type of window for a particular instrument. Since the spectral density of interest is in general nonwhite, we must determine both the "true value" and a way to normalize the fractional error of the estimate as a function of the number of samples. This can be done by making use of the above theoretical analysis that shows that the variance should decrease as the square root of the number of samples since they are approximately statistically independent (in fact, exactly so in the cases of white and random-walk noise). As an

example we have chosen to take $N_b = 1000$ blocks of the various power-law noise types examined in III B above and compare the value of the spectral estimate with that obtained from $N_b = 32$ blocks (see Figure 4). Since the variance of the 1000 block data is about 32 times smaller than that of the 32 block data, it can serve as an accurate estimate of the "true value." Let $\hat{S}_{1000}(f_j)$ represent this quantity at the j -th channel (frequency). By subtracting the 1000 block data from the 32 block data at the j -th channel, we then have one estimate of the error for the 32 block data; by repeating this procedure over N_c different channels and N_r different replications, we can obtain accurate estimates of the variance for the 32 block data. Let $\hat{S}_{32i}(f_j)$ represent the spectral estimate for the 32 block data at the j -th channel and the i -th replication. To compensate for the variation in the level of the spectral estimates with channel, it is necessary to divide the error at the j -th channel by the "true value" $\hat{S}_{1000}(f_j)$. The mean square fractional error of the 32 block data for the noise type under study is given by

$$\sigma_{32}^2 = \frac{1}{N_r N_c} \sum_{i=1}^{N_r} \sum_j \left(\frac{\hat{S}_{32i}(f_j) - \hat{S}_{1000}(f_j)}{\hat{S}_{1000}(f_j)} \right)^2 \approx \frac{\text{var}\{\hat{S}_{32i}(f_j)\}}{S^2(f_j)}$$

It is assumed that all channels with bias — as indicated in Table 1 — have been excluded in the summation over j . It is also important that the changes in the spectral density not exceed the dynamic range of the digitizer because under this condition the quantization errors — in addition to causing biases in the spectral estimates as discussed earlier — can lead to situations where the variance does not improve as N_b increases. These values can be scaled to any number of blocks N_b if care is taken to avoid these quantization errors. Upper and lower approximate 67% confidence limits for $S(f_j)$ — the true spectral density at channel j — using Hanning, uniform and the proprietary "flattened peak" windows for N_b approximately independent blocks are given by

$$\hat{S}(f_j)(1 \pm V(\alpha, N_b))$$

where $\hat{S}(f_j)$ is the spectral estimate given by Equation (1) and $V(\alpha, N_b)$ is the fractional variance given in Table 2 for f^α and $\alpha = 0, -2, -3$ and -4 (these results were obtained by averaging over $N_r N_c = 1200$ channels). The variances obtained are very close to those obtained from standard statistical analysis for white noise, i.e.,

$$\hat{S}(f_j) \left(1 \pm \frac{1}{\sqrt{N_b}} \right)$$

Table 2. Confidence Intervals for FFT Spectral Estimates

power law noise type	uniform	Hanning	flattened peak
f^0	$1.02/\sqrt{N_b}$	$0.98/\sqrt{N_b}$	$0.98/\sqrt{N_b}$
f^{-2}	$1.02/\sqrt{N_b}$	$1.04/\sqrt{N_b}$	$1.04/\sqrt{N_b}$
f^{-3}	unusable	$1.04/\sqrt{N_b}$	$1.04/\sqrt{N_b}$
f^{-4}	unusable	$1.04/\sqrt{N_b}$	$1.04/\sqrt{N_b}$

V. Conclusions

We have introduced experimental techniques to evaluate the statistical properties of FFT spectral estimates for common noise types found in oscillators, amplifiers, mixers and similar

devices, and we have compared these with theoretical calculations. We have used these techniques to study the biases and variances of FFT spectral estimates using the uniform, Hanning, and a proprietary "flattened peak" window. The theoretical analysis was greatly hampered because the instrument manufacturer does not disclose the exact form of the "flattened peak" window or the normalization procedure for the other windows. Nevertheless, we obtained fair agreement between the theoretical and the experimental analysis. The variances of the spectral estimation were virtually identical to a few percent for f^0 to f^{-4} noise except for the uniform window which is incapable of measuring noise which falls off faster than f^{-2} . There was a very large difference in the biases of the first few channels for the three windows. The Hanning window showed significant biases in the first 3 channels while the proprietary "flattened peak" window showed large biases for f^{-4} noise even up to channel 13. The Hanning window therefore yields useful information over three times wider frequency range than the proprietary "flattened peak" window. In the particular instrument studied, the proprietary "flattened peak" window is the best choice for estimating the height of a narrow band source, while the Hanning window is by far the best choice for spectral analysis of common noise types found in oscillators, amplifiers, mixers, etc. We have also shown that the 67% confidence levels for spectral estimation as a function of the number of contiguous nonoverlapping blocks, N_b , is approximately given by

$$S = S_m \left(1 \pm \frac{0.95}{\sqrt{N_b}} \right)$$

for white noise (f^0) and by

$$S = S_m \left(1 \pm \frac{1.04}{\sqrt{N_b}} \right)$$

for noise types f^{-2} to f^{-4} . This agrees to within 4% of that found by standard statistical analysis for white noise. Using this data one can now determine the number of samples necessary to estimate — to a given level of statistical uncertainty — the spectrum of the various noise types commonly found in oscillators, amplifiers, mixers, etc.

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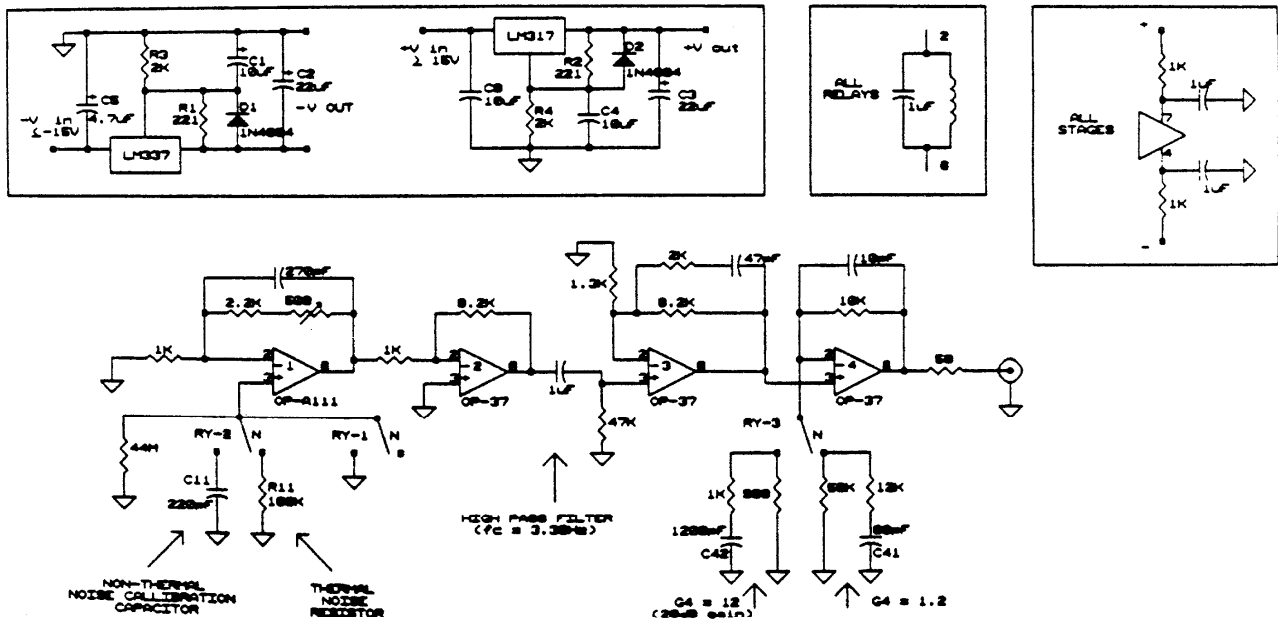


Figure 5. Circuit diagram of a precision noise source.

Appendix. Precision Noise Source

Figure 5 shows the circuit diagram of a precision noise source whose spectral density can be determined from first principles to ± 0.2 dB over the frequency range from 20 Hz to 20 kHz. The spectral density is basically given by the Johnson noise of the 10^5 ohm resistor, $V_n^2 = 4kTR$, where T is in Kelvin, and k is Boltzmann's constant. Corrections due to the input noise voltage and noise current of the amplifier amount to about 0.2 dB

for the circuit elements shown. All resistors are precision 1% metal film resistors. The output level can be switched from -100 dBV/Hz to -80 dBV/Hz. By adjusting the noise-gain capacitors one can make the noise spectrum flat to within 0.3 dB out to 200 kHz. There is also provision to measure the input noise voltage of the amplifier by shorting the input to ground or the combined noise voltage and noise current by switching a 220 pF capacitor into the input instead of the 10^5 ohm noise resistor.

A MODIFIED "ALLAN VARIANCE" WITH INCREASED OSCILLATOR CHARACTERIZATION ABILITY

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Summary

Heretofore, the "Allan Variance," $\sigma_y^2(\tau)$, has become the de facto standard for measuring oscillator instability in the time-domain. Often oscillator frequency instabilities are reasonably modelable with a power law spectrum: $S_y(f) \sim f^{-\alpha}$, where y is the normalized frequency, f is the Fourier frequency, and α is a constant over some range of Fourier frequencies. It has been shown that for power law spectrum $\sigma_y^2(\tau) \sim \tau^\mu$, and that $\mu = -\alpha - 1$ for $-3 < \alpha < +1$, where τ is the nominal sample time over which each value of y is measured. The modified "Allan Variance" developed in this paper yields $\mu \cong -\alpha - 1$ for all α in the range $-3 < \alpha$, which removes the previous ambiguity: $\mu = -2$ for $+1 < \alpha$. In other words, with the modified "Allan Variance" one can easily distinguish between white phase noise ($\alpha = +2$) and flicker phase noise ($\alpha = +1$) -- commonly occurring for the short term instabilities of quartz crystal oscillators and active hydrogen masers.

Key Words. Flicker Noise; Frequency Stability; Oscillator Noise Modeling; Power Law Spectrum; Time-Domain Stability; White Noise.

Introduction

The random fluctuations in precision oscillators may often be characterized by a power law spectrum:

$$S_y(f) = h_\alpha f^{-\alpha}, \quad (1)$$

where y is the normalized frequency deviation, f is the Fourier frequency, h_α is the intensity of the particular noise process, and α is constant over some range of f . The typical values of α are: +2 (white noise phase modulation, PM); +1 (flicker noise PM); 0 (white noise frequency modulation, FM); -1 (flicker noise FM); and -2 (random walk FM). The Allan variance, as it has come to be known,¹ has been demonstrated as a very useful statistical tool for characterizing these various random processes with the exception that if $\alpha = +1$ or +2, the dependence on τ is nominally

the same, i.e., $\sim \tau^{-2}$. It is not at all uncommon for white PM and flicker PM to occur in precision oscillators for τ of the order of one second and shorter. The modified Allan variance, as developed in this paper, depends as τ^{-2} for $\alpha = +1$ and as τ^{-3} for $\alpha = +2$. This yields a clear distinction in the time domain between these heretofore somewhat ambiguous processes.

Definition of "Allan Variance" and Related Concepts

Define y , the normalized frequency deviation, as

$$y(t) = \frac{v(t) - v_0}{v_0} \quad (2)$$

where $v(t)$ is the output frequency of the oscillator being studied, and v_0 is nominally the same frequency, but of a reference oscillator assumed for the moment without loss of generality to be better than the test oscillator. The time deviation from some arbitrary origin ($t = 0$) is the integral of the frequency deviations (from that origin):

$$x(t) = \int_0^t y(t') \cdot dt' \quad (3)$$

The i^{th} average frequency deviation over an interval, τ , is

$$\bar{y}_i = \frac{x_{i+1} - x_i}{\tau} \quad (4)$$

where the assumption is made that the time deviation measurements are nominally spaced τ apart.

The "Allan Variance" is defined as:

$$\sigma_y^2(\tau) = \frac{1}{2} \langle (\bar{y}_{i+1} - \bar{y}_i)^2 \rangle, \quad (5)$$

where the brackets " $\langle \rangle$ " denote infinite time average. Using equation (4), one may write:

$$\sigma_y^2(\tau) = \frac{1}{2\tau^2} \langle (x_{i+2} - 2x_{i+1} + x_i)^2 \rangle. \quad (6)$$

It has been shown that typically $\sigma_y^2(\tau)$ varies as τ^μ , and that $\mu = -\alpha - 1$ for $-3 \leq \alpha \leq +1$.^{1,2} Hence, we see one of the dimensions of usefulness of $\sigma_y^2(\tau)$; i.e., ascertaining the dependence on τ allows an estimate of α (the power law spectral type of noise). However, if $\alpha \geq +1$, then $\mu \cong -2$, and the τ dependence becomes somewhat ambiguous as to the type of noise in this region. It is interesting to note that in the region $\alpha \geq +1$, $\sigma_y^2(\tau)$ is bandwidth (f_h) dependent: i.e., the bandwidth of the measurement system will affect the value of $\sigma_y(\tau)$, and furthermore, one may use the bandwidth dependence³ to determine the value of α (see also Appendix Ref. 2).

Development of the Modified Allan Variance

One may also write $\sigma_y^2(\tau)$ in terms of a generalized autocovariance function:

$$\sigma_y^2(\tau) = \frac{1}{2\tau^2} [4U_x(\tau) - U_x(2\tau)], \quad (7)$$

where

$$U_x(\tau) = 2[R_x(0) - R_x(\tau)], \quad (8)$$

and where

$$R_x(\tau) = \langle x(t+\tau) \cdot x(t) \rangle, \quad (9)$$

the classical autocovariance function of $x(t)$. Using the Fourier transforms of generalized functions, one may determine the coefficients relating the power spectral density to $\sigma_y^2(\tau)$. Ref. 1 gives these relationships. It is of interest to note that $U_x(\tau)$ has the following approximate form in the region $\alpha \geq +1$ (see Appendix Ref. 2):

$$U_x(\tau) \sim a(\alpha) \left[\left| \frac{1}{2\pi f_h} \right|^{-\alpha+1} - |\tau|^{-\alpha+1} \right] \quad (10)$$

Hence, one notes that by changing the reciprocal bandwidth as well as τ , one affects $\sigma_y^2(\tau)$ in similar ways, depending on the value of α . From this, one should be able to deduce the value of α , since the bandwidth dependence becomes stronger for α moving positive from +1, and the τ dependence becomes stronger as α moves negative from +1. One can change the bandwidth in the hardware or in the software. In the past, it has typically been done in the hardware.³ James Snyder⁴ has shown that it is relatively easy to change the bandwidth in the data processing by a clever technique and we have followed his lead. In particular, we have chosen a new variance analysis scheme which coincides with the Allan variance at the minimum sample time, τ_0 , (i.e., minimum data spacing), but which changes the bandwidth in the software as the sample time, τ , is changed.

Each reading of the time deviation, x_i , has associated with it an intrinsic nominal (hardware) measurement system bandwidth, f_h . Define $\tau_h = \frac{1}{2\pi f_h}$; and similarly we may define a software bandwidth, $f_s = f_h/n$, which is $1/n$ times narrower than the hardware bandwidth. This software bandwidth can be realized by averaging n adjacent x_i 's; $\tau_s = n\tau_h$, where $\tau_s = 1/f_s$. We have defined a modified Allan variance which allows the reciprocal software bandwidth to change linearly with the sample time, τ :

$$\text{Mod } \sigma_y^2(\tau) = \frac{1}{2\tau^2} \left\langle \left[\frac{1}{n} \sum_{i=1}^n (x_{i+2n} - 2x_{i+n} + x_i) \right]^2 \right\rangle \quad (11)$$

where $\tau = n\tau_0$. Eq. 11 clearly coincides with Eq. 6 for $n = 1$. One can see that, in general, we have formed a second difference of three time readings with each of the three being an average of n of the x_i 's (with non-overlapping averages). As n increases, the (software) bandwidth decreases and this bandwidth varies just as $f_s = f_h/n$.

For a finite data set of N readings of x_i ($i = 1$ to N), we may write an estimate:

$$\text{Mod } \sigma_y^2(\tau) = \frac{1}{2\tau^2 n^2 (N-3n+1)} \cdot$$

$$\sum_{j=1}^{N-3n+1} \left[\sum_{i=j}^{n+j-1} (x_{i+2n} - 2x_{i+n} + x_i) \right]^2 \quad (12)$$

Eq. 12 is easy to program, but takes more time to compute than for $\sigma_y(\tau)$. This is only of significance for the smaller computer or handheld calculator.

Comparisons, Tests, and Examples of Usage of the Modified Allan Variance

We simulated various power law noise processes, and applied Eq. 12. Shown in Fig. 1 are the resulting τ -dependences of the modified Allan variances for $\alpha = -2, -1, 0, +1$, and $+2$. The solid lines drawn are the anticipated or theoretical slopes for the particular noise process. One sees excellent agreement for white noise PM and for flicker noise PM, and nominal agreement for the others.

One can express Eq. 11 in terms of the generalized autocovariance function:

$$\begin{aligned} \text{Mod } \sigma_y^2(\tau) = & \frac{1}{2\tau^2 n^2} \left\{ [4U_x(\tau) - U_x(2\tau)] \cdot n \right. \\ & + \sum_{i=1}^{n-1} (n-i) [-6U_x(i\tau_0) + 4U_x((n+i)\tau_0) \\ & + 4U_x((n-i)\tau_0) - U_x((2n+i)\tau_0) \\ & \left. - U_x((2n-i)\tau_0)] \right\} \quad (13) \end{aligned}$$

In the range $-3 \leq \alpha \leq +1$, one may write:

$$U_x(\tau) = a(\alpha) \cdot \tau^{-\alpha+1}, \quad (14)$$

which when substituted in Eq. 13, and using Eq. 7, yields

$$\begin{aligned} \text{Mod } \sigma_y^2(\tau) = & \sigma_y^2(\tau) \left\{ \frac{1}{n} + \frac{1}{n^2 4n^{-\alpha+1} - (2n)^{-\alpha+1}} \cdot \right. \\ & \sum_{i=1}^{n-1} (n-i) \cdot \left[-6i^{-\alpha+1} + 4(n+i)^{-\alpha+1} \right. \\ & \left. \left. - (2n+i)^{-\alpha+1} + 4(n-i)^{-\alpha+1} - (2n-i)^{-\alpha+1} \right] \right\} \quad (15) \end{aligned}$$

Since we know that $\sigma_y^2(\tau)$ is well behaved in this range and $\mu = -\alpha - 1$, it is of interest to look at the ratio:

$$R(n) = \text{Mod } \sigma_y^2(\tau) / \sigma_y^2(\tau) \cdot$$

As stated before, at $n = 1$ ($\tau = \tau_0$) the ratio is unity. One can evaluate Eq. 15 with a computer. A reasonable empirical fit may be formed, which is good to 0.5% or better of Eq. 15:

$$R(n) = \frac{q+pn^E - p}{qn^E} \quad (16)$$

which approaches p/q asymptotically as n approaches infinity, and is within 1% of p/q for $n \geq 8$. Listed in Table 1 are the empirical values of p , q , and E and the quality of fit for the appropriate power law noise processes.

TABLE 1

Noise Type	α	p	q	E	fit	$\text{Mod } \sigma_y^2(\tau) / \sigma_y^2(\tau)$ $\tau \gg \tau_0$
White FM	0	1	2	2	perfect	.707
Flicker FM	-1	99.9	148	2.35	1%	.821
Random Walk FM	-2	33	40	2.35	<1%	.908
Flicker Walk FM	-3	1	1	--	perfect	1

The results of Table 1 are in reasonable agreement with simulated results of Fig. 1(a) through 1(e). The last row in Table 1, "flicker walk" frequency modulation, is out of the range of applicability of α , but the ratio, $R(n)$, is still convergent.

The $U_x(\tau)$ function for flicker noise PM is extremely complicated and has not been developed, but one can arrive at an empirical value for it. The $U_x(\tau)$ function is derivable for the other power law spectral processes. Table 2 gives the relationships between the time domain measure $\text{Mod } \sigma_y^2(\tau)$ and its power law spectral counterpart, given in Eq. 1. Also listed in the right hand column of Table 2 are the asymptotic values of $R(n)$:

TABLE 2

Noise Type	α	$\text{Mod } \sigma_y^2(\tau)$	Comment	$R(n)$
White PM	+2	$h_2 \cdot \frac{3 f_h}{(2\pi)^2 n \tau^2}$	Exact	1
Flicker PM	+1	$h_1 \cdot \frac{1.038 + 32n(\omega_h \tau)}{(2\pi)^2 \tau^2}$	Empirical	1
White FM	0	$h_0 \cdot \frac{R(n)}{2\tau}$	Exact	0.5
Flicker FM	-1	$h_{-1} \cdot 22n(2) \cdot R(n)$	Empirical; Exact Available	0.674
Random Walk FM	-2	$h_{-2} \cdot \frac{(2\pi)^2}{6} \cdot R(n)$	Empirical; Exact Available	0.824

*

It is clear from Table 2 that $\text{Mod } \sigma_y^2(\tau)$ is very useful for white PM and flicker noise PM, but for $\alpha < +1$ the conventional Allan variance, $\sigma_y^2(\tau)$, gives both an easier-to-interpret and an easier-to-calculate measure of stability.

It is interesting to make a graph of α versus μ for both the ordinary Allan variance and the modified Allan variance. Shown in Fig. 2 is such a graph. This graph allows one to determine power law spectra for non-interger as well as interger values of α . The dashed line for the modified Allan variance has been intentionally moved to the left in Fig. 2 because for small values of n the value of μ will appear to be slightly more negative than for $\sigma_y^2(\tau)$, even though for large n , they both approach the same slope (i.e., the same values of μ). In fact, in the asymptotic limit, the equation relating μ and α for the modified Allan variance is

$$\alpha = -\mu - 1, \text{ for } -3 < \alpha < +3. \quad (17)$$

* See Appendix Note # 34

The value of $\mu = -4$ for $\alpha = +3$ was verified empirically with simulated data, and it appears that for $\alpha > +3$, μ remains at -4 .

A direct application for using the modified Allan variance recently arose in the analysis of atomic clock data as received from a Global Positioning System (GPS) satellite. We were interested in knowing the short-term characteristics of the newly developed, high-accuracy NBS/GPS receiver, as well as the propagation fluctuations. Fig. 3 shows both $\sigma_y^2(\tau)$ and $\text{Mod } \sigma_y^2(\tau)$ for comparison. Using $\text{Mod } \sigma_y^2(\tau)$, we can tell that the fundamental limiting noise process involved in the system is white noise PM with the exciting result that averaging for four minutes can allow one to ascertain time difference to better than one nanosecond excluding other systematic effects.

Conclusion

We have developed a supplemental measure, the "Modified Allan Variance" ($\text{Mod } \sigma_y^2(\tau)$), which has very useful properties when analyzing oscillator or signal stability in the presence of white noise phase modulation or flicker noise phase modulation. It also works reasonably well as a stability measure for other commonly occurring noise processes in precision oscillators.

We would recommend that for most time domain analysis, $\sigma_y^2(\tau)$ should be the first choice. If $\sigma_y^2(\tau)$ depends on τ as τ^{-1} , then the modified Allan variance can be used as a substitute to help remove the ambiguity as to the noise processes.

Acknowledgments

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SIMULATED NOISE

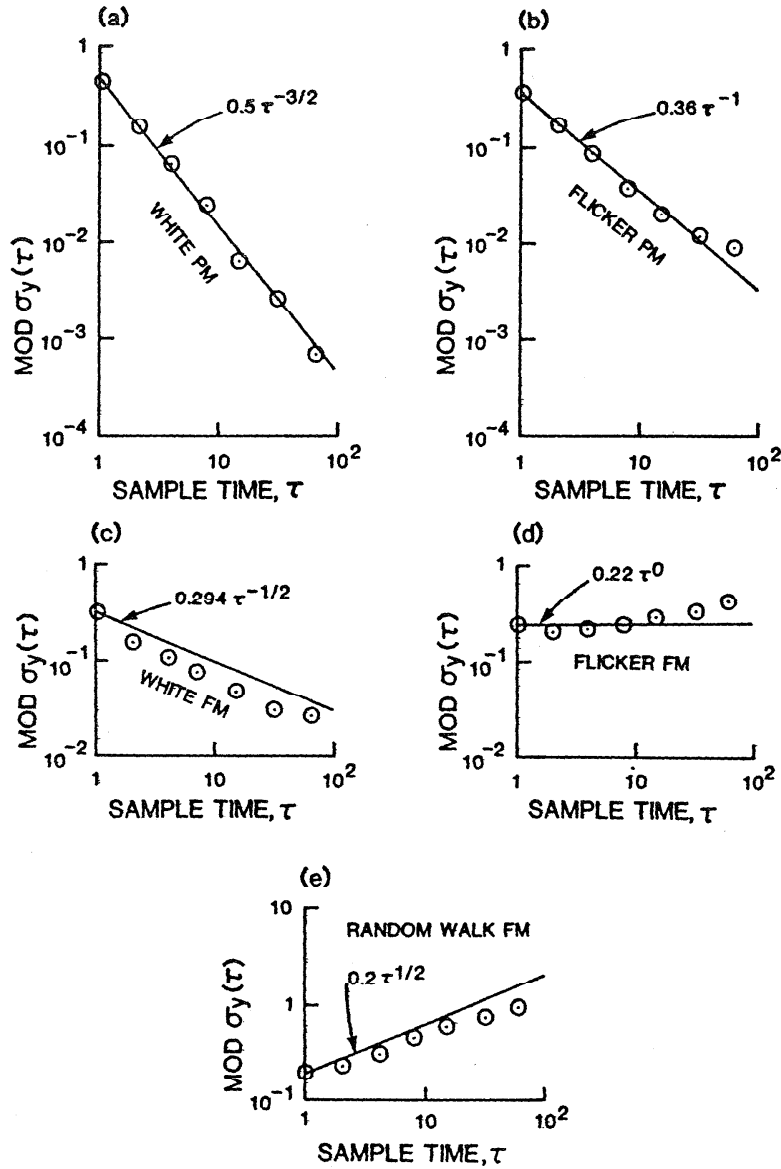


Fig 1a-e. $\text{Mod } \sigma_y(\tau)$ using Eq. 12 was calculated for different sample times for independently generated and simulated noise processes, which were white phase noise, flicker phase noise, white frequency noise, flicker frequency noise, and random walk frequency noise, respectively. $\text{Mod } \sigma_y(\tau)$ was computed for 399 data points in each case. One sees the excellent fit to the theory for white phase noise and flicker phase noise, an important new contribution in the ability to characterize oscillators having these noise processes.

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**Characterization of Frequency Stability: Analysis of
the Modified Allan Variance and Properties of
Its Estimate**

PAUL LESAGE AND THÉOPHANE AYI

Abstract—An analytical expression for the modified Allan variance is given for each component of the model usually considered to describe the frequency or phase fluctuations in frequency standards. The relation between the Allan variance and the modified Allan variance is specified and compared with that of a previously published analysis. The uncertainty on the estimate of the modified Allan variance calculated from a finite set of measurement data is discussed.

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I. INTRODUCTION

Many works [1]-[5] have been devoted to the characterization of the frequency stability of ultrastable frequency sources and have shown that the frequency noise of a generator can be easily characterized by means of the "two-sample variance" [2] of frequency fluctuations, which is also known as the "Allan variance" [2] in the special case where the dead time between samples is zero.

An algorithm for frequency measurements has been developed by J. J. Snyder [6], [7]. It increases the resolution of frequency meters, in the presence of white phase noise. It has been considered in detail by D. W. Allan and J. A. Barnes [8]. They have defined a function called the "modified Allan variance" and they have analyzed its properties for the commonly encountered components of phase or frequency fluctuations [3]. For that purpose, the authors of [8] have expressed the modified Allan variance in terms of the autocorrelation of the phase fluctuations. For each noise component, they have computed the modified Allan variance and deduced an empirical expression for the ratio between the modified Allan variance and the Allan variance.

In this paper, we show that the analytical expression of this ratio can be obtained directly, even for the noise components for which the autocorrelation of phase functions is not defined from the mathematical point of view. We give the theoretical expressions and compare them with those published in [8].

The precision of the estimate of the modified Allan variance is discussed and results related to white phase and white frequency noises are presented.

II. BACKGROUND AND DEFINITIONS

In the time domain, the characterization of frequency stability is currently achieved by means of the two-sample variance [2] $\langle \sigma_y^2(2, T, \tau) \rangle$ of fractional frequency fluctuations. It is defined as

$$\langle \sigma_y^2(2, T, \tau) \rangle = \frac{1}{2} \langle (\bar{y}_{k+1} - \bar{y}_k)^2 \rangle \quad (1)$$

where the quantity \bar{y}_k is the average value of the fractional frequency fluctuations $y(t)$ over the time interval $(t_k, t_k + \tau)$ such that

$$\bar{y}_k = \frac{1}{\tau} \int_{t_k}^{t_k + \tau} y(t) dt. \quad (2)$$

In (2), t_k represents the moment at which the k th observation time interval starts. We have

$$t_k = t_0 + kT, \quad T \geq \tau \quad (3)$$

where t_0 is an arbitrary time origin, k is a positive integer, and T is the time interval between the beginning of two successive observations.

In all the following, we assume that the dead time between samples is zero. We then have

$$T = \tau. \quad (4)$$

In this special case, the two-sample variance is well known as the Allan variance $\sigma_y^2(\tau)$

$$\sigma_y^2(\tau) = \langle \sigma_y^2(2, \tau, \tau) \rangle. \quad (5)$$

The relation between the Allan variance and $y(t)$ can be expressed as

$$\sigma_y^2(\tau) = \frac{1}{2\tau^2} \left\langle \left(\int_{t_k}^{t_k + 2\tau} y(t) dt - \int_{t_k}^{t_k + \tau} y(t) dt \right)^2 \right\rangle. \quad (6)$$

Equation (6) shows that $\sigma_y^2(\tau)$ is proportional to the true variance of the output of a linear filter with input signal $y(t)$ and impulse response $h_1(t)$ in Fig. 1.

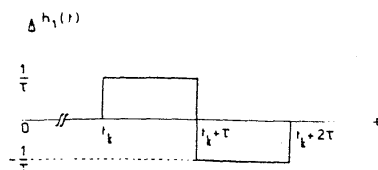


Fig. 1. Variations with time of the linear filter impulse response which represents the signal processing for the Allan variance calculation.

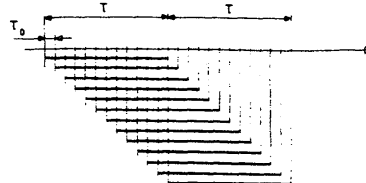


Fig. 2. Illustration of the algorithm considered for the measurement of periodic signal frequency.

The fractional frequency fluctuations $y(t)$ are actually well described by a conventional model which consists of a set of five independent noise processes [2]. Taking into account the finite bandwidth of the processed signal and assuming a single pole filter, the one-sided power spectral density $S_y(f)$ of $y(t)$ can be written as

$$S_y(f) = h_\alpha \frac{f^\alpha}{1 + \left(\frac{f}{f_c}\right)^2} \quad (7)$$

where coefficients h_α do not depend on f . The integer α equals 2, 1, 0, -1, and -2. f_c is the 3-dB bandwidth of the hardware filter.

III. THE MODIFIED ALLAN VARIANCE

The main property of the algorithm developed by J. J. Snyder is to increase the precision on the measure of periodic signal frequency, in presence of white phase noise [7]. It consists in dividing a time interval τ into n cycles of clock period τ_0 such as

$$\tau = n\tau_0. \quad (8)$$

Therefore, from a given observation time interval of duration 2τ , n overlapping time intervals of duration τ can be obtained, as depicted in Fig. 2. Another property of this algorithm is to reduce the total observation time by a factor $n/2$.

Following this way, Allan and Barnes have introduced the "modified Allan variance" [8] such as

$$\text{Mod } \sigma_y^2(\tau) = \frac{1}{2\tau^2} \left\langle \left[\frac{1}{n} \sum_{i=1}^n \left\{ \int_{t_0 + (i+n)\tau_0}^{t_0 + (i+2n)\tau_0} y(t) dt - \int_{t_0 + i\tau_0}^{t_0 + (i+n)\tau_0} y(t) dt \right\} \right]^2 \right\rangle. \quad (9)$$

It can be easily seen from (9) that the calculation of each statistical sample involved in the definition of $\text{Mod } \sigma_y^2(\tau)$ requires a signal observation of duration 3τ .

The impulse response $h_n(t)$ of the equivalent linear filter consists in finite sum of n shifted impulse responses $h_1(t)$. We have

$$h_n(t) = \frac{1}{n} \sum_{i=1}^n h_1(t - i\tau_0). \quad (10)$$

TABLE I
ANALYTICAL EXPRESSION FOR THE MODIFIED ALLAN VARIANCE WITHIN
CONDITION $2\pi f_c \tau_0 \gg 1$

NOISE TYPE	α	Mod $\sigma_y^2(\tau)$
WHITE P M	2	$\frac{3 h_2 f_c}{8 n \pi^2}$
FLICKER P M	1	$\frac{h_1}{4 \pi^2 n^2 \tau^2} \left[3n \ln(2\pi f_c \tau) + \sum_{k=1}^{n-1} (n-k) \left\{ 4 \ln \left(\frac{n^2}{k^2} - 1 \right) - 2n \left(\frac{4n^2}{k^2} - 1 \right) \right\} \right]$
WHITE F M	0	$\frac{h_0}{2} \times \frac{n^2 + 1}{2n^2}$
FLICKER F M	-1	$\frac{2h_{-1} \ln 2}{n^2} \left[\frac{4n^2 - 3n + 1}{2} + \frac{1}{n^2 \ln 2} \times \sum_{k=1}^{n-1} (n-k) \times \left\{ \frac{n}{2} \left[(k+2n) \ln(k+2n) - (k-2n) \ln(2n-k) \right] + \frac{1}{2} (k+n)(k-2n) \ln(k+n) + \frac{1}{2} (k-n)(k+2n) \ln k-n + 3k^2 \ln k - k \left[(n+2k) \ln \left(k + \frac{n}{2} \right) - (n-2k) \ln \left k - \frac{n}{2} \right \right] \right\} \right]$
RANDOM WALK F M	-2	$\frac{33}{40} + \frac{1}{8n^2} + \frac{1}{20n^4}$

*

In order to illustrate (10), variations with time of the shifted functions $h_1(t - i\tau_0)$ and of the impulse response $h_n(t)$ are represented in Fig. 3(a) and (b), respectively, for $n = 10$.

For $n = 1$, the Allan variance and the modified one are equal. We have

$$\text{Mod } \sigma_y^2(\tau) = \sigma_y^2(\tau). \tag{11}$$

One can express (9) in terms of the spectral density $S_y(f)$. We have

$$\begin{aligned} \text{Mod } \sigma_y^2(\tau) = & \frac{2}{n^2 \pi^2 \tau^2} \cdot \left\{ n \int_0^\infty \frac{1}{f^2} S_y(f) \sin^4(\pi f n \tau_0) df \right. \\ & + 2 \sum_{k=1}^{n-1} (n-k) \int_0^\infty \frac{1}{f^2} S_y(f) \cos(2\pi k f \tau_0) \\ & \left. \cdot \sin^4(\pi f n \tau_0) df \right\}. \tag{12} \end{aligned}$$

It should be noted that the integrals involved in (12) are convergent for each noise component. The analytical expression or the modified Allan variance can therefore be deduced directly from this equation.

In the following, it is assumed that the condition $2\pi f_c \tau_0 \gg 1$ is fulfilled. This means that the hardware bandwidth of the measurement system must be much larger than the reference clock frequency.

We have calculated the modified Allan variance for each noise component. Results are reported in Table I. It appears that the analytical expression for $\text{Mod } \sigma_y^2(\tau)$ is relatively simple except for flicker phase and flicker frequency noises where it is given as a finite sum of functions depending on n . In order to compare the Allan variance with the modified one, we

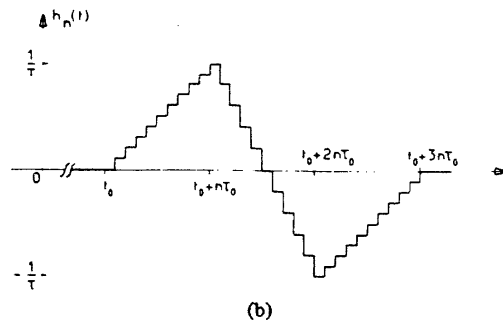
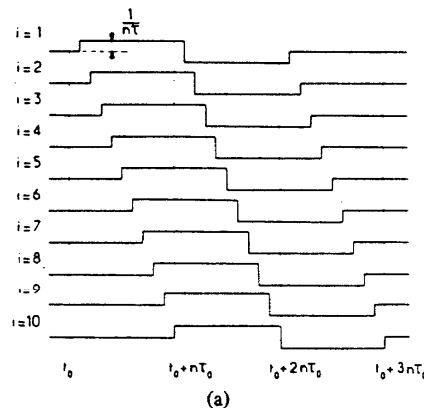


Fig. 3. (a) The impulse response $h_n(t)$, associated with the modified Allan variance calculation, represented as a sum of n shifted impulse response $h_1(t)$. It is assumed $n = 10$. (b) Variations with time of the impulse response $h_n(t)$, in the special case where $n = 10$.

* See Appendix Note # 35

TABLE II
ANALYTICAL EXPRESSIONS AND ASYMPTOTICAL VALUE FOR $R(n)$
(Results Are Valid within Condition $2\pi f_c \tau_0 \gg 1$)

α	$R(n)$	$\lim_{n \rightarrow \infty} R(n)$
2	$\frac{1}{n}$	0
1	$\frac{1}{n^2} \left[n + \frac{1}{3 \pi n (2\pi f_c \tau_0)} \times \sum_{k=1}^{n-1} (n-k) \left\{ 4 \pi n \left(\frac{n^2}{k^2} - 1 \right) - \pi n \left(\frac{4n^2}{k^2} - 1 \right) \right\} \right]$	0
0	$\frac{n^2 + 1}{2n^2}$	0.5
-1	$\frac{1}{n^2} \left[\frac{4n^2 - 3n + 1}{2} + \frac{1}{n^2 \pi n^2} \times \sum_{k=1}^{n-1} (n-k) \times \left\{ \frac{n}{2} \left[(k+2n) \pi n (k+2n) - (k-2n) \pi n (2n-k) \right] + \frac{1}{2} (k+n)(k-2n) \pi n (k+n) + \frac{1}{2} (k-n)(k+2n) \pi n k-n + 3k^2 \pi n k - k \left[(n+2k) \pi n (k + \frac{n}{2}) - (n-2k) \pi n k - \frac{n}{2} \right] \right\} \right]$	0.787
-2	$\frac{33}{40} + \frac{1}{8n^2} + \frac{1}{20n^4}$	0.825

consider the ratio $R(n)$ defined in [8] as

$$R(n) = \text{Mod } \sigma_y^2(\tau) / \sigma_y^2(\tau) \tag{13}$$

The analytical expressions for $R(n)$, deduced from Table I, are reported in Table II. One can see that $R(n)$ does not depend on the product $f_c \tau_0$, except for flicker phase noise modulation. The asymptotic values of $R(n)$ are also listed in Table II.

Fig. 4 depicts the variations of $R(n)$ with n . It shows that, for large values of n , white phase and flicker phase noise modulations have different dependences. As outlined in [8], this gives a means to easily distinguish these two noise processes, in the time domain. For large n , and for $\alpha = 0, -1, -2$, $R(n)$ remains a constant. Consequently the Allan variance can be deduced from the modified one, for these noise processes.

A comparison with results of [8] shows a good agreement for $\alpha = 2, 0$ and -2 . But, for $\alpha = 1$ and -1 , our expressions for the modified Allan variance and ratio $R(n)$ disagree, especially for flicker phase noise modulation. This discrepancy might be due to the fact that in [8], $\text{Mod } \sigma_y^2(\tau)$ is expressed in terms of the autocorrelation function of phase fluctuations which is not defined for $\alpha = 1$.

IV. UNCERTAINTY ON THE ESTIMATE OF THE MODIFIED ALLAN VARIANCE

Equation (9) shows that the definition of the modified Allan variance theoretically implies an infinite set of time intervals.

Practically, one can only estimate this quantity from a finite set of m successive cycles similar to the one depicted in Fig. 3(b).

Let $\text{Mod } \hat{\sigma}_y^2(\tau)$ be the estimated modified Allan variance (EMAV) such as

$$\text{Mod } \hat{\sigma}_y^2(\tau) = \frac{1}{2\tau^2 n^2} \times \frac{1}{m} \sum_{k=1}^m \left\{ \sum_{i=1}^n A_{i,k} \right\}^2 \tag{14}$$

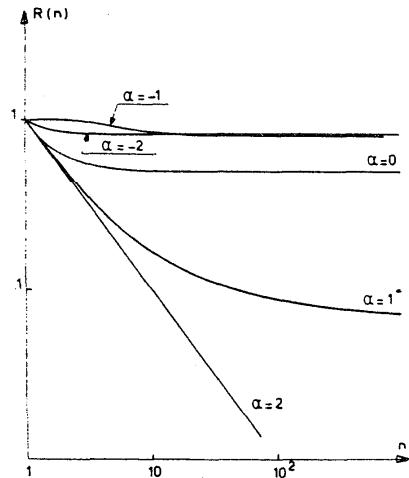


Fig. 4. Variations with n of the ratio $R(n)$, for fractional frequency fluctuations with power law spectrum $S_y(f) = h_\alpha \cdot (1/1 + (f/f_c)^2)^{\alpha}$, within condition $2\pi f_c \tau_0 \gg 1$. (*For $\alpha = +1$, $R(n)$ is a function of f_c and τ_0 . The reported variations are for $2\pi f_c \tau_0 = 10^4$.)

where

$$A_{i,k} = \int_{t_0 + (i+n)\tau_0 + (k-1)3\pi\tau_0}^{t_0 + (i+2n)\tau_0 + (k-1)3\pi\tau_0} y(t) dt - \int_{t_0 + i\tau_0 + (k-1)3\pi\tau_0}^{t_0 + (i+n)\tau_0 + (k-1)3\pi\tau_0} y(t) dt. \tag{15}$$

* See Appendix Note # 35

The EMAV is a random function of m . Its calculation requires an observation time of duration $3m\tau$.

We consider ϵ , the fractional deviation of the EMAV relative to the modified Allan variance defined as follows:

$$\epsilon = \frac{\text{Mod } \hat{\sigma}_y^2(\tau) - \text{Mod } \sigma_y^2(\tau)}{\text{Mod } \sigma_y^2(\tau)} \quad (16)$$

The standard deviation $\sigma(\epsilon)$ of ϵ defines the relative uncertainty on the measurement of the modified Allan variance, due to the finite number of averaging cycles. We have

$$\sigma(\epsilon) = \frac{1}{\text{Mod } \sigma_y^2(\tau)} \{ \sigma^2 [\text{Mod } \hat{\sigma}_y^2(\tau)] \}^{1/2} \quad (17)$$

where $\sigma^2 [\text{Mod } \hat{\sigma}_y^2(\tau)]$ denotes the true variance of the EMAV such as

$$\sigma^2 [\text{Mod } \hat{\sigma}_y^2(\tau)] = \langle [\text{Mod } \hat{\sigma}_y^2(\tau)]^2 \rangle - [\text{Mod } \sigma_y^2(\tau)]^2 \quad (18)$$

We assume that the fluctuations $y(t)$ are normally distributed [10]. One can therefore express $([\text{Mod } \hat{\sigma}_y^2(\tau)]^2)$ as

$$\begin{aligned} m^2 \langle [\text{Mod } \hat{\sigma}_y^2(\tau)]^2 \rangle &= (m^2 + 2m) [\text{Mod } \sigma_y^2(\tau)]^2 \\ &+ 4 \sum_{p=1}^{m-1} (m-p) \left\{ 2 \sum_{i=1}^{n-1} (n-i) I_n + n I_0 \right\}^2 \end{aligned} \quad (19)$$

where I_n are integrals which depend on n and on the noise process. We have

$$\begin{aligned} 8\pi^2 \tau^2 n^2 I_n &= \int_0^\infty \frac{S_y(f)}{f^2} \cos 6\pi n p f \tau_0 \\ &\times \{ 6 \cos 2\pi f \tau_0 i - 4 \cos 2\pi f \tau_0 (i+n) \\ &- 4 \cos 2\pi f \tau_0 (i-n) + \cos 2\pi f \tau_0 (i+2n) \\ &+ \cos 2\pi f \tau_0 (i-2n) \} df. \end{aligned} \quad (20)$$

For each noise component, the expression for $\sigma(\epsilon)$ can be deduced from the calculation of integrals involved in (20). These expressions are generally lengthy and complicated except for white phase and white frequency noise modulations, where integrals I_n equal zero. We have limited the present analysis to these two noise components. We get for $\sigma(\epsilon)$

$$\sigma(\epsilon) = \frac{2}{m}, \quad \text{for } \alpha = 2 \text{ and } 0. \quad (21)$$

We now compare (21) with previously published results related to the estimate of the Allan variance [5]. For a given time observation of duration $3m\tau$, it can be easily deduced from [5] that the relative uncertainty on the estimate of the Allan variance varies asymptotically as $1.14 m^{-1/2}$ and $1.0 m^{-1/2}$ for $\alpha = 2$ and 0 , respectively. For these two noise components, the uncertainty on the EMAV is larger than the uncertainty on the estimated Allan variance, but of the same order of magnitude.

V. CONCLUSION

We have calculated the analytical expression for the modified Allan variance for each component of the model usually considered to characterize random frequency fluctuations in precision oscillators. These expressions have been compared with previously published results and the link between the Allan variance and the modified Allan variance has been specified.

The uncertainty on the estimate of the modified Allan vari-

ance has been studied and numerical values have been reported for white phase and white frequency noise modulations.

In conclusion, the modified Allan variance appears to be well suited for removing the ambiguity between white and flicker phase noise modulation. Nevertheless, the calculation of the modified Allan variance requires signal processing which is complicated, compared to the Allan variance. In the presence of white or flicker phase noise, the Allan variance cannot be easily deduced from the modified Allan variance. Furthermore, for a given source exhibiting different noise components, the determination of the Allan variance from the modified one is difficult to perform. For most of time-domain measurements, the use of the Allan variance is preferred.

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THE MEASUREMENT OF LINEAR FREQUENCY DRIFT IN OSCILLATORS

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ABSTRACT

A linear drift in frequency is an important element in most stochastic models of oscillator performance. Quartz crystal oscillators often have drifts in excess of a part in ten to the tenth power per day. Even commercial cesium beam devices often show drifts of a few parts in ten to the thirteenth per year. There are many ways to estimate the drift rates from data samples (e.g., regress the phase on a quadratic; regress the frequency on a linear; compute the simple mean of the first difference of frequency; use Kalman filters with a drift term as one element in the state vector; and others). Although most of these estimators are unbiased, they vary in efficiency (i.e., confidence intervals). Further, the estimation of confidence intervals using the standard analysis of variance (typically associated with the specific estimation technique) can give amazingly optimistic results. The source of these problems is not an error in, say, the regressions techniques, but rather the problems arise from correlations within the residuals. That is, the oscillator model is often not consistent with constraints on the analysis technique or, in other words, some specific analysis techniques are often inappropriate for the task at hand.

The appropriateness of a specific analysis technique is critically dependent on the oscillator model and can often be checked with a simple "whiteness" test on the residuals. Following a brief review of linear regression techniques, the paper provides guidelines for appropriate drift estimation for various oscillator models, including estimation of realistic confidence intervals for the drift.

I. INTRODUCTION

Almost all oscillators display a superposition of random and deterministic variations in frequency and phase. The most typical model used is^[1]:

$$X(t) = a + b \cdot t + D_r \cdot t^2/2 + \phi(t) \quad (1) \quad *$$

where $X(t)$ is the time (phase) error of the oscillator (or clock) relative to some standard; a , b , and D_r are constants for the particular clock; and $\phi(t)$ is the random part. $X(t)$ is a random variable by virtue of its dependence on $\phi(t)$.

* See Appendix Note # 36

Even though one cannot predict future values of $X(t)$ exactly, there are often significant autocorrelations within the random parts of the model. These correlations allow forecasts which can significantly reduce clock errors. Errors in each element of the model (Eq. 1) contribute their own uncertainties to the prediction. These time uncertainties depend on the duration of the forecast interval, τ , as shown below in Table 1:

TABLE 1. GROWTH OF TIME ERRORS

<u>MODEL ELEMENT NAME</u>	<u>CLOCK PARAMETER</u>	<u>RMS TIME ERROR</u>
Initial Time Error	a	Constant
Initial Freq Error	b	$\sim \tau$
Frequency Drift	Dr	$\sim \tau^2$
Random Variations	$\phi(t)$	$\sim \tau^{3/2*}$

*The growth of time uncertainties due to the random component, $\phi(t)$, can have various time dependencies. The three-halves power-law shown here is a "worst case" model.[2]

One of the most significant points provided by Table 1 is that eventually, the linear drift term in the model over-powers all other uncertainties for sufficiently long forecast intervals! While one can certainly measure (i.e., estimate) the drift coefficient, Dr , and make corrections, there must always remain some uncertainty in the value used. That is, the effect of a drift correction based on a measurement of Dr , is to reduce (hopefully!) the magnitude of the drift error, but not remove it. Thus, even with drift corrections, the drift term eventually dominates all time uncertainties in the model.

As with any random process, one wants not only the point estimate of a parameter, but one also wants the confidence interval. For example, one might be happy to know that a particular value (e.g., clock time error) can be estimated without bias, he may still want to know how large an error range he should expect. Clearly, an error in the drift estimate (see Eq. 1) leads directly to a time error and hence the drift confidence interval leads directly to a confidence interval for the forecast time.

II. LEAST SQUARES REGRESSION OF PHASE ON A QUADRATIC

A conventional least squares regression of oscillator phase data on a quadratic function reveals a great deal about the general problems. A slight modification of Eq. 1 provides a conventional model used in regression analysis[3]:

$$X(t) = a + b \cdot t + c \cdot t^2 + \phi(t) \quad (2)$$

where $c = Dr/2$. In regression analysis, it is customary to use the symbol "Y" as the dependent variable and "X" as the independent variable. This is in conflict with usage in time and frequency where "X" and "Y" (time error and frequency error, respectively) are dependent on a coarse measure of time, t , the independent variable. This paper will follow the time and frequency custom even though this may cause some confusion in the use of regression analysis text books.

The model given by Eq. 2 is complete if the random component, $\phi(t)$, is a white noise (i.e., random, uncorrelated, normally distributed, zero mean, and finite variance).

III. EXAMPLE

One must emphasize here that ALL results regarding parameter error magnitudes and their distributions are totally dependent on the adequacy of the model. A primary source of errors is often autocorrelation of the residuals (contrary to the explicit model assumptions). While simple visual inspection of the residuals is often sufficient to recognize the autocorrelation problem, "whiteness tests" can be more objective and precise.

This section analyzes a set of 94 hourly values of the time difference between two oscillators. Figure 1 is a plot of the time difference (measured in microseconds) between the two oscillators. The general curve of the data along with the general expectation of frequency drift in crystal oscillators leads one to try the quadratic behavior (Model #1; models #2 and #3 discussed below). While it is not common to find white phase noise on oscillators at levels indicated on the plot, that assumption will be made temporarily. The results of the regression are summarized in a conventional Analysis of Variance, Table 2.

TABLE 2. ANALYSIS OF VARIANCE QUADRATIC FIT TO PHASE
(Units: seconds squared)

<u>SOURCE</u>	<u>SUM OF SQUARES</u>	<u>d.f.</u>	<u>MEAN SQUARE</u>
Regression	2.32E-9	3	
Residuals	4.26E-12	91	4.68E-14
Total	2.329E-9	94	2.48E-11

Coefficient of simple determination 0.99713

Parameters:

$$\begin{aligned} \hat{a} &= 1.10795E-5 \text{ (seconds)} & t\text{-ratio} &= 161.98 \\ \hat{b} &= 1.4034E-10 \text{ (sec/sec)} & t\text{-ratio} &= 152.02 \\ \hat{c} &= -3.7534E-16 \text{ (sec/sec}^2\text{)} & t\text{-ratio} &= -143.51 \end{aligned}$$

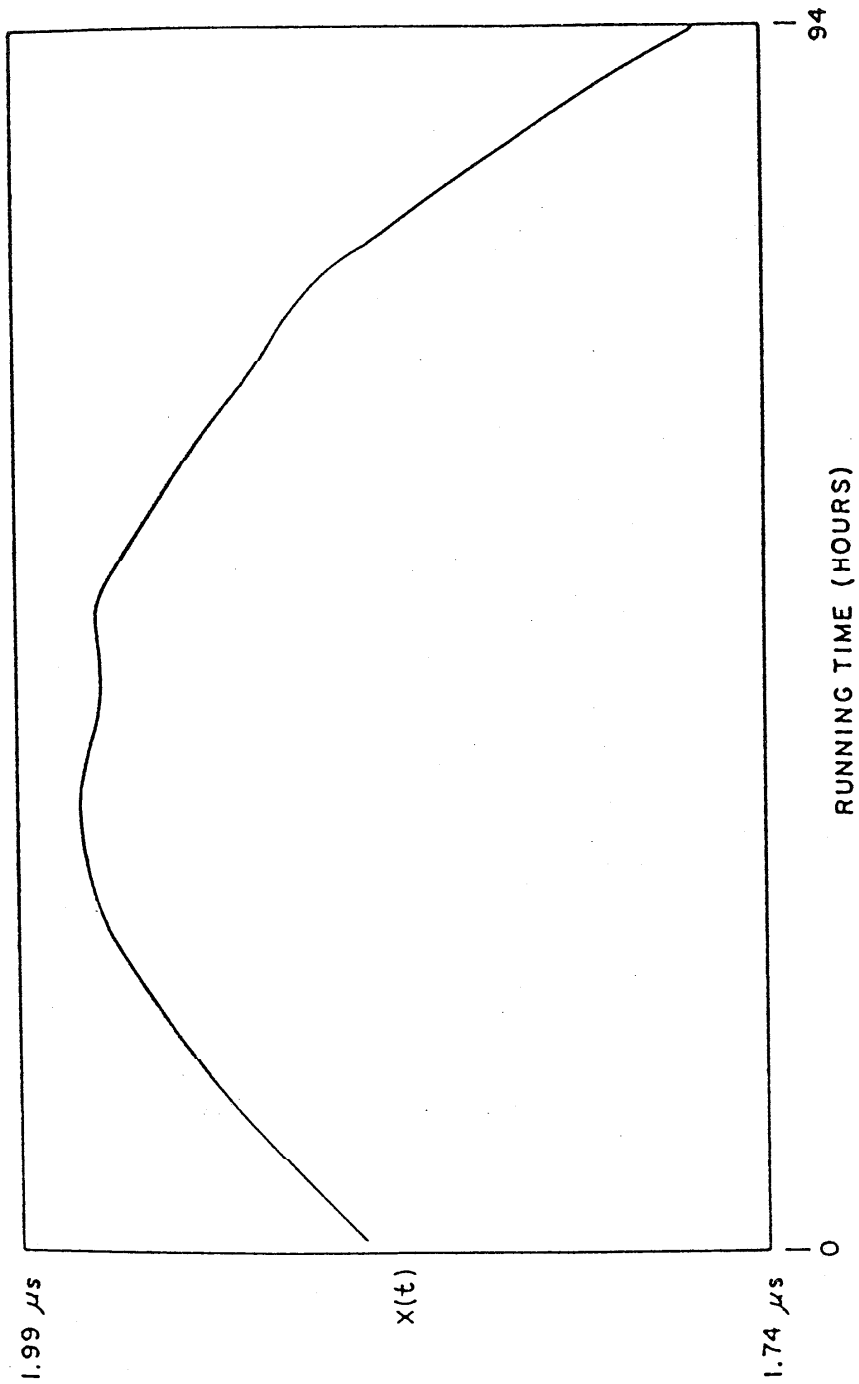


Fig. 1 - Phase difference

$$\hat{D}_r = 2\hat{c} = -7.507E-16 \text{ (about } -6.5E-11 \text{ per day)}$$

$$\text{Std. Error} = 0.05231E-16$$

(Note: \hat{a} is the value estimated for the a-parameter, etc.)

The Analysis of Variance, Table 2, above suggests an impressive fit of the data to a quadratic function, with 99.71% of the variations in the data "explained" by the regression. The estimated drift coefficient, \hat{D}_r , is $-7.507E-16$ (sec/sec²) or about $-6.5E-11$ per day --- 143 times the indicated standard error of the estimate. However, Figure 2, a plot of the residuals, reveals significant autocorrelations even visually and without sensitive tests. (The autocorrelations can be recognized by the essentially smooth variations in the plot. See Fig. 5 as an example of a more nearly white data set.) It is true that the regression reduced the peak-to-peak deviations from about 18 microseconds to less than one microsecond. It is also true that the drift rate is an unbiased estimate of the actual drift rate, but the model assumptions are NOT consistent with the autocorrelation visible in Fig. 2. This means that the confidence intervals for the parameters are not reliable. In fact, the analysis to follow will show just how extremely optimistic these intervals really are.

At this point we can consider at least two other simple analysis schemes which might provide more realistic estimates of the drift rate and its variance. Each of the two analysis schemes has its own implicit model; they are:

- (2) Regress the beat frequency on a straight line.
(Model: Linear frequency drift and white FM.)
- (3) Remove a simple average from the second difference of the phase.
(Model: Linear frequency drift and random walk FM.)

Continuing with scheme 2, above, the (average) frequency, $\bar{Y}(t)$, is the first difference of the phase data divided by the time interval between successive data points. The regression model is:

$$\bar{Y}(t) = b + D_r \cdot t + \epsilon(t) \tag{3} \quad *$$

where $\epsilon(t) = [\phi(t + \tau_0) - \phi(t)]/\tau_0$. Following standard regression procedures as before, the results are summarized in another Analysis of Variance Table, Table 3.

TABLE 3. ANALYSIS OF VARIANCE LINEAR FIT TO FREQUENCY
(Units: sec²/sec²)

<u>SOURCE</u>	<u>SUM OF SQUARES</u>	<u>d.f.</u>	<u>MEAN SQUARE</u>
Regression	1.879E-18	2	
Residuals	2.084E-20	91	2.29E-22
Total	1.899E-18	93	2.042E-20

(Note: Taking the first differences of the original data set reduces the number of data points from 94 to 93.)

* See Appendix Note # 37

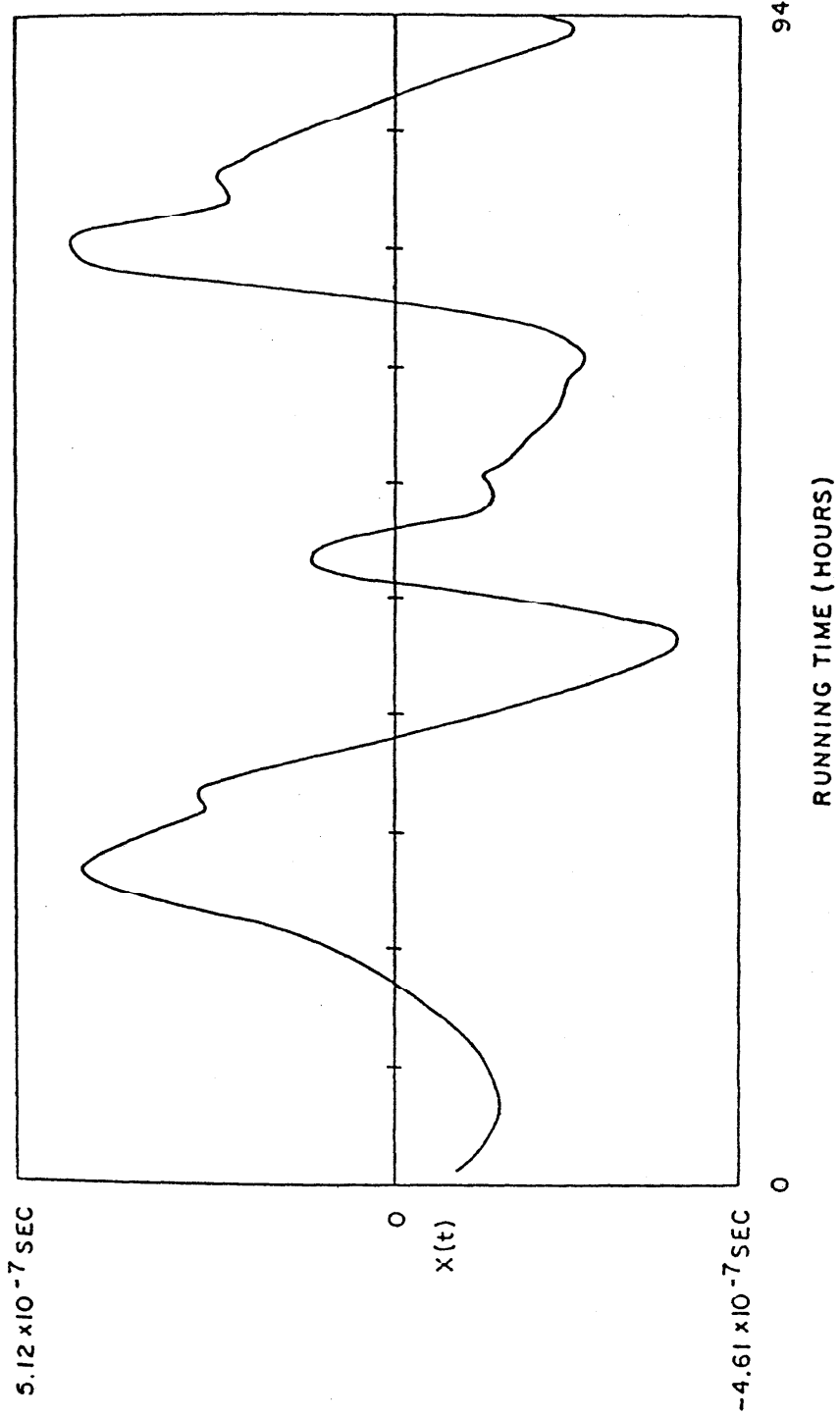


Fig. 2 - Residuals from quadratic fit to phase

Coefficient of simple determination 0.9605

Parameters:

$$\hat{b} = 1.049E-10 \text{ (sec/sec)} \quad t\text{-ratio} = 3.32$$

$$\hat{D}_r = -7.635E-16 \text{ (sec/sec}^2\text{)} \quad t\text{-ratio} = -4.70$$

Std. Error = 0.1624E-16 (sec/sec²)

While the drift rate estimates for the two regressions are comparable in value (-7.507E-16 and -7.635E-16), the standard errors of the drift estimates have gone from 0.052E-16 to 0.162E-16 (a factor of 3). The linear regression's coefficient of simple determination is 96.05% compared to 99.17% for the quadratic fit. Figure 3 shows the residuals from the linear fit and they appear more nearly white. A cumulative periodogram^[4] is a more objective test of whiteness, however. The periodogram, Fig. 4, does not find the residuals acceptable at all.

IV. DRIFT AND RANDOM WALK FM

In the absence of noise, the second difference of the phase would be a constant, $D_r \cdot \tau_0^2$. If one assumes that the second difference of the noise part is white, then one has the classic problem of estimating a constant (the drift term), in the presence of white noise (the second difference of the phase noise). Of course, the optimum estimate of the drift term is just the simple mean of the second difference divided by τ_0^2 . The results are summarized below, Table 4:

TABLE 4. SIMPLE MEAN OF SECOND DIFFERENCE PHASE

Simple mean	$\hat{D}_r = -6.709E-16$	$t\text{-ratio} = -2.45$
Degrees of Freedom	= 91	
Standard Deviation	$\hat{s} = 26.2405E-16$	
Standard Deviation of the Mean	= 2.7358E-16	

Figure 5 shows the second difference of the phase after the mean was subtracted. Visually, the data appear reasonably white, and the periodogram, Fig. 6, cannot reject the null hypothesis of whiteness. Now the standard error of the drift term is 2.735E-16, 52 times larger than that computed for the quadratic fit! Indeed, the estimated drift term is only 2.45 times its standard error.

V. SUMMARY OF TESTS

The analyses reported above were all performed on a single data set. In order to verify any conclusions, all three analyses used above were performed a total of four times on four different data sets from the same pair of oscillators. Table 5 summarizes the results:

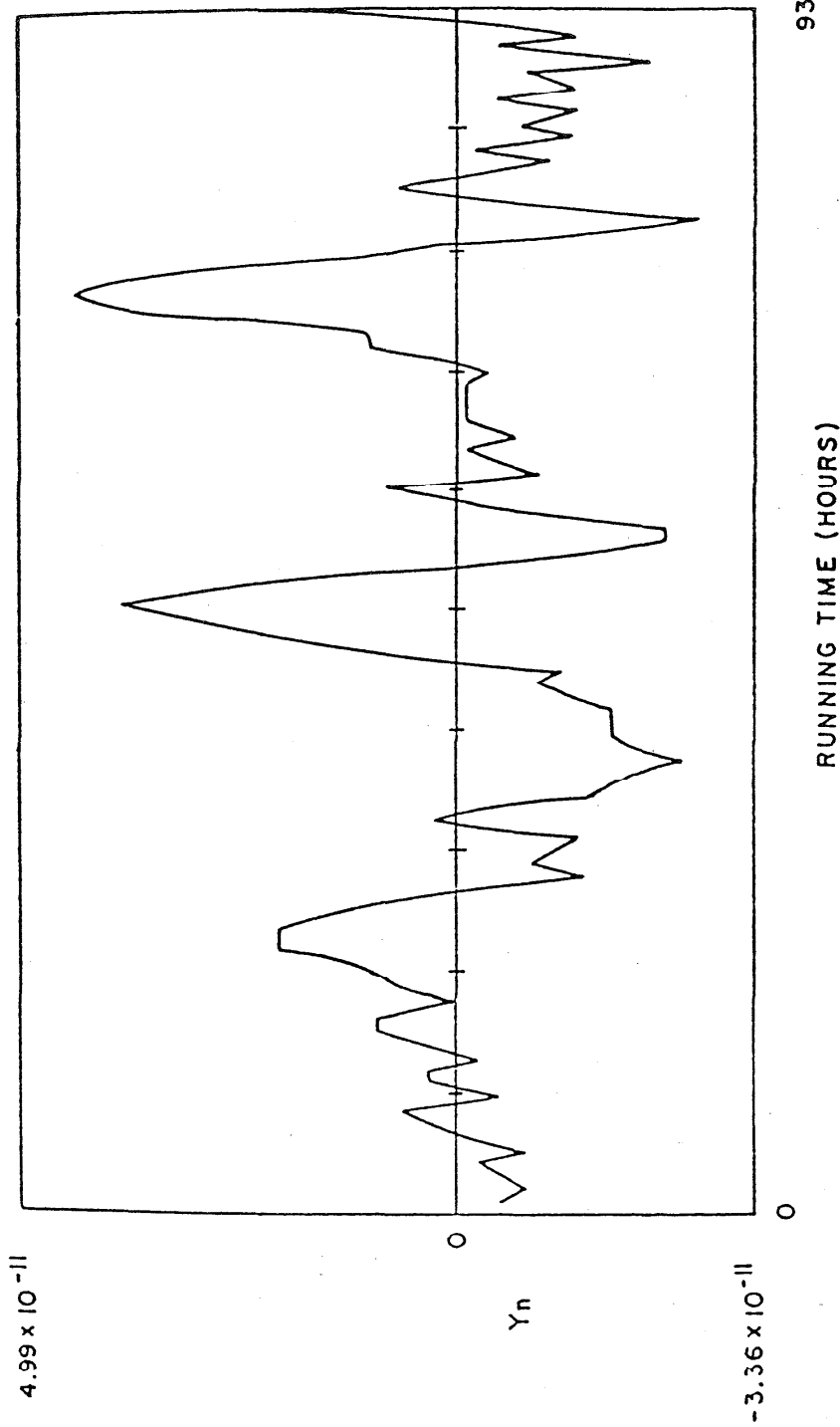


Fig. 3 - Frequency residual
(1st diff-lin)

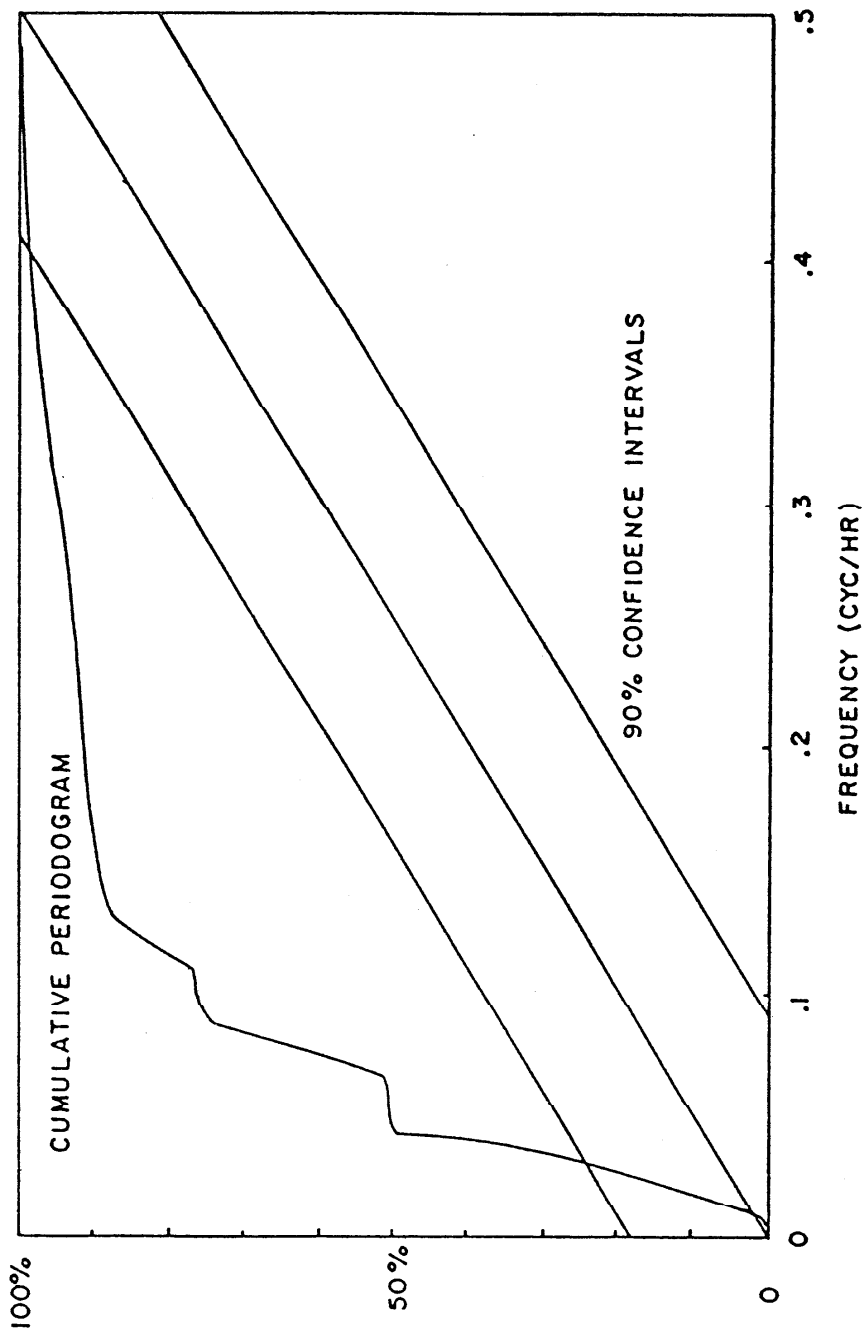


Fig. 4 - First diff - linear

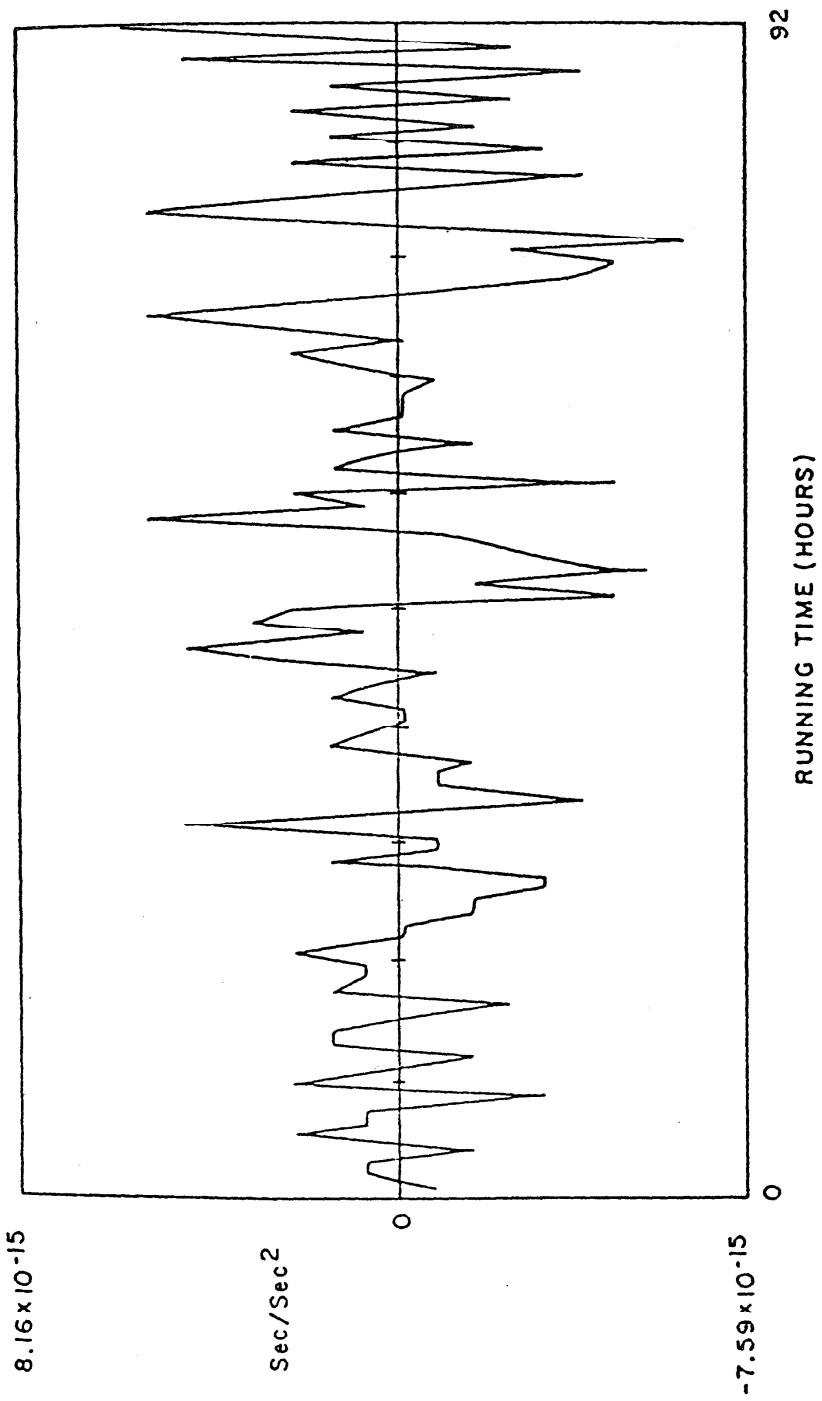


Fig. 5 - Second diff of phase

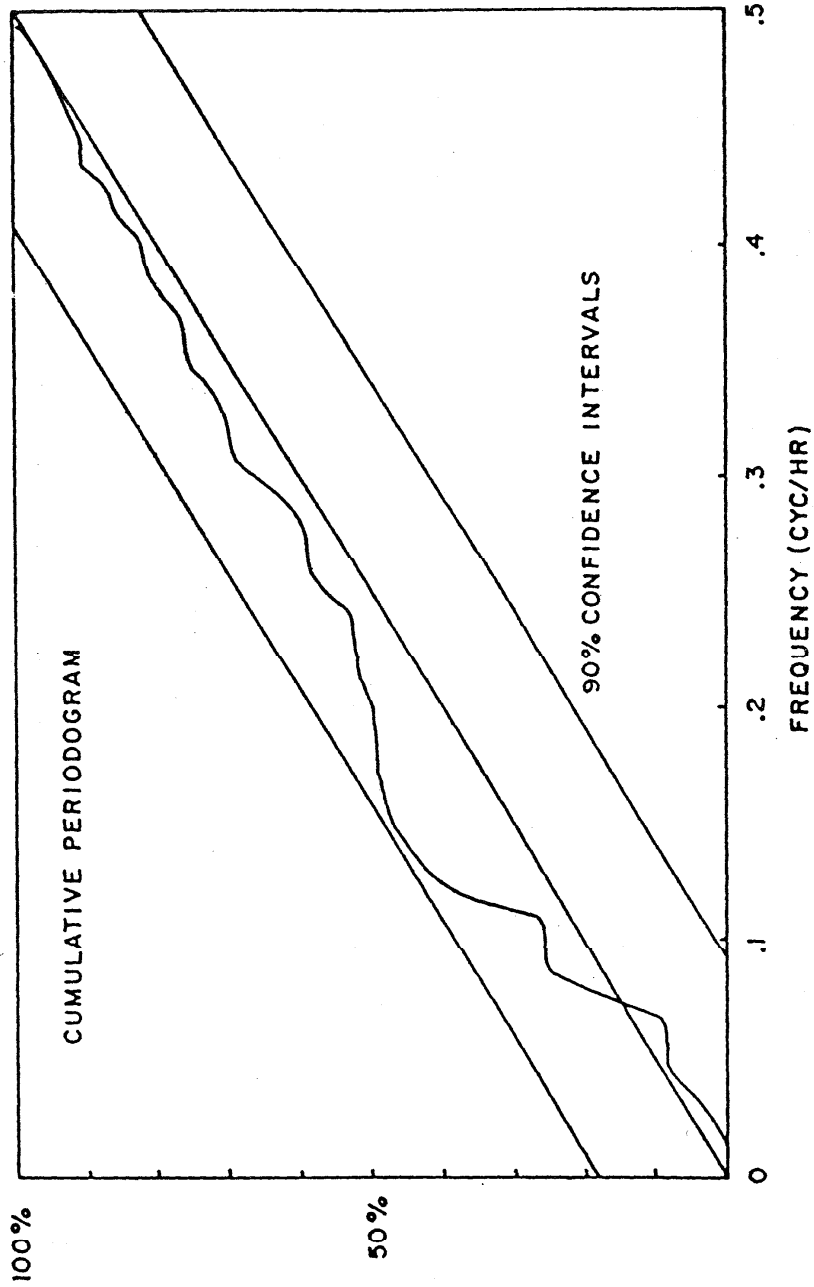


Fig. 6 - Second diff - avg.

TABLE 5. SUMMARY OF DRIFT ESTIMATES
(Units of $1.E-16$ sec/sec²)

<u>ESTIMATION PROCEDURE & MODEL</u>	<u>DRIFT ESTIMATE</u>	<u>COMPUTED STANDARD ERROR</u>	<u>PASS WHITENESS TESTS?</u>
Quad Fit (White PM and Drift)	-7.507	.0523	No
	-8.746	.0493	No
	-6.479	.0645	No
	-6.468	.0880	No
1st Difference Linear Fit (White FM and Drift)	-7.635	.162	No
	-8.558	.206	No
	-6.443	.192	No
	-6.253	.295	No
Second Difference Less Mean (Random Walk FM and Drift)	-6.710	2.736	Yes
	-7.462	9.335	No
	-6.870	3.424	Yes
	-6.412	3.543	Yes

One can calculate the sample means and variances of the drift estimates for each of the three procedures listed in Table 5, and compare these "external" estimates with those values listed in the table under "Computed Standard Error," the "internal" estimates. Of course the sample size is small and we do not expect high precision in the results, but some conclusions can be drawn. The comparisons are shown in Table 6.

It is clear that the quadratic fit to the data displays a very optimistic internal estimate for the standard deviation of the drift rate. Other conclusions are not so clear cut, but still some things can be said. Considering Table 5, the "2nd Diff - Mean" residuals passed the whiteness test three times out of four. The external estimate of the drift standard deviation lies between the internal estimates based on the first and second differences. Since the oscillators under test were crystal oscillators, one expects flicker FM to be present at some level. One also expects the flicker FM behavior to lie between white FM and random walk FM. This may be the explanation of the observed standard deviations, noted in Table 6.

TABLE 6. STANDARD DEVIATIONS

<u>PROCEDURE (Model)</u>	<u>EXTERNAL ESTIMATE (Std. Dev. of Drift Estimates from Col. #2, Table 5)</u>	<u>INTERNAL ESTIMATE (RMS Computed Std. Dev. Col. #3, Table 5)</u>
Quad Fit (White PM)	1.08	0.065
1st Diff - Lin (White FM)	1.08	0.203
2nd Diff - Mean (Rand Wlk FM)	0.44	5.45

VI. DISCUSSION

In all three of the analysis procedures used above, more parameters than just the frequency drift rate were estimated. Indeed, this is generally the case. The estimated parameters included the drift rate, the variance of the random (white) noise component, and other parameters appropriate to the specific model (e.g., the initial frequency offset for the first two models). If these other parameters could be known precisely by some other means, then methods exist to exploit this knowledge and get even better estimates of the drift rate. The real problems, however, seem to require the estimate of several parameters in addition to the drift rate, and it is not appropriate to just ignore unknown model parameters.

To this point, we have considered only three, rather ideal oscillator models, and seldom does one encounter such simplicity. Typical models for commercial cesium beam frequency standards include white FM, random walks FM, and frequency drift. Unfortunately, none of the three estimation routines discussed above are appropriate to such a model. This problem has been solved in some of the recent work of Jones and Tryon^{[5],[6]}. Their estimation routines are based on Kalman Filters and maximum likelihood estimators and these methods are appropriate for the more complex models. For details, the reader is referred to the works of Jones and Tryon.

Still left untreated are the models which, in addition to drift and other noises, incorporate flicker noises, either in PM or FM or both. In principle, the methods of Jones and Tryon could be applied to Kalman Filters which incorporate empirical flicker models^[7]. To the author's knowledge, however, no such analyses have been reported.

VII. CONCLUSIONS

There are two primary conclusions to be drawn:

- (1) The estimation of the linear frequency drift rates of oscillators and the inclusion of realistic confidence intervals for these

estimates are critically dependent on the adequacy of the model used and, hence the adequacy of the analysis procedures.

- (2) The estimation of the drift rate must be carried along with the estimation of any and all other model parameters which are not known precisely from other considerations (e.g., initial frequency and time offsets, phase noise types, etc.)

More and more, scientists and engineers require clocks which can be relied on to maintain accuracy relative to some master clock. Not only is it important to know that on the average the clock runs well, but it is essential to have some measure of time imprecision as the clock ages. For example, the uncertainties might be expressed as, say,, 90% certain that the clock will be within 5 microseconds of the master two weeks after synchronization. Such measures are what statisticians call "interval estimates" (in contrast to point estimates) and their estimations require interval estimates of the clock's model parameters. Clearly, the parameter estimation routines must be reliable and based on sound measurement practices. Some inappropriate estimation routines can be applied to clocks and oscillators and give dangerously optimistic forecasts of performance.

APPENDIX A

**REGRESSION ANALYSIS
(Equally Spaced Data)**

We begin with the continuous model equation:

$$X(t) = a + b \cdot t + c \cdot t^2 + \phi(t) \tag{A1}$$

We assume that the data is in the form of discrete readings of the dependent variable $X(t)$ at the regular intervals given by:

$$t = n\tau_0 \tag{A2}$$

Equation (A1) can then be written in the obvious form:

$$X_n = a + \tau_0 b \cdot n + \tau_0^2 c \cdot n^2 + \phi(n\tau_0) \tag{A3}$$

for $n = 1, 2, 3 \dots, N$.

Next, we define the matrices:

$$\underline{N} = \begin{bmatrix} 1 & 1 & 1 \\ & 1 & 2 & 4 \\ & & 1 & 3 & 9 \\ & & & \dots & \dots & \dots \\ & & & & \dots & \dots \\ & & & & & 1 & N & N^2 \end{bmatrix} \quad \underline{X} = \begin{bmatrix} X_1 \\ X_2 \\ X_3 \\ \dots \\ X_N \end{bmatrix}$$

$$\underline{T} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \tau_0 & 0 \\ 0 & 0 & \tau_0^2 \end{bmatrix}$$

$$(\underline{NT})' \underline{X} = \begin{bmatrix} \sum_1^N X_n \\ \tau_0 \sum_1^N X_n n \\ \tau_0^2 \sum_1^N X_n n^2 \end{bmatrix} = \underline{T} \cdot \begin{bmatrix} \sum_1^N X_n \\ \sum_1^N X_n n \\ \sum_1^N X_n n^2 \end{bmatrix} = \underline{T} \cdot \begin{bmatrix} S_x \\ S_{nx} \\ S_{n^2x} \end{bmatrix}$$

where four quantities must be calculated from the data:

$$S_x = \sum_{n=1}^N X_n \quad S_{nnx} = \sum_{n=1}^N X_n n^2$$

$$S_{nx} = \sum_{n=1}^N X_n n \quad S_{xx} = \sum_{n=1}^N X_n^2$$

Define

$$\underline{B} = \begin{bmatrix} a \\ b \\ c \end{bmatrix}$$

With these definitions, Eq. A3 can be rewritten in the matrix form:

$$\underline{X} = \underline{N} \underline{T} \underline{B} + \underline{\epsilon} \quad (\text{A4})$$

and the coefficients, B, which minimize the squared errors are given by:

$$\underline{\hat{B}} = \begin{bmatrix} \hat{a} \\ \hat{b} \\ \hat{c} \end{bmatrix} = \underline{T} \cdot (\underline{N}'\underline{N})^{-1} \cdot (\underline{N}'\underline{X}) \quad (\text{A5})$$

The advantage of evenly spaced data for these regressions is that, with a bit of algebra, the matrix, $(\underline{N}'\underline{N})^{-1}$, can be written down in closed form:

$$(\underline{N}'\underline{N})^{-1} = \begin{bmatrix} A & B & C \\ B & D & E \\ C & E & F \end{bmatrix} \cdot 1 / G \quad (\text{A6})$$

where

$$* A = 3 [3 (N + 1) + 2]$$

$$B = -18 (2N + 1)$$

$$C = 30$$

$$D = 12 (2N + 1) (8N + 11) / [(N + 1) (N + 2)]$$

$$E = -180 / (N + 2)$$

$$F = 180 / [(N + 1) (N + 2)]$$

* See Appendix Note # 40

and

$$G = N (N - 1) (N - 2)$$

Also, the inverse of T is just:

$$\underline{T}^{-1} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1/\tau_0 & 0 \\ 0 & 0 & 1/\tau_0^2 \end{bmatrix}$$

The complete solution for the regression parameters can be summarized as follows:

There are four quantities which must calculate from the data:

$$S_x = \sum_{n=1}^N X_n \quad S_{nnx} = \sum_{n=1}^N X_n n^2$$

$$S_{nx} = \sum_{n=1}^N X_n n \quad S_{xx} = \sum_{n=1}^N X_n^2$$

for $n = 1, 2, 3, \dots, N$. Based on these four quantities, the regression parameters are calculated from the seven following equations:

$$\hat{a} = (A S_x + \tau_0 B S_{nx} + \tau_0^2 C S_{nnx}) / G$$

$$\hat{b} = (B S_x + \tau_0 D S_{nx} + \tau_0^2 E S_{nnx}) / (G \tau_0)$$

$$\hat{c} = (C S_x + \tau_0 E S_{nx} + \tau_0^2 F S_{nnx}) / (G \tau_0^2)$$

$$\hat{\sigma}^2 = (S_{xx} - \hat{a} S_x - \tau_0 \hat{b} S_{nx} - \tau_0^2 \hat{c} S_{nnx}) / (N - 3)$$

$$\hat{\sigma}_a^2 = \hat{\sigma}^2 A / G$$

$$\hat{\sigma}_b^2 = \hat{\sigma}^2 D / (G \tau_0^2)$$

$$\hat{\sigma}_c^2 = \hat{\sigma}^2 F / (G \tau_0^4)$$

where the coefficients A, B, C, etc., are given by:

$$A = 3 [3N (N + 1) + 2]$$

$$B = -18 (2N + 1)$$

$$C = 30$$

$$D = 12 (2N + 1) (8N + 11) / [(N + 1) (N + 2)]$$

$$E = -180 / (N + 2)$$

$$F = 180 / [(N + 1) (N + 2)]$$

and

$$G = N (N - 1) (N - 2).$$

In matrix form, the error variance for forecast values is:

$$\text{Var} (\hat{X}_k) = \hat{\sigma}^2 [1 + \underline{N}_k' (\underline{N}'\underline{N})^{-1} \underline{N}_k]$$

where $\underline{N}_k' = [1 \ n_0 \ n_0]$ and $\tau_0 \ n_0$ is the date for the forecast point, \hat{X}_k . That is, $n_0 = N + K$ and K is the number of lags past the last data point at lag N .

APPENDIX B

REGRESSIONS ON LINEAR AND CUBIC FUNCTIONS

The matrixes $(N'N)^{-1}$ for the linear fit and cubic fit, which correspond to Eq. A6 in Appendix A are as follows:

For the linear fit:

$$(N'N)^{-1} = \begin{bmatrix} A & B \\ B & C \end{bmatrix} 1 / D$$

where

$$A = 2 (2N + 1)$$

$$B = -6$$

$$C = 12 / (N + 1)$$

and

$$D = N (N - 1)$$

For the cubic fit:

$$(N'N)^{-1} = \begin{bmatrix} A & B & C & D \\ B & E & F & G \\ C & F & H & I \\ D & G & I & J \end{bmatrix} 1 / K$$

where

$$A = 8 (2N + 1) (N + N + 3)$$

$$B = -20 (6N^2 + 6N + 5)$$

$$C = 120 (2N + 1)$$

$$D = -140$$

$$E = 200 (6N^4 + 27 N^3 + 42 N^2 + 30 N + 11) / L$$

$$F = -300 (N + 1) (3N + 5) (3N + 2) / L$$

$$G = 280 (6N^2 + 15N + 11) / L$$

$$H = 360 (2N + 1) (9N + 13) / L$$

$$I = -4200 (N + 1) / L$$

$$J = 2800 / L$$

and

$$K = N (N - 1) (N - 2) (N - 3)$$

$$L = (N + 1) (N + 2) (N + 3).$$

The restrictions on these equations are that the data is evenly spaced beginning with $n = 1$ to $n = N$, and no missing values. For error estimates (and their distributions) to be valid, the residuals must be random, uncorrelated, (i.e., white).

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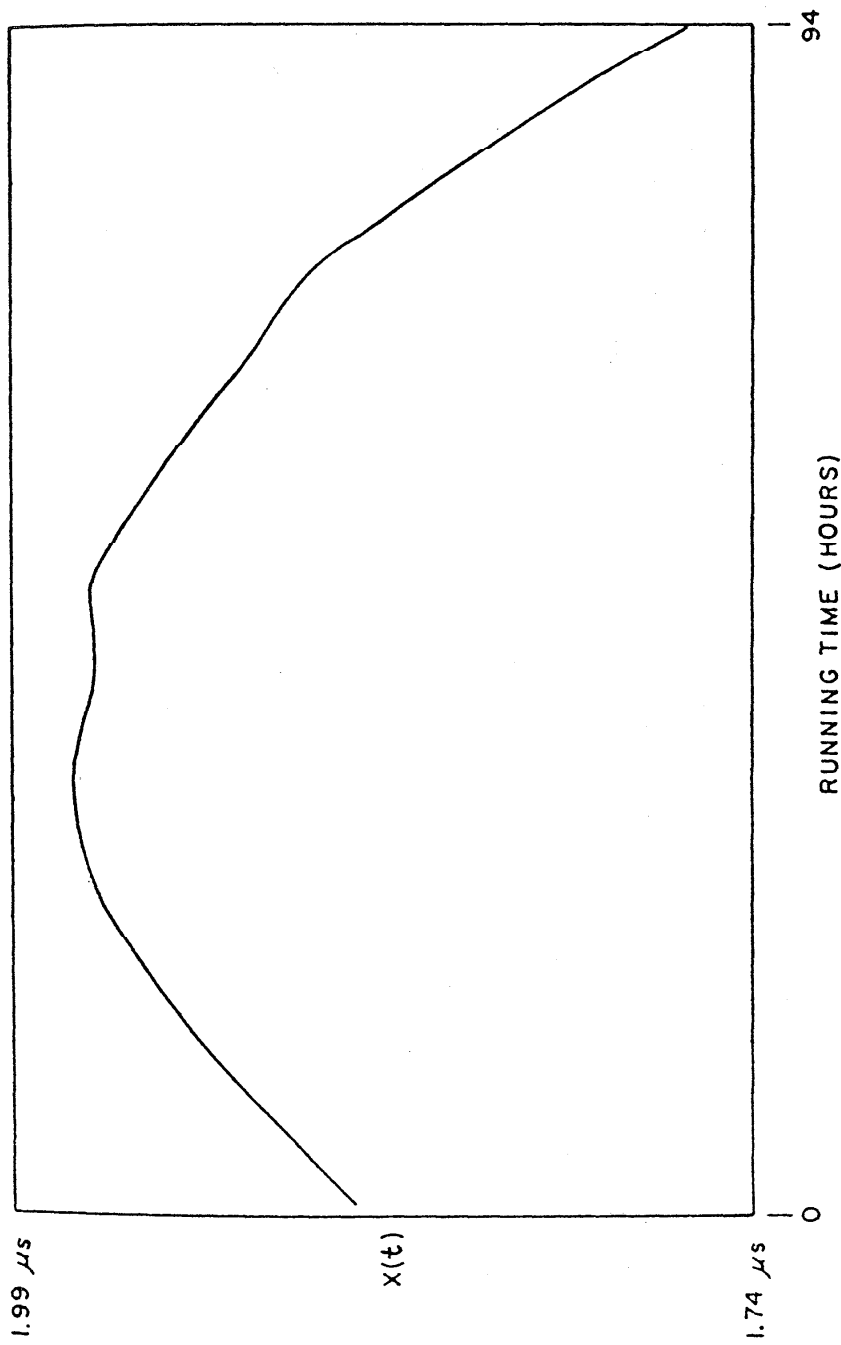


Fig. 1 - Phase difference

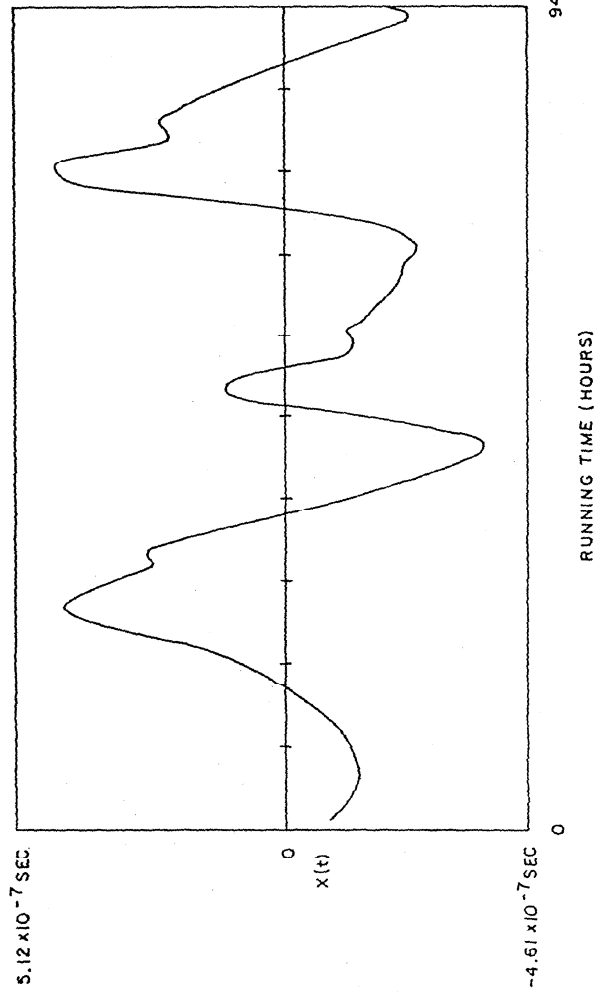
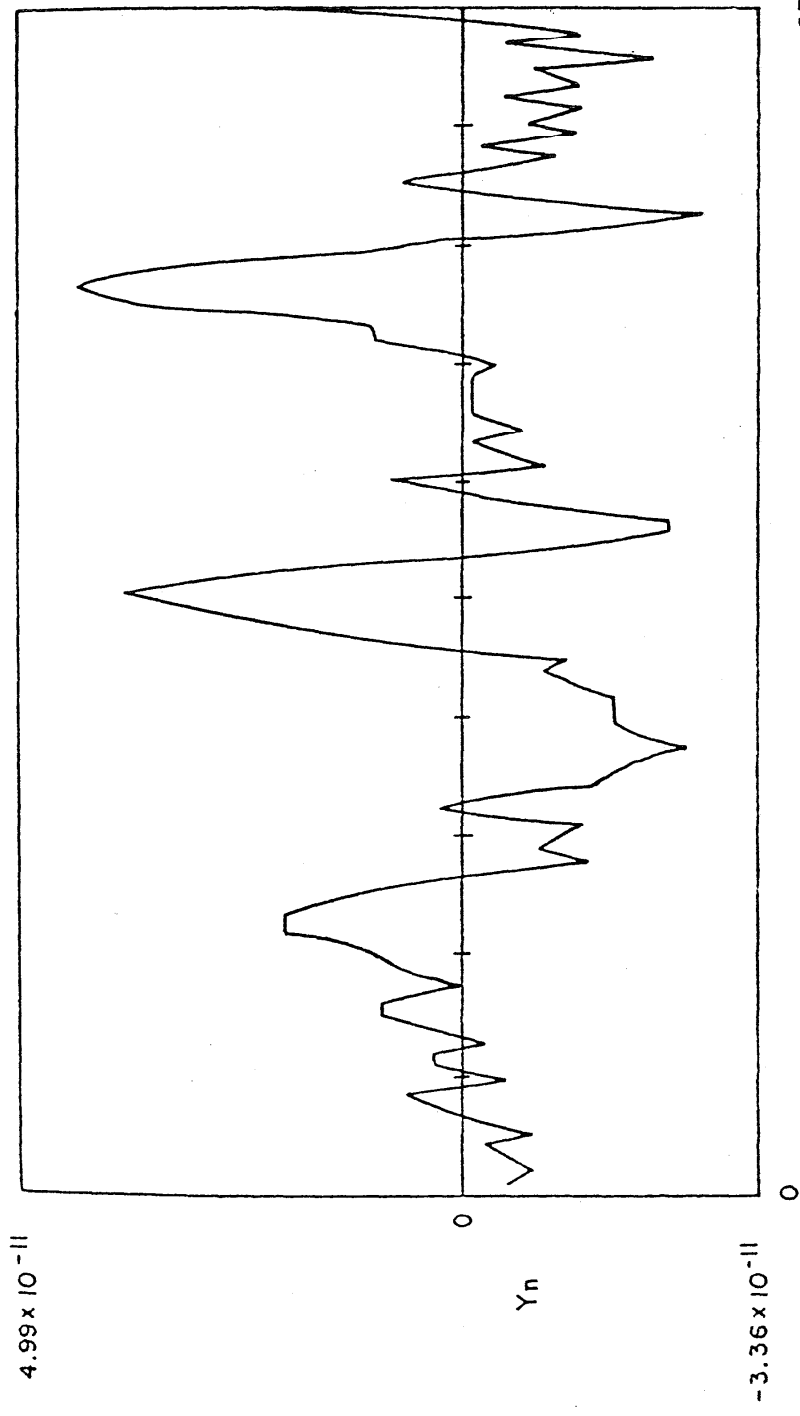


Fig. 2 - Residuals from quadratic fit to phase



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RUNNING TIME (HOURS)

Fig. 3 - Frequency residual
(1st diff-lin)

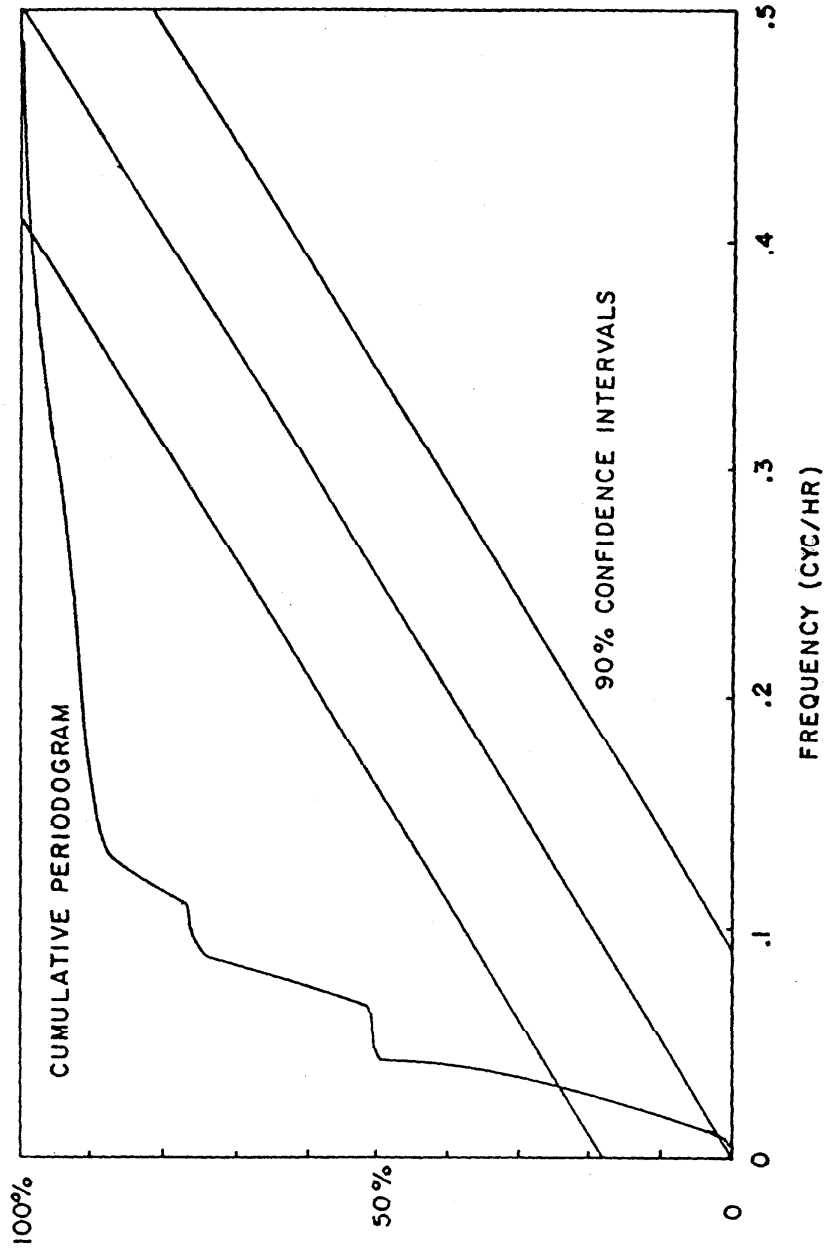


Fig. 4 - First diff - linear

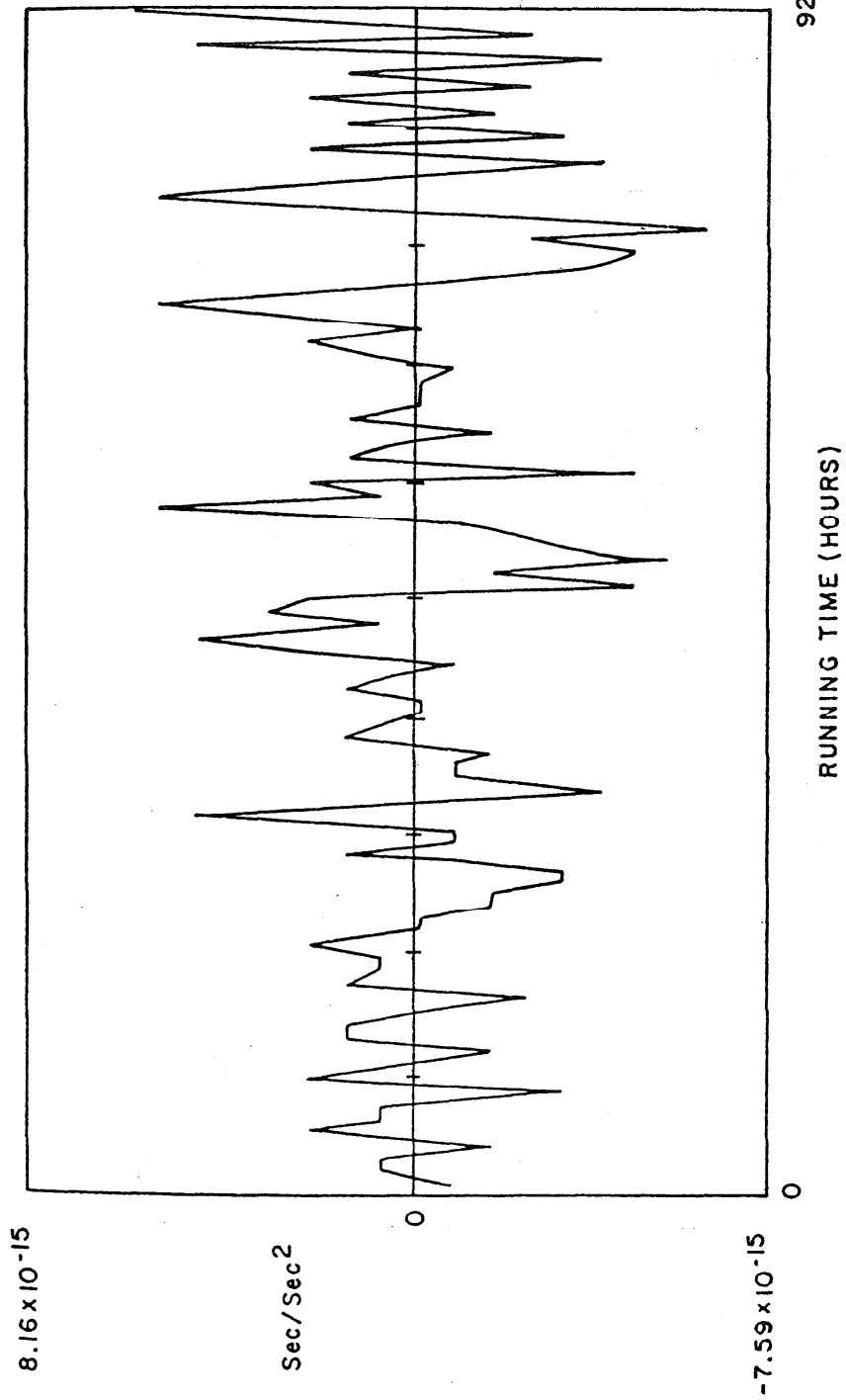


Fig. 5 - Second diff of phase

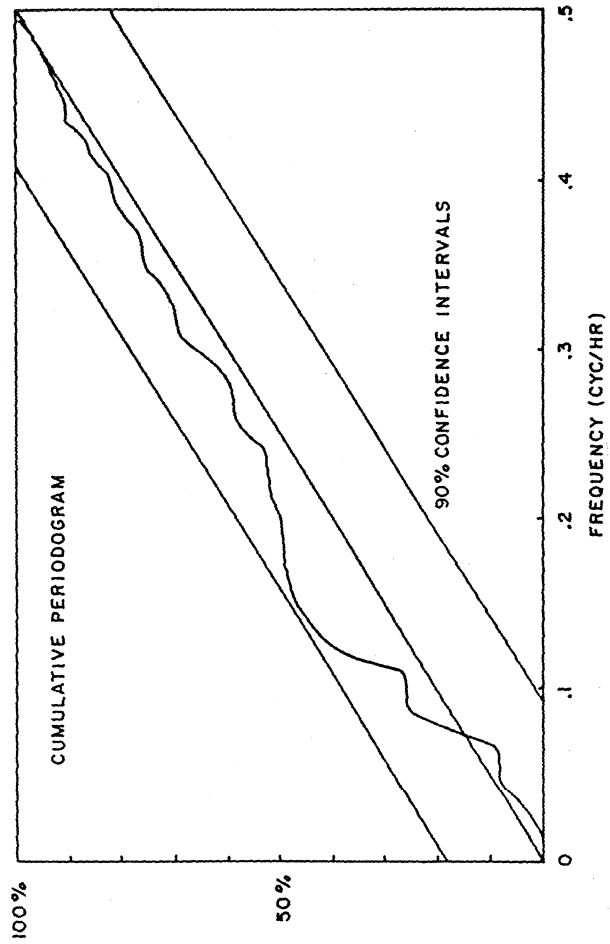
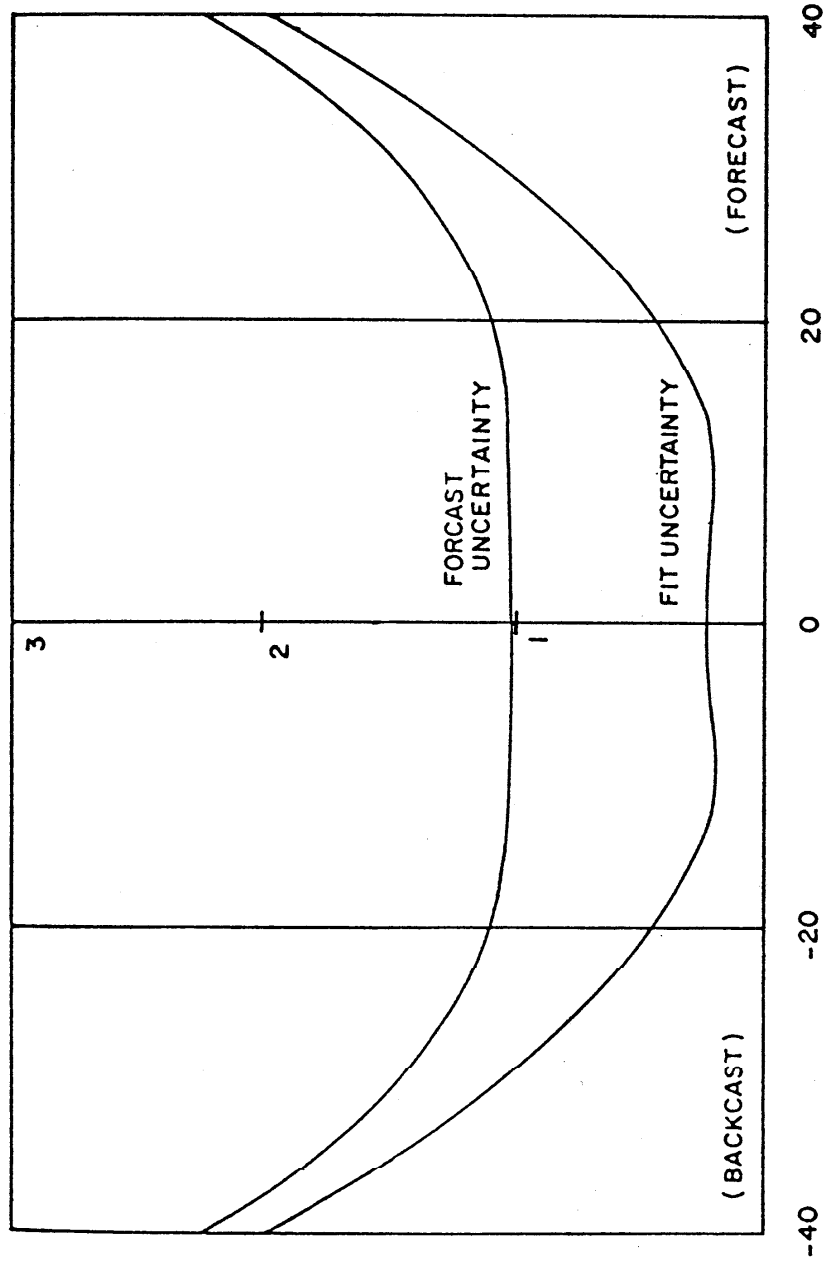
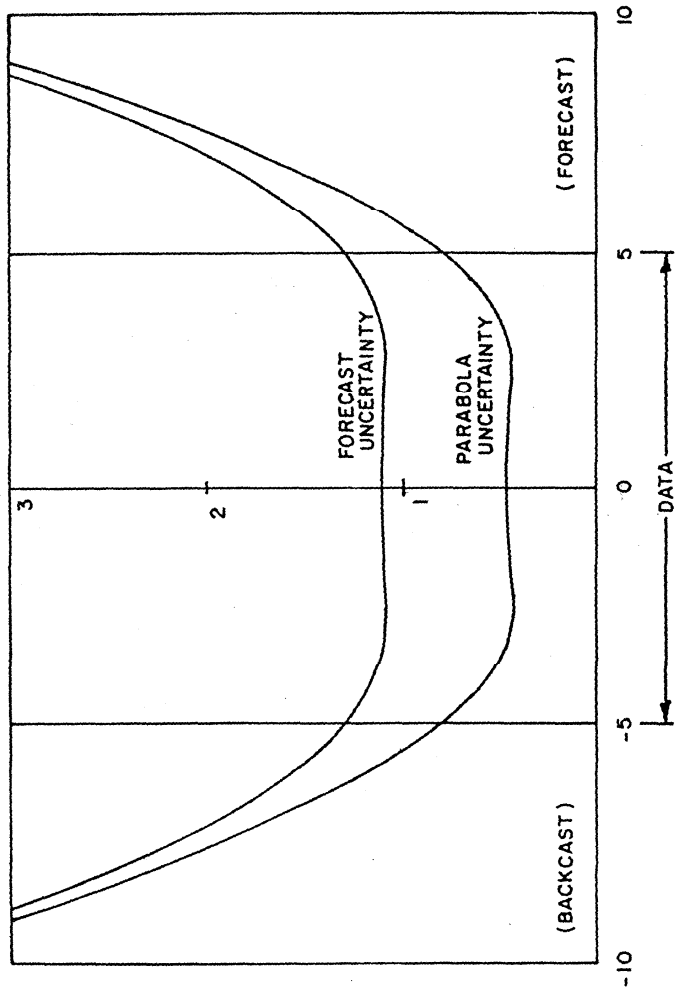


Fig. 6 ~ Second diff ~ avg.

FREQUENCY DRIFT AND WHITE PM
STANDARD DEVIATIONS



FREQUENCY DRIFT AND WHITE PM
STANDARD DEVIATIONS



QUESTIONS AND ANSWERS

MR. ALLAN:

We have found the second difference drift estimator to be useful for the random walk FM process. One has to be careful when you apply it, because if you use overlapping second differences, for example, if you have a cesium beam, and you have white noise FM out to several days and then random walk FM beyond that point, and you are making hourly or 2 hour measurements, if you use the mean of the second differences, you can show that all the middle terms cancel, and in fact you are looking at the frequency at the beginning of the run and the frequency at the end of the run to compute the frequency drift and that's very poor.

DR. BARNES:

My comment would be: There again the problem is in the model, and not in the arithmetic. The models applied here were 3 very simple models, very simple, simpler than you will run into in life. It was pure random walk plus a drift or pure frequency noise plus a drift, or pure random walk of frequency noise plus a drift and it did not approach at all any of the noise complex models where you would have both white frequency and random walk frequency and a drift. That's got to be handled separately, I'm not even totally sure how to perform it in all cases at this point.

MR. McCASKILL:

I would like to know if you would comment on the value of the sample time and the reason why, of course, is that the mean second difference does depend on the sample time? For instance, if you wanted to estimate the aging rate or change in linear change of frequency, what value of sample time would you use in order to make that correction. So, really, the question is, how does the sample time enter into your calculations?

DR. BARNES:

At least I will try to answer in part. I don't know in all cases, I'm sure, but if you have a complex noise process where you have at short term different noise behavior than in long term, it may benefit you to take a longer sampling time and effectively not look at the short term, and then one of the simple models might apply. I honestly haven't looked in great detail at how to choose the sample time. It is an interesting question.

MR. McCASKILL:

Well, let me go further, because we had the benefit of being out at NBS and talking with Dr. Allan earlier and he suggested that we use the mean second difference, and the only problem is if we want to calculate or correct for our aging rate at a tau of five days or ten days, and you calculate the mean second difference, you come up with exactly the Allan Variance. What appears is that in order to come up with the number for the aging rate, you have to calculate the mean second difference using a sample time whenever you take your differences of longer than, let's say, ten days sample time. So we use, of course, the regression model, but we use on the order of two or three weeks in order to calculate an aging rate correction for, say, something like the rubidium in the NAVSTAR 3 clock. It looks like in order to come up with a valid value for the Allan Variance of five or ten day sample time you have to calculate the aging rate at a longer, maybe two or three times longer sample time.

DR. BARNES:

Dave Allan, do you think you can answer that?

MR. ALLAN:

Not to go into details, but if you assume that in the longer term you have random walk frequency modulation as the predominant noise process in a clock, which seems to be true for rubidium, cesium and hydrogen, you can do a very simple thing. You can take the full data length and take the time at the beginning, the time in the middle and the time in the end and construct a second difference, and that's your drift.

There is still the issue of the confidence interval on that. If you really want to verify your confidence interval, you have to have enough data to do a regression. You need enough data to test to be sure the model is good.

DR. WINKLER:

That argument is fourteen years old, because we have been criticized here. I still believe that for a practical case where you depend on measurements which are contaminated and maybe even contaminated by arbitrarily large errors if they are digital. These, theoretical advantages that you have outlined may not be as important as the benefits which you get when you make a least square regression. In this case you can immediately identify the wrong data. Otherwise, if you put the data into an algorithm you may not know how much your data is contaminated. So for practical applications, the first model even though theoretically it is poor, it still gives reasonable estimates of the drift and you have residuals which let you identify wrong phase values immediately.

DR. BARNES:

I think that's true and you may have different reasons to do regression analysis, if your purpose is to measure a drift and understand the confidence intervals, then I think what has been presented is reasonable, if you have as your purpose to look--to see if there are indications of funny behavior in a curve that has such strong curvature or drift that you can't get it on graph paper without doing that, I think it is a very reasonable thing to do. I think looking at the data is one of the healthiest things any analyst can do.

VARIANCES BASED ON DATA WITH DEAD TIME BETWEEN THE MEASUREMENTS

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The accepted definition of frequency stability in the time domain is the two-sample variance (or Allan variance). It is based on the measurement of average frequencies over adjacent time intervals, with no "dead time" between the intervals. The primary advantages of the Allan variance are that (1) it is convergent for many encountered noise models for which the conventional variance is divergent; (2) it can distinguish between many important and different spectral noise types; (3) the two-sample approach relates to many practical implementations; for example, the rms change of an oscillator's frequency from one period to the next; and (4) Allan variances can be easily estimated at integer multiples of the sample interval.

In 1974 a table of bias functions which related variance estimates with various configurations of number of samples and dead time to the Allan variance was published [1]. The tables were based on noises with pure power-law spectral densities.

Often situations occur that unavoidably have dead time between measurements, but still the conventional variances are not convergent. Some of these applications are outside of the time-and-frequency field. Also, the dead times are often distributed throughout a given average, and this distributed dead time is not treated in the 1974 tables.

This paper reviews the bias functions $B_1(N,r,\mu)$, and $B_2(r,\mu)$ and introduces a new bias function, $B_3(2,M,r,\mu)$, to handle the commonly occurring cases of the effect of distributed dead time on the computed variances. Some convenient and easy-to-interpret asymptotic limits are reported. A set of tables for the bias functions are included at the end of this paper.

Key words: Allan variance; bias functions; data sampling and dead time; dead time between the measurement; definition of frequency stability; distributed dead time; two-sample variance

1. Introduction

The sample mean and variance indicate respectively the approximate magnitude of a quantity and its uncertainty. For many situations a continuous function of time is sampled, or measured, at fairly regular intervals. Sampling is not always instantaneous. It takes a finite time and provides an "average reading." If the underlying process (or noise) is random and uncorrelated in time, then the fluctuations are said to be "white" noise. In this situation, the sample mean and variance calculated by the conventional formulas,

$$m = \frac{1}{N} \sum_{n=1}^N \bar{y}_n, \tag{1}$$
$$s^2 = \frac{1}{N-1} \sum_{n=1}^N (\bar{y}_n - m)^2,$$

provide the needed information. The "bar" over the y in eq (1) above denotes the average over a finite time interval. In time and frequency work, y is defined as the average fractional (or normalized) frequency deviation from nominal over an interval τ and at some specified measurement time. As in science generally, the physical model determines the appropriate mathematical model. For the white noise model, the sample mean and variance are the mainstays of most analyses.

Although white noise is a common model for many physical processes, more general noise models are being identified and used. In precise time and frequency measurement, for example, there are two quantities of great interest: instantaneous frequency and phase. These two quantities by definition are exactly related by a differential. (We are NOT considering Fourier frequencies at this point.) That is, the instantaneous frequency is the time rate of change of phase. Thus, if we were employing a model of white frequency-modulation (white FM) noise, then the phase noise is the integral of the white FM noise, commonly called a Brownian motion or random walk. Therefore, depending on whether we are currently interested in phase or frequency, the sample mean and variance may or may not be appropriate.

By definition, white noise has a power spectral density (PSD) that is constant with Fourier frequency. Since random walk noise is the integral of white noise, the power spectral density of a random walk varies as $1/f^2$ (where f is the Fourier frequency) [2]. We encounter noise models whose power spectral densities are various power laws of their Fourier frequencies. Flicker noise is very common and is defined as a noise whose power spectral density varies as $1/f$ over a relevant spectral range. If an oscillator's instantaneous frequency is well modeled by flicker noise, then its phase would be the integral of the flicker noise. It would have a PSD which varied as $1/f^3$.

Noise models whose PSD's are power laws of the Fourier frequency but not integer exponents are possible as well but not as common. This paper considers power-law PSD's of a quantity $y(t)$; $y(t)$ is a continuous sample function which can be measured at regular intervals. For noises whose PSD's vary as f^α with $\alpha < -1$ at low frequencies, the conventional sample mean and variance given in eq (1) do not converge as N gets large [2, 3]. This lack of convergence renders the sample mean and variance ineffective and often misleading in some situations.

Although the sample mean and variance have limitations, other time-domain statistics can be convergent and quite useful. The quantities that we consider in this paper depend significantly on the details of the sampling procedures. Indeed, each sampling scheme has its own bias, and this is the motivation for the bias functions discussed in this paper.

2. The Allan Variance

Recognizing that for particular types of noise, the conventional sample variance fails to converge as the number of samples, N , grows, Allan suggested that we set $N = 2$ and average many of these two-sample variances to get a convergent and stable measure of the spread of the quantity in question [3]. This is what has come to be called the Allan variance.

More specifically let us consider a sample function of time as indicated in figure 1. A measurement consists of averaging $y(t)$ over the interval τ . The next measurement begins at a time T after the beginning of the previous measurement interval. There is no logical reason why T must be as large as τ or larger--if $T < \tau$, then the second measurement begins before the first is completed, which is unusual but possible. When $T = \tau$, there is no dead time between measurements.

The accepted definition of the Allan variance is the expected value of a two-sample variance with no dead time between successive measurements. In symbols, the Allan variance is given by

$$\sigma_y^2(\tau) = \frac{1}{2}E[(\bar{y}_{n+1} - \bar{y}_n)^2], \quad (2)$$

where there is no dead time between the two sample averages for the Allan variance and the $E[\cdot]$ denotes the expectation operator.

3. The Bias Function $B_1(N, r, \mu)$

Define N to be the number of sample averages of $y(t)$ used in eq (1) to estimate a sample variance ($N = 2$ for an Allan variance). Also define r to be the ratio of T to τ ($r = 1$ when there is no dead time between measurements). The parameter μ is related to the exponent of the power law of the PSD of the process $y(t)$. If α is the exponent in the power-law spectrum for $y(t)$, then the Allan variance varies as τ raised to the μ power, where α and μ are related as shown in figure 2 [2-4]. We can use estimates of μ to infer α , the spectral type. The ambiguity in α for $\mu = -2$ has been resolved by using a modified $\sigma_y^2(\tau)$ [5-7].

Often data cannot be taken without dead time between sample averages, and it is useful to consider other than two-sample variances. We will define the bias function $B_1(N, r, \mu)$ by the ratio,

$$B_1(N, r, \mu) = \frac{\sigma^2(N, T, \tau)}{\sigma^2(2, T, \tau)}, \quad (3)$$

where $\sigma^2(N, T, \tau)$ is the expected sample variance given in eq (1) and based on N measurements at intervals T and averaged over a time τ and $r = T/\tau$. In words, $B_1(N, r, \mu)$ is the ratio of the expected variance for N measurements to the expected variance for two samples (everything else held constant). The variances on the right in eq (3) depend implicitly on the noise type even though μ or α are not shown as independent variables. The noise-type parameter, μ , is shown as an independent variable for all of the bias functions in this paper, because the values of the ratio of these variances explicitly depend on μ as will be derived later in the paper. Allan showed that if N and r are held constant, then the α, μ relationship shown in figure 2 is the same; that is, we can still infer the spectral type from the τ dependence using the equation $\alpha = -\mu - 1$, $-2 \leq \mu < 2$ [3].

4. The Bias Function $B_2(r, \mu)$

The bias function $B_2(r, \mu)$ is defined in [1] by the relation,

$$B_2(r, \mu) = \frac{\sigma^2(2, T, \tau)}{\sigma^2(2, \tau, \tau)} = \frac{\sigma^2(2, T, \tau)}{\sigma_y^2(\tau)}. \quad (4)$$

In words, $B_2(r, \mu)$ is the ratio of the expected two-sample variance with dead time to that without dead time (with $N = 2$ and τ the same for both variances). A plot of the $B_2(r, \mu)$ function is shown in figure 3. The bias functions B_1 and B_2 represent biases relative to $N = 2$ rather than infinity; that is, the ratio of the N sample variance (with or without dead time) to the Allan variance and the ratio of the two-sample dead-time variance to the Allan variance respectively.

5. The Bias Function $B_3(N, M, r, \mu)$

Consider the case where a great many measurements are available with dead time between each pair of measurements ($T_0 > \tau_0$). The measurements are averaged over the time interval τ_0 , the spacing between the beginning of one measurement to the next is T_0 , and it may not be convenient to retake the data. We might want to estimate the Allan variance at, say, multiples M of

the averaging time τ_0 . If we average groups of the measurements of $y(t)$, then the dead times between the original measurements are distributed periodically throughout the new average measurements (see figure 4). Define

$$\bar{y}_i = \frac{1}{M} \sum_{n=i}^{M+i-1} \bar{y}_n, \quad (5)$$

where \bar{y}_i are the raw or original measurements based on dead time $T_0 - \tau_0$.

Also define the two-sample variance with distributed dead time as

$$\sigma^2(2, M, T, \tau) = \frac{1}{2} E[(\bar{y}_i - \bar{y}_{i+M})^2], \quad (6)$$

with $\tau = M\tau_0$ and $T = MT_0$.

We can now define B_3 as the ratio of the N-sample variance with distributed dead time to the N-sample variance with dead time accumulated at the end as in figure 1:

$$B_3(N, M, r, \mu) = \frac{\sigma^2(N, M, T, \tau)}{\sigma^2(N, T, \tau)}. \quad (7)$$

Although $B_3(N, M, r, \mu)$ is defined for general N, the tables in the Appendix confine treatment to the case where $N = 2$. There is little value in extending the tables to include general N. Though the variances on the right in eq (7) depend explicitly on N, T and τ , the ratio $B_3(N, M, r, \mu)$ depends on the ratio $r = T/\tau$, and on μ as developed later in this paper.

In words, $B_3(2, M, r, \mu)$ is the ratio of the expected two-sample variance with periodically distributed dead time, as shown in figure 4, to the expected two-sample variance with all the dead time grouped together as shown in figure 1. Both the numerator and the denominator have the same total averaging time and dead time, but they are apportioned differently. The product $B_2(r, \mu) \cdot B_3(2, M, r, \mu)$ is the distributed dead-time variance over the Allan variance for a particular T, τ , M and μ .

Some useful asymptotic forms of B_3 can be found. In the case of large M and $M > r$, we may write that

$$B_3 \approx \frac{1 + \mu}{3}, \quad 1 \leq \mu \leq 2, \quad (8)$$

$$B_3 \approx \frac{4 \ln(2)}{2 \ln(r) + 3}, \quad \mu = 0.$$

One simple and important conclusion from these two equations is that for the cases of flicker FM noise and random-walk FM noise, the τ^μ dependence for large τ is the same whether or not there is periodically distributed dead time. The values of the variances differ only by a constant, and in the latter case the constant is 1. This conclusion is also true for white FM noise, and in this case the constant is also 1.

In the cases $r \gg 1$ and $-2 \leq \mu \leq -1$, we may write for the asymptotic behavior of B_3

$$B_3 \approx M^\alpha, \quad \alpha = -\mu - 1, \quad (9)$$

as was determined empirically. In this region of power-law spectrum the B_3 function has an M^α dependence for an f^α spectrum.

6. The Bias Functions

The bias functions can be written fairly simply by first defining the function,

$$F(A) = 2A^{\mu+2} - (A+1)^{\mu+2} - |(A-1)|^{\mu+2}. \quad (10)$$

The bias functions become

$$B_1(N, r, \mu) = \frac{1 + \sum_{n=1}^{N-1} \frac{N-n}{N(N-1)} \cdot F(nr)}{1 + \frac{1}{2} F(r)}, \quad (11)$$

$$B_2(r, \mu) = \frac{1 + \frac{1}{2}F(r)}{2(1-2^\mu)}, \quad (12)$$

as given in [1], and

$$B_3(2, M, r, \mu) = \frac{2M + M \cdot F(Mr) - \sum_{n=1}^{M-1} (M-n)[2F(nr) - F((M+n)r) - F((M-n)r)]}{(M^{\mu+2})[F(r) + 2]}, \quad (13)$$

as indicated in the appendix.

For $\mu = 0$, eqs (11), (12), and (13) are the indeterminate form 0/0 and must be evaluated by l'Hôpital's rule. Special attention must also be given when expressions of the form 0^0 arise. We verified a random sampling of the table entries using noise simulation and Monte Carlo techniques. No errors were detected. The results in this paper differ some from those in [8], which suggests that there may be some mistakes. Tables for the three bias functions are listed at the end of the paper (note that the computer print-out did not have a symbol for Greek mu = μ).

7. Examples of the Use of the Bias Functions

The spectral type, that is, the value of μ , may be inferred by varying τ , the sample time. However, another useful way of determining the value of μ is by using $B_1(N, r, \mu)$ as follows: calculate an estimate of $\sigma_y^2(N, T, \tau)$ and $\sigma_y^2(2, T, \tau)$ and hence $B_1(N, r, \mu)$; then use the tables to infer the value of μ .

Suppose one has an experimental value for $\sigma_y^2(N_1, T_1, \tau_1)$ and its spectral type is known, that is, μ is known. Suppose also that one wishes to know the variance at some other set of measurement parameters, N_2, T_2, τ_2 . An unbiased estimate of $\sigma_y^2(N_2, T_2, \tau_2)$ may be calculated by the equation:

$$\sigma_y^2(N_2, T_2, \tau_2) = \left(\frac{\tau_2}{\tau_1}\right)^\mu \left(\frac{B_1(N_2, r_2, \mu)B_2(r_2, \mu)}{B_1(N_1, r_1, \mu)B_2(r_1, \mu)}\right) \sigma_y^2(N_1, T_1, \tau_1) \quad (14)$$

where $r_1 = T_1/\tau_1$ and $r_2 = T_2/\tau_2$.

Since the time-domain definition for frequency stability is the Allan variance, it behooves us, where possible, to relate other variances to the Allan variance. If we have an N-sample variance on data with dead-time T- τ and we know the power-law spectral type (the value of μ), then we may write

$$\sigma_y^2(\tau) = \frac{\sigma_y^2(N, T, \tau)}{B_1(N, r, \mu) B_2(r, \mu)}. \quad (15)$$

If we have an N-sample variance where each data entry is an average of M samples with distributed dead time, then we may write

$$\sigma_y^2(\tau) = \frac{\sigma_y^2(N, M, T, \tau)}{B_1(N, r, \mu) B_2(r, \mu) B_3(N, M, r, \mu)}. \quad (16)$$

8. Conclusion

For some important power-law spectral density models often used in characterizing precision oscillators ($S_y(f) \sim f^\alpha$, $\alpha = -2, -1, 0, +1, +2$), we have studied the effects on variances when there is dead time between the frequency samples, and the frequency samples are averaged to increase the integration time. Since dead time between measurements is a common problem throughout metrology, the analysis here has broader applicability than just to time and frequency. Specifically, this kind of analysis has been used with gage blocks and standard volt cells--showing that the classical variance may be non-convergent in some cases [9].

Heretofore, the Allan variance has been shown to have some convenient theoretical properties in relation to power-law spectra as the integration or sample time is varied (if $\sigma_y^2(\tau) \sim \tau^\alpha$, then $\alpha = -\mu - 1$, $-2 < \mu \leq 2$). Since $\sigma_y(\tau)$, by definition, is estimated from data with no dead time, the sample or integration time can be unambiguously changed to investigate the τ dependence. From our analysis, we have concluded that for the asymptotic limit of several samples being averaged with dead time present in the data, the τ dependence of the variances is the same. The $\alpha = -\mu - 1$ relationship still remains valid for white FM noise ($\mu = -1$, $\alpha = 0$), flicker FM noise ($\mu = 0$, $\alpha = -1$), and for

random-walk FM noise ($\mu = +1$, $\alpha = -2$). The asymptotic limit is approached as the product of number of samples averaged and the initial data sample time, τ_0 , becomes larger than the dead time ($M > r$). The variances so obtained differ only by a constant, which can be calculated as given in this paper.

A knowledge of the appropriate power-law spectral model is required to translate a distributed dead-time variance to the corresponding value of the Allan variance. In principle, the power-law spectral model can be estimated from the τ^μ dependence, using the variance analysis on the data as outlined above.

9. References

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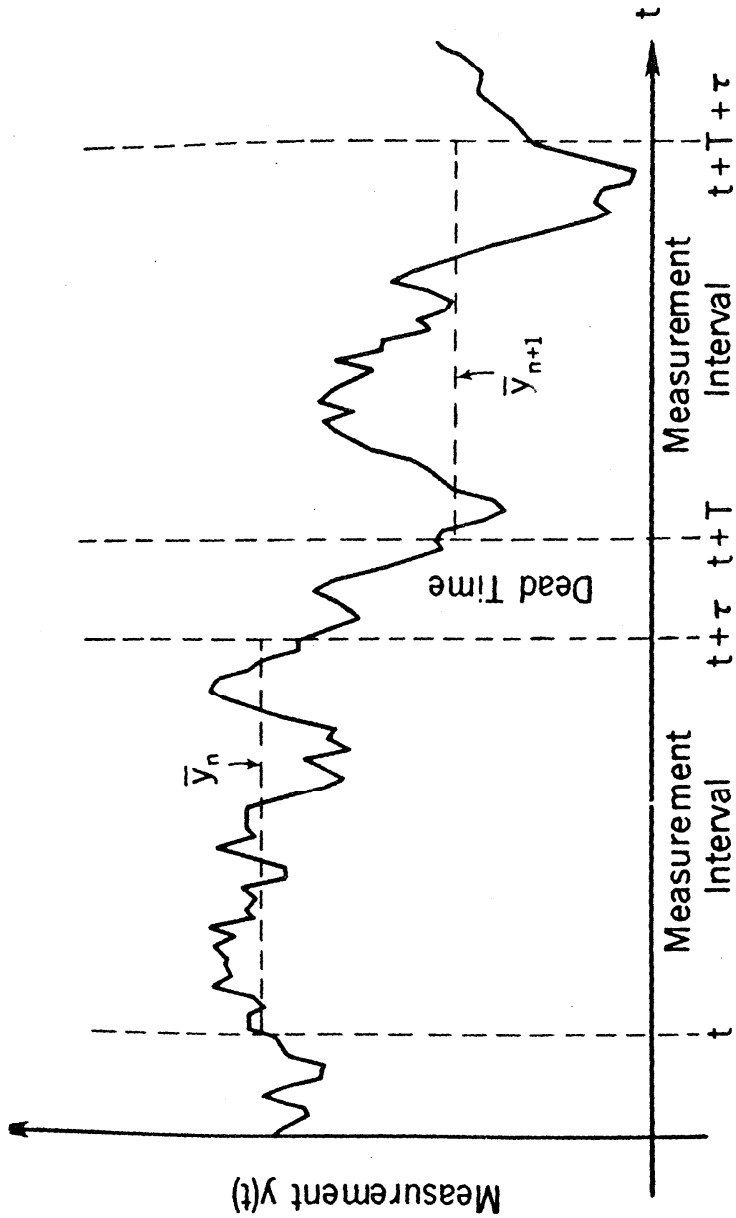
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Table 1. Table of some bias function identities

$B_1(2, r, \mu)$	= 1
$B_1(N, r, 2)$	= $(N(N+1))/6$
$B_1(N, 1, 1)$	= $N/2$
$B_1(N, 1, \mu)$	= $(N(1-N^\mu))/[(2(N-1)(1-2^\mu))]$ for $\mu \neq 0$ = $N \ln(N)/[2(N-1) \ln(2)]$ for $\mu=0$
* $B_1(N, 1, \mu)$	= 1 for $\mu < 0$ = $[2/(N(N-1))] \sum_{n=1}^{N-1} (N-n) \cdot n^\mu$ for $\mu > 0$
$B_1(N, r, -1)$	= 1 if $r \geq 1$
$B_1(N, r, -2)$	= 1 if $r \neq 1$ or 0
$B_2(0, \mu)$	= 0
$B_2(1, \mu)$	= 1
$B_2(r, 2)$	= r^2
$B_2(r, 1)$	= $(3r - 1)/2$ if $r \geq 1$
$B_2(r, -1)$	= r if $0 \leq r \leq 1$ = 1 if $r \geq 1$
$B_2(r, -2)$	= 0 if $r=0$ = 1 if $r=1$ = $2/3$ otherwise
$B_3(2, M, 1, \mu)$	= 1
$B_3(2, M, r, -2)$	= M
$B_3(2, r, \mu)$	= 1
$B_3(2, M, r, 2)$	= 1
$B_3(2, M, r, -1)$	= 1 for $r \geq 1$

* See Appendix Note #41



TWO MEASUREMENTS OF A SET WITH DEAD TIME

Figure 1. Illustration of two fractional frequency samples with dead time, $T-r$, between the samples. This is two of a set of adjacent frequency measurements, each averaged over an interval r , needed to calculate a two-sample variance from a data set.

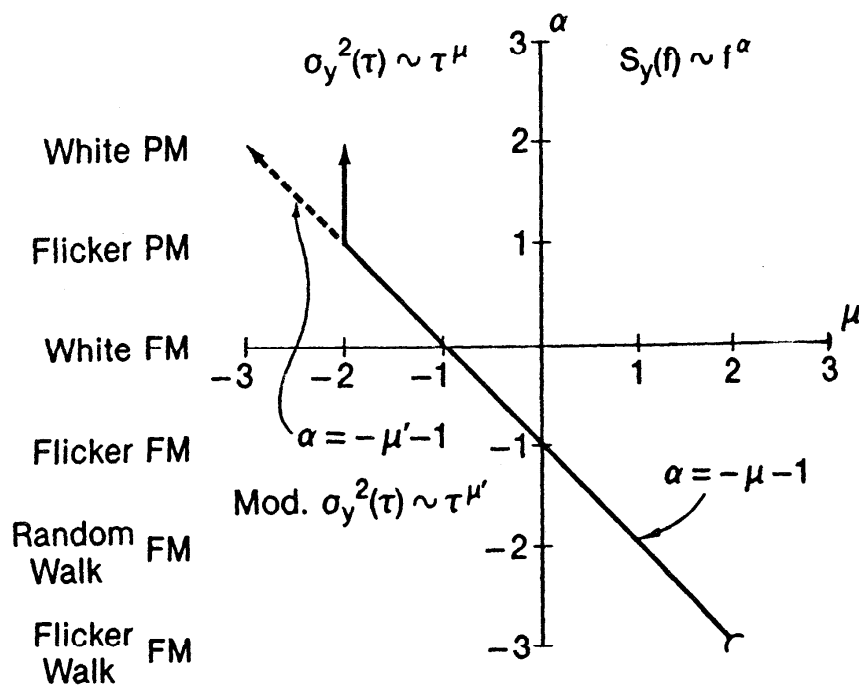
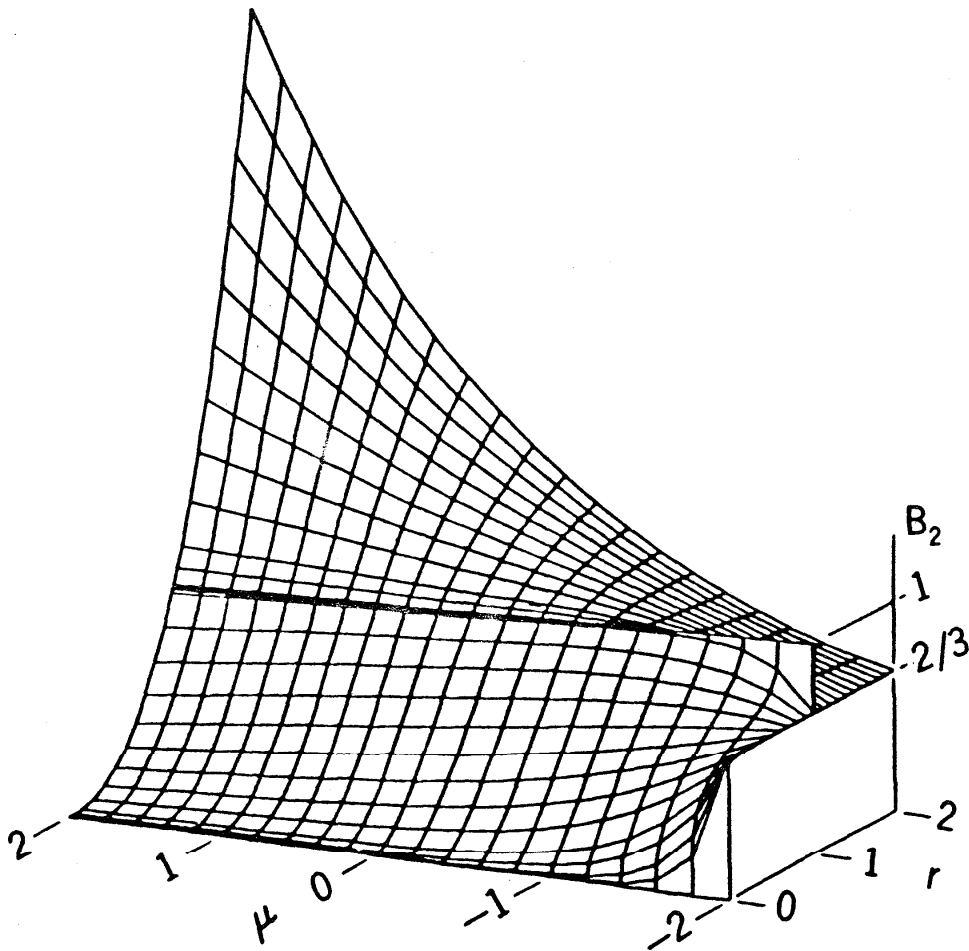


Figure 2. A plot of the relationship between the frequency-domain power-law spectral-density exponent α and the time-domain two-sample Allan variance exponent μ ($\alpha = -\mu - 1$, $-2 \leq \mu < 2$ and $\alpha \geq 1$ for $\mu = -2$). Also shown is the similar relationship between α and the modified Allan variance with exponent on τ of μ' ($\alpha = -\mu' - 1$, $-4 \leq \mu' < 2$). The pointing arrows indicate the mu-alpha relationship (α vs. μ or μ') for which the particular variance applies.



THE BIAS FUNCTION, $B_2(r, \mu)$

Figure 3. A three dimensional plot of the bias function $B_2(r, \mu)$, where $r = T/\tau$, and the dead time is $T - \tau$. The "fin" at $r = 1$ and $\mu = -2$ approaches zero width as the measurement bandwidth approaches infinity (see appendix ref. [3]).

Appendix

With reference to figure 1, the frequency sampling window has an equivalent phase sampling window. The intent is to evaluate the variance, $S(M)$, of the sampled phase function in terms of the phase autocorrelation function, $R(\tau)$. The process here is to correctly account for terms and cross-terms coming from squaring and averaging the samples for each M . The $B_3(2, M, r, \mu)$ function can then be obtained from the relation,

$$B_3(2, M, r, \mu) = \frac{S(M)}{S(1) \cdot M^{\mu+2}},$$

for appropriate M , r , and μ . The denominator is just the two-sample variance with dead time for MT and $M\tau$ (in accordance with the definition of $B_3(2, M, r, \mu)$). The factors common to the numerator and denominator are ignored in the following.

For $M = 1$, the variance $S(1)$ is just

$$S(1) = 4 \cdot R(0) - 4 \cdot R(\tau) - 4 \cdot R(T) + 2 \cdot R(T+\tau) + 2R(T-\tau),$$

where use has been made of the definition of the autocorrelation function,

$$R(T) = E[\phi(t) \cdot \phi(t+T)].$$

It is convenient to define a function $G(T)$ as

$$G(T) = 2 \cdot R(T) - R(T+\tau) - R(T-\tau).$$

Similarly, $S(2)$ can now be written in the form,

$$S(2) = 8 \cdot R(0) - 8 \cdot R(\tau) + 2 \cdot G(T) - 4 \cdot G(2T) - 2 \cdot G(3T).$$

Following this procedure, we can verify that the general $S(M)$ is just

$$S(M) = 4 \cdot M \cdot R(0) - 4 \cdot M \cdot R(\tau) - 2 \cdot M \cdot G(MT)$$

$$+ 2 \sum_{n=1}^{M-1} (M-n) [2 \cdot G(nT) - G((M+n)T) - G((M-n)T)].$$

Following the work of Barnes and Allan [2,3], we can define the function $U(\tau)$ by the relation,

$$U(\tau) = 2 \cdot R(0) - 2 \cdot R(\tau),$$

and also define

$$F(nr) = G(nT)/U(\tau),$$

where $r = T/\tau$. The function $U(\tau)$ for power-law power spectral densities has the form,

$$U(\tau) = \frac{|\tau|^{\mu+2}}{4-2^{\mu+2}},$$

which yields

$$F(nr) = 2 \cdot (nr)^{\mu+2} - (nr+1)^{\mu+2} - |nr-1|^{\mu+2}.$$

Finally, the working relation can be written as

$$B_3(2, M, r, \mu) = \frac{2 \cdot M + M \cdot F(Mr) - \sum_{n=1}^{M-1} (M-n) [2 \cdot F(nr) - F((M+n)r) - F((M-n)r)]}{[2 + F(r)] \cdot M^{\mu+2}}$$

B1(N,r,mu) for r = .01

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.091E+00	1.210E+00	1.360E+00	1.545E+00	1.772E+00	2.089E+00	2.349E+00	2.456E+00	2.494E+00	2.512E+00
-1.6	1.199E+00	1.487E+00	1.893E+00	2.455E+00	3.239E+00	4.440E+00	5.505E+00	6.002E+00	6.198E+00	6.309E+00
-1.4	1.328E+00	1.856E+00	2.688E+00	3.991E+00	6.054E+00	9.500E+00	1.282E+01	1.454E+01	1.532E+01	1.585E+01
-1.2	1.482E+00	2.346E+00	3.880E+00	6.598E+00	1.145E+01	2.026E+01	2.947E+01	3.486E+01	3.754E+01	3.980E+01
-1	1.667E+00	3.000E+00	5.667E+00	1.100E+01	2.167E+01	4.255E+01	6.628E+01	8.190E+01	9.064E+01	1.000E+02
-.8	1.884E+00	3.860E+00	8.303E+00	1.828E+01	4.041E+01	8.691E+01	1.443E+02	1.868E+02	2.137E+02	2.519E+02
-.6	2.134E+00	4.959E+00	1.205E+01	2.977E+01	7.296E+01	1.697E+02	2.995E+02	4.074E+02	4.851E+02	6.423E+02
-.4	2.407E+00	6.279E+00	1.703E+01	4.655E+01	1.247E+02	3.106E+02	5.814E+02	8.357E+02	1.043E+03	1.715E+03
-.2	2.677E+00	7.714E+00	2.296E+01	6.836E+01	1.976E+02	5.220E+02	1.036E+03	1.580E+03	2.084E+03	5.581E+03
0	2.912E+00	9.075E+00	2.909E+01	9.282E+01	2.857E+02	7.951E+02	1.672E+03	2.719E+03	3.826E+03	
.2	3.089E+00	1.018E+01	3.448E+01	1.162E+02	3.763E+02	1.097E+03	2.444E+03	4.263E+03	6.453E+03	
.4	3.203E+00	1.095E+01	3.857E+01	1.354E+02	4.572E+02	1.390E+03	3.287E+03	6.172E+03	1.014E+04	
.6	3.268E+00	1.143E+01	4.133E+01	1.495E+02	5.221E+02	1.649E+03	4.142E+03	8.419E+03	1.514E+04	
.8	3.302E+00	1.170E+01	4.303E+01	1.590E+02	5.708E+02	1.869E+03	4.991E+03	1.104E+04	2.189E+04	
1	3.319E+00	1.185E+01	4.403E+01	1.653E+02	6.064E+02	2.055E+03	5.842E+03	1.412E+04	3.109E+04	
1.2	3.327E+00	1.192E+01	4.461E+01	1.693E+02	6.330E+02	2.215E+03	6.717E+03	1.782E+04	4.383E+04	
1.4	3.330E+00	1.196E+01	4.494E+01	1.720E+02	6.530E+02	2.359E+03	7.640E+03	2.234E+04	6.169E+04	
1.6	3.332E+00	1.198E+01	4.514E+01	1.738E+02	6.688E+02	2.493E+03	8.637E+03	2.793E+04	8.697E+04	
1.8	3.333E+00	1.199E+01	4.526E+01	1.750E+02	6.819E+02	2.623E+03	9.737E+03	3.493E+04	1.231E+05	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = .03

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.091E+00	1.211E+00	1.366E+00	1.573E+00	1.827E+00	1.950E+00	1.995E+00	2.010E+00	2.015E+00	2.016E+00
-1.6	1.199E+00	1.491E+00	1.910E+00	2.532E+00	3.341E+00	3.782E+00	3.963E+00	4.030E+00	4.053E+00	4.064E+00
-1.4	1.329E+00	1.863E+00	2.719E+00	4.133E+00	6.080E+00	7.268E+00	7.815E+00	8.044E+00	8.135E+00	8.191E+00
-1.2	1.484E+00	2.355E+00	3.918E+00	6.768E+00	1.093E+01	1.378E+01	1.525E+01	1.595E+01	1.626E+01	1.651E+01
-1	1.667E+00	3.000E+00	5.667E+00	1.100E+01	1.928E+01	2.560E+01	2.929E+01	3.127E+01	3.229E+01	3.333E+01
-.8	1.878E+00	3.825E+00	8.141E+00	1.753E+01	3.304E+01	4.626E+01	5.499E+01	6.029E+01	6.341E+01	6.770E+01
-.6	2.114E+00	4.838E+00	1.147E+01	2.706E+01	5.441E+01	8.049E+01	9.997E+01	1.134E+02	1.225E+02	1.403E+02
-.4	2.362E+00	6.005E+00	1.566E+01	3.993E+01	8.524E+01	1.335E+02	1.744E+02	2.066E+02	2.314E+02	3.099E+02
-.2	2.604E+00	7.242E+00	2.047E+01	5.581E+01	1.259E+02	2.094E+02	2.897E+02	3.619E+02	4.256E+02	8.571E+02
0	2.819E+00	8.430E+00	2.546E+01	7.348E+01	1.748E+02	3.096E+02	4.568E+02	6.082E+02	7.611E+02	
.2	2.992E+00	9.460E+00	3.013E+01	9.127E+01	2.286E+02	4.324E+02	6.856E+02	9.834E+02	1.328E+03	
.4	3.118E+00	1.027E+01	3.412E+01	1.077E+02	2.835E+02	5.750E+02	9.868E+02	1.541E+03	2.277E+03	
.6	3.203E+00	1.087E+01	3.728E+01	1.220E+02	3.372E+02	7.357E+02	1.376E+03	2.364E+03	3.868E+03	
.8	3.257E+00	1.128E+01	3.966E+01	1.339E+02	3.887E+02	9.152E+02	1.879E+03	3.580E+03	6.554E+03	
1	3.290E+00	1.155E+01	4.141E+01	1.437E+02	4.382E+02	1.117E+03	2.533E+03	5.395E+03	1.114E+04	
1.2	3.309E+00	1.172E+01	4.267E+01	1.518E+02	4.867E+02	1.348E+03	3.392E+03	8.128E+03	1.903E+04	
1.4	3.320E+00	1.184E+01	4.360E+01	1.588E+02	5.352E+02	1.615E+03	4.535E+03	1.228E+04	3.276E+04	
1.6	3.327E+00	1.191E+01	4.431E+01	1.650E+02	5.851E+02	1.930E+03	6.068E+03	1.866E+04	5.684E+04	
1.8	3.331E+00	1.196E+01	4.487E+01	1.706E+02	6.374E+02	2.304E+03	8.142E+03	2.849E+04	9.936E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = .1

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.093E+00	1.226E+00	1.438E+00	1.547E+00	1.579E+00	1.586E+00	1.586E+00	1.585E+00	1.584E+00	1.583E+00
-1.6	1.205E+00	1.522E+00	2.021E+00	2.329E+00	2.449E+00	2.490E+00	2.502E+00	2.505E+00	2.505E+00	2.504E+00
-1.4	1.337E+00	1.907E+00	2.832E+00	3.477E+00	3.772E+00	3.892E+00	3.938E+00	3.954E+00	3.960E+00	3.962E+00
-1.2	1.491E+00	2.396E+00	3.931E+00	5.125E+00	5.751E+00	6.046E+00	6.177E+00	6.235E+00	6.260E+00	6.278E+00
-1	1.667E+00	3.000E+00	5.375E+00	7.429E+00	8.653E+00	9.312E+00	9.652E+00	9.825E+00	9.912E+00	1.000E+01
-.8	1.860E+00	3.720E+00	7.204E+00	1.055E+01	1.281E+01	1.419E+01	1.502E+01	1.550E+01	1.577E+01	1.616E+01
-.6	2.065E+00	4.540E+00	9.423E+00	1.461E+01	1.859E+01	2.137E+01	2.326E+01	2.452E+01	2.537E+01	2.702E+01
-.4	2.273E+00	5.430E+00	1.199E+01	1.971E+01	2.642E+01	3.177E+01	3.592E+01	3.910E+01	4.153E+01	4.922E+01
-.2	2.472E+00	6.344E+00	1.481E+01	2.585E+01	3.673E+01	4.664E+01	5.542E+01	6.314E+01	6.989E+01	1.156E+02
0	2.653E+00	7.236E+00	1.779E+01	3.300E+01	5.003E+01	6.771E+01	8.566E+01	1.037E+02	1.219E+02	
.2	2.810E+00	8.065E+00	2.080E+01	4.111E+01	6.690E+01	9.750E+01	1.331E+02	1.741E+02	2.214E+02	
.4	2.940E+00	8.804E+00	2.376E+01	5.015E+01	8.816E+01	1.397E+02	2.084E+02	2.993E+02	4.194E+02	
.6	3.043E+00	9.445E+00	2.663E+01	6.014E+01	1.149E+02	1.999E+02	3.295E+02	5.265E+02	8.254E+02	
.8	3.123E+00	9.989E+00	2.938E+01	7.118E+01	1.487E+02	2.862E+02	5.269E+02	9.466E+02	1.678E+03	
1	3.184E+00	1.045E+01	3.204E+01	8.347E+01	1.918E+02	4.114E+02	8.523E+02	1.735E+03	3.500E+03	
1.2	3.230E+00	1.084E+01	3.463E+01	9.728E+01	2.470E+02	5.943E+02	1.394E+03	3.232E+03	7.453E+03	
1.4	3.265E+00	1.118E+01	3.721E+01	1.130E+02	3.183E+02	8.636E+02	2.304E+03	6.103E+03	1.613E+04	
1.6	3.293E+00	1.147E+01	3.982E+01	1.309E+02	4.113E+02	1.263E+03	3.843E+03	1.166E+04	3.535E+04	
1.8	3.315E+00	1.174E+01	4.252E+01	1.517E+02	5.330E+02	1.858E+03	6.467E+03	2.250E+04	7.830E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = .3

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.124E+00	1.226E+00	1.259E+00	1.265E+00	1.265E+00	1.263E+00	1.262E+00	1.261E+00	1.261E+00	1.260E+00
-1.6	1.254E+00	1.479E+00	1.564E+00	1.589E+00	1.594E+00	1.594E+00	1.593E+00	1.591E+00	1.591E+00	1.590E+00
-1.4	1.388E+00	1.759E+00	1.925E+00	1.988E+00	2.009E+00	2.014E+00	2.015E+00	2.015E+00	2.014E+00	2.013E+00
-1.2	1.527E+00	2.068E+00	2.351E+00	2.479E+00	2.534E+00	2.556E+00	2.565E+00	2.569E+00	2.570E+00	2.571E+00
-1	1.667E+00	2.405E+00	2.850E+00	3.087E+00	3.209E+00	3.271E+00	3.302E+00	3.318E+00	3.326E+00	3.333E+00
-.8	1.806E+00	2.769E+00	3.435E+00	3.847E+00	4.093E+00	4.239E+00	4.325E+00	4.376E+00	4.405E+00	4.447E+00
-.6	1.943E+00	3.160E+00	4.122E+00	4.804E+00	5.275E+00	5.595E+00	5.812E+00	5.958E+00	6.056E+00	6.252E+00
-.4	2.076E+00	3.577E+00	4.927E+00	6.024E+00	6.889E+00	7.564E+00	8.086E+00	8.487E+00	8.795E+00	9.776E+00
-.2	2.203E+00	4.021E+00	5.877E+00	7.597E+00	9.147E+00	1.053E+01	1.175E+01	1.282E+01	1.377E+01	2.017E+01
0	2.325E+00	4.492E+00	7.000E+00	9.651E+00	1.238E+01	1.515E+01	1.796E+01	2.079E+01	2.363E+01	
.2	2.440E+00	4.994E+00	8.338E+00	1.236E+01	1.710E+01	2.260E+01	2.897E+01	3.632E+01	4.479E+01	
.4	2.550E+00	5.530E+00	9.940E+00	1.599E+01	2.411E+01	3.493E+01	4.926E+01	6.823E+01	9.330E+01	
.6	2.654E+00	6.106E+00	1.187E+01	2.089E+01	3.472E+01	5.581E+01	8.786E+01	1.365E+02	2.103E+02	
.8	2.754E+00	6.729E+00	1.422E+01	2.756E+01	5.098E+01	9.187E+01	1.632E+02	2.874E+02	5.037E+02	
1	2.852E+00	7.406E+00	1.708E+01	3.673E+01	7.618E+01	1.552E+02	3.132E+02	6.292E+02	1.261E+03	
1.2	2.947E+00	8.147E+00	2.059E+01	4.942E+01	1.157E+02	2.678E+02	6.170E+02	1.419E+03	3.261E+03	
1.4	3.042E+00	8.963E+00	2.493E+01	6.709E+01	1.781E+02	4.703E+02	1.241E+03	3.272E+03	8.631E+03	
1.6	3.137E+00	9.868E+00	3.030E+01	9.185E+01	2.774E+02	8.381E+02	2.535E+03	7.673E+03	2.324E+04	
1.8	3.234E+00	1.087E+01	3.699E+01	1.267E+02	4.366E+02	1.511E+03	5.245E+03	1.823E+04	6.343E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

Bi(N,r,mu) for r = 1

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	8.333E-01	7.500E-01	7.083E-01	6.875E-01	6.771E-01	6.719E-01	6.693E-01	6.680E-01	6.673E-01	6.667E-01
-1.8	8.581E-01	7.827E-01	7.431E-01	7.226E-01	7.122E-01	7.068E-01	7.042E-01	7.028E-01	7.021E-01	7.014E-01
-1.6	8.866E-01	8.221E-01	7.864E-01	7.672E-01	7.570E-01	7.517E-01	7.490E-01	7.476E-01	7.468E-01	7.461E-01
-1.4	9.193E-01	8.700E-01	8.410E-01	8.245E-01	8.154E-01	8.105E-01	8.079E-01	8.065E-01	8.058E-01	8.051E-01
-1.2	9.569E-01	9.284E-01	9.105E-01	8.997E-01	8.933E-01	8.897E-01	8.877E-01	8.866E-01	8.860E-01	8.854E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.050E+00	1.088E+00	1.117E+00	1.137E+00	1.150E+00	1.160E+00	1.165E+00	1.169E+00	1.171E+00	1.175E+00
-.6	1.107E+00	1.197E+00	1.271E+00	1.327E+00	1.370E+00	1.401E+00	1.422E+00	1.438E+00	1.448E+00	1.470E+00
-.4	1.172E+00	1.333E+00	1.476E+00	1.599E+00	1.700E+00	1.782E+00	1.847E+00	1.898E+00	1.938E+00	2.065E+00
-.2	1.247E+00	1.502E+00	1.754E+00	1.994E+00	2.216E+00	2.418E+00	2.599E+00	2.759E+00	2.900E+00	3.863E+00
0	1.333E+00	1.714E+00	2.133E+00	2.581E+00	3.048E+00	3.528E+00	4.016E+00	4.509E+00	5.005E+00	
.2	1.432E+00	1.982E+00	2.658E+00	3.471E+00	4.432E+00	5.555E+00	6.858E+00	8.363E+00	1.010E+01	
.4	1.546E+00	2.320E+00	3.391E+00	4.846E+00	6.801E+00	9.407E+00	1.287E+01	1.744E+01	2.350E+01	
.6	1.677E+00	2.750E+00	4.424E+00	7.006E+00	1.096E+01	1.698E+01	2.614E+01	4.005E+01	6.114E+01	
.8	1.827E+00	3.299E+00	5.894E+00	1.045E+01	1.841E+01	3.230E+01	5.652E+01	9.872E+01	1.722E+02	
1	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	
1.2	2.198E+00	4.900E+00	1.104E+01	2.506E+01	5.717E+01	1.308E+02	2.999E+02	6.881E+02	1.580E+03	
1.4	2.426E+00	6.059E+00	1.546E+01	3.999E+01	1.044E+02	2.738E+02	7.202E+02	1.897E+03	5.003E+03	
1.6	2.688E+00	7.555E+00	2.191E+01	6.479E+01	1.938E+02	5.833E+02	1.762E+03	5.331E+03	1.615E+04	
1.8	2.988E+00	9.490E+00	3.138E+01	1.063E+02	3.646E+02	1.260E+03	4.372E+03	1.519E+04	5.286E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

Bi(N,r,mu) for r = 1.01

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.156E-01	8.682E-01	8.425E-01	8.288E-01	8.217E-01	8.181E-01	8.162E-01	8.153E-01	8.148E-01	8.143E-01
-1.6	9.098E-01	8.566E-01	8.264E-01	8.098E-01	8.009E-01	7.963E-01	7.938E-01	7.926E-01	7.920E-01	7.913E-01
-1.4	9.286E-01	8.840E-01	8.573E-01	8.419E-01	8.333E-01	8.287E-01	8.262E-01	8.249E-01	8.242E-01	8.235E-01
-1.2	9.599E-01	9.331E-01	9.160E-01	9.056E-01	8.995E-01	8.960E-01	8.940E-01	8.929E-01	8.924E-01	8.917E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.048E+00	1.085E+00	1.113E+00	1.133E+00	1.147E+00	1.155E+00	1.161E+00	1.165E+00	1.167E+00	1.170E+00
-.6	1.104E+00	1.193E+00	1.265E+00	1.321E+00	1.362E+00	1.393E+00	1.414E+00	1.429E+00	1.440E+00	1.461E+00
-.4	1.168E+00	1.327E+00	1.468E+00	1.589E+00	1.689E+00	1.770E+00	1.835E+00	1.885E+00	1.924E+00	2.050E+00
-.2	1.243E+00	1.495E+00	1.743E+00	1.980E+00	2.200E+00	2.400E+00	2.579E+00	2.737E+00	2.877E+00	3.829E+00
0	1.329E+00	1.706E+00	2.120E+00	2.563E+00	3.025E+00	3.500E+00	3.984E+00	4.472E+00	4.963E+00	
.2	1.428E+00	1.972E+00	2.642E+00	3.447E+00	4.400E+00	5.512E+00	6.804E+00	8.296E+00	1.002E+01	
.4	1.541E+00	2.309E+00	3.371E+00	4.814E+00	6.754E+00	9.340E+00	1.277E+01	1.731E+01	2.332E+01	
.6	1.672E+00	2.738E+00	4.400E+00	6.963E+00	1.089E+01	1.687E+01	2.597E+01	3.978E+01	6.073E+01	
.8	1.822E+00	3.285E+00	5.864E+00	1.039E+01	1.830E+01	3.212E+01	5.619E+01	9.814E+01	1.712E+02	
1	1.995E+00	3.985E+00	7.966E+00	1.593E+01	3.185E+01	6.369E+01	1.274E+02	2.547E+02	5.095E+02	
1.2	2.194E+00	4.885E+00	1.100E+01	2.497E+01	5.695E+01	1.303E+02	2.987E+02	6.854E+02	1.573E+03	
1.4	2.422E+00	6.045E+00	1.542E+01	3.988E+01	1.041E+02	2.730E+02	7.180E+02	1.892E+03	4.988E+03	
1.6	2.685E+00	7.542E+00	2.187E+01	6.466E+01	1.934E+02	5.822E+02	1.758E+03	5.321E+03	1.611E+04	
1.8	2.986E+00	9.482E+00	3.135E+01	1.062E+02	3.643E+02	1.259E+03	4.367E+03	1.518E+04	5.280E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

Bi(N,r,mu) for r = 1.1

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.527E-01	9.243E-01	9.082E-01	8.994E-01	8.948E-01	8.924E-01	8.911E-01	8.905E-01	8.902E-01	8.898E-01
-1.6	9.415E-01	9.048E-01	8.831E-01	8.709E-01	8.642E-01	8.607E-01	8.588E-01	8.578E-01	8.574E-01	8.569E-01
-1.4	9.493E-01	9.157E-01	8.948E-01	8.825E-01	8.755E-01	8.717E-01	8.696E-01	8.685E-01	8.680E-01	8.674E-01
-1.2	9.696E-01	9.482E-01	9.341E-01	9.254E-01	9.201E-01	9.171E-01	9.154E-01	9.145E-01	9.140E-01	9.134E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.040E+00	1.072E+00	1.096E+00	1.114E+00	1.126E+00	1.134E+00	1.139E+00	1.142E+00	1.144E+00	1.147E+00
-.6	1.088E+00	1.166E+00	1.230E+00	1.281E+00	1.319E+00	1.346E+00	1.366E+00	1.379E+00	1.389E+00	1.408E+00
-.4	1.147E+00	1.288E+00	1.416E+00	1.526E+00	1.617E+00	1.692E+00	1.751E+00	1.797E+00	1.832E+00	1.948E+00
-.2	1.217E+00	1.444E+00	1.671E+00	1.889E+00	2.091E+00	2.275E+00	2.440E+00	2.586E+00	2.714E+00	3.593E+00
0	1.298E+00	1.643E+00	2.026E+00	2.435E+00	2.863E+00	3.304E+00	3.752E+00	4.205E+00	4.661E+00	
.2	1.393E+00	1.898E+00	2.521E+00	3.272E+00	4.160E+00	5.199E+00	6.404E+00	7.797E+00	9.402E+00	
.4	1.504E+00	2.223E+00	3.219E+00	4.575E+00	6.398E+00	8.828E+00	1.205E+01	1.632E+01	2.197E+01	
.6	1.633E+00	2.640E+00	4.213E+00	6.639E+00	1.035E+01	1.602E+01	2.463E+01	3.771E+01	5.754E+01	
.8	1.783E+00	3.177E+00	5.637E+00	9.954E+00	1.750E+01	3.068E+01	5.365E+01	9.366E+01	1.634E+02	
1	1.957E+00	3.870E+00	7.696E+00	1.535E+01	3.065E+01	6.126E+01	1.225E+02	2.449E+02	4.898E+02	
1.2	2.158E+00	4.766E+00	1.069E+01	2.422E+01	5.521E+01	1.263E+02	2.894E+02	6.640E+02	1.524E+03	
1.4	2.391E+00	5.929E+00	1.508E+01	3.897E+01	1.016E+02	2.666E+02	7.011E+02	1.847E+03	4.870E+03	
1.6	2.660E+00	7.443E+00	2.154E+01	6.366E+01	1.904E+02	5.730E+02	1.731E+03	5.237E+03	1.586E+04	
1.8	2.972E+00	9.418E+00	3.111E+01	1.053E+02	3.614E+02	1.249E+03	4.333E+03	1.506E+04	5.238E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

Bi(N,r,mu) for r = 2

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.901E-01	9.836E-01	9.796E-01	9.774E-01	9.761E-01	9.755E-01	9.752E-01	9.750E-01	9.749E-01	9.748E-01
-1.6	9.845E-01	9.737E-01	9.669E-01	9.630E-01	9.607E-01	9.595E-01	9.589E-01	9.586E-01	9.584E-01	9.582E-01
-1.4	9.837E-01	9.719E-01	9.641E-01	9.593E-01	9.565E-01	9.550E-01	9.541E-01	9.537E-01	9.534E-01	9.532E-01
-1.2	9.886E-01	9.799E-01	9.738E-01	9.699E-01	9.675E-01	9.661E-01	9.653E-01	9.648E-01	9.646E-01	9.643E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.019E+00	1.035E+00	1.048E+00	1.058E+00	1.064E+00	1.069E+00	1.072E+00	1.074E+00	1.075E+00	1.077E+00
-.6	1.046E+00	1.090E+00	1.126E+00	1.156E+00	1.178E+00	1.195E+00	1.207E+00	1.215E+00	1.221E+00	1.233E+00
-.4	1.084E+00	1.168E+00	1.246E+00	1.315E+00	1.373E+00	1.420E+00	1.457E+00	1.487E+00	1.509E+00	1.583E+00
-.2	1.133E+00	1.277E+00	1.425E+00	1.568E+00	1.702E+00	1.824E+00	1.934E+00	2.031E+00	2.117E+00	2.704E+00
0	1.195E+00	1.427E+00	1.688E+00	1.971E+00	2.267E+00	2.573E+00	2.884E+00	3.198E+00	3.515E+00	
.2	1.273E+00	1.629E+00	2.075E+00	2.615E+00	3.256E+00	4.005E+00	4.876E+00	5.882E+00	7.042E+00	
.4	1.369E+00	1.901E+00	2.644E+00	3.659E+00	5.025E+00	6.847E+00	9.267E+00	1.247E+01	1.670E+01	
.6	1.486E+00	2.264E+00	3.485E+00	5.371E+00	8.262E+00	1.267E+01	1.937E+01	2.955E+01	4.498E+01	
.8	1.628E+00	2.750E+00	4.733E+00	8.215E+00	1.431E+01	2.494E+01	4.347E+01	7.576E+01	1.320E+02	
1	1.800E+00	3.400E+00	6.600E+00	1.300E+01	2.580E+01	5.140E+01	1.026E+02	2.050E+02	4.098E+02	
1.2	2.007E+00	4.271E+00	9.409E+00	2.114E+01	4.800E+01	1.096E+02	2.510E+02	5.756E+02	1.321E+03	
1.4	2.255E+00	5.439E+00	1.366E+01	3.512E+01	9.142E+01	2.395E+02	6.299E+02	1.659E+03	4.374E+03	
1.6	2.552E+00	7.012E+00	2.014E+01	5.935E+01	1.773E+02	5.334E+02	1.611E+03	4.874E+03	1.476E+04	
1.8	2.908E+00	9.132E+00	3.006E+01	1.017E+02	3.487E+02	1.205E+03	4.179E+03	1.453E+04	5.053E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = 4

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.974E-01	9.957E-01	9.946E-01	9.940E-01	9.936E-01	9.935E-01	9.934E-01	9.933E-01	9.933E-01	9.933E-01
-1.6	9.952E-01	9.918E-01	9.896E-01	9.884E-01	9.876E-01	9.873E-01	9.870E-01	9.869E-01	9.869E-01	9.868E-01
-1.4	9.941E-01	9.898E-01	9.869E-01	9.851E-01	9.840E-01	9.834E-01	9.831E-01	9.829E-01	9.829E-01	9.828E-01
-1.2	9.953E-01	9.916E-01	9.890E-01	9.873E-01	9.862E-01	9.856E-01	9.853E-01	9.851E-01	9.850E-01	9.849E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.010E+00	1.019E+00	1.026E+00	1.031E+00	1.035E+00	1.037E+00	1.039E+00	1.040E+00	1.041E+00	1.042E+00
-.6	1.027E+00	1.053E+00	1.075E+00	1.093E+00	1.107E+00	1.117E+00	1.124E+00	1.129E+00	1.132E+00	1.140E+00
-.4	1.054E+00	1.109E+00	1.160E+00	1.205E+00	1.243E+00	1.274E+00	1.299E+00	1.318E+00	1.333E+00	1.382E+00
-.2	1.092E+00	1.194E+00	1.298E+00	1.399E+00	1.494E+00	1.580E+00	1.658E+00	1.727E+00	1.788E+00	2.204E+00
0	1.144E+00	1.318E+00	1.514E+00	1.726E+00	1.949E+00	2.179E+00	2.414E+00	2.651E+00	2.890E+00	
.2	1.214E+00	1.495E+00	1.848E+00	2.275E+00	2.783E+00	3.377E+00	4.067E+00	4.865E+00	5.784E+00	
.4	1.303E+00	1.742E+00	2.356E+00	3.196E+00	4.326E+00	5.834E+00	7.837E+00	1.049E+01	1.399E+01	
.6	1.415E+00	2.082E+00	3.129E+00	4.748E+00	7.229E+00	1.101E+01	1.676E+01	2.550E+01	3.875E+01	
.8	1.555E+00	2.548E+00	4.303E+00	7.385E+00	1.278E+01	2.219E+01	3.860E+01	6.718E+01	1.170E+02	
1	1.727E+00	3.182E+00	6.091E+00	1.191E+01	2.355E+01	4.682E+01	9.336E+01	1.865E+02	3.726E+02	
1.2	1.938E+00	4.045E+00	8.826E+00	1.974E+01	4.472E+01	1.020E+02	2.336E+02	5.356E+02	1.229E+03	
1.4	2.194E+00	5.219E+00	1.303E+01	3.340E+01	8.686E+01	2.275E+02	5.981E+02	1.575E+03	4.154E+03	
1.6	2.504E+00	6.819E+00	1.952E+01	5.745E+01	1.715E+02	5.160E+02	1.558E+03	4.714E+03	1.428E+04	
1.8	2.879E+00	9.006E+00	2.960E+01	1.000E+02	3.431E+02	1.186E+03	4.112E+03	1.429E+04	4.972E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = 8

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.993E-01	9.988E-01	9.985E-01	9.983E-01	9.982E-01	9.981E-01	9.981E-01	9.981E-01	9.981E-01	9.981E-01
-1.6	9.984E-01	9.973E-01	9.966E-01	9.962E-01	9.960E-01	9.958E-01	9.958E-01	9.957E-01	9.957E-01	9.957E-01
-1.4	9.978E-01	9.961E-01	9.950E-01	9.944E-01	9.940E-01	9.937E-01	9.936E-01	9.936E-01	9.935E-01	9.935E-01
-1.2	9.980E-01	9.963E-01	9.952E-01	9.945E-01	9.940E-01	9.938E-01	9.936E-01	9.935E-01	9.935E-01	9.934E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.006E+00	1.011E+00	1.014E+00	1.017E+00	1.020E+00	1.021E+00	1.022E+00	1.022E+00	1.023E+00	1.023E+00
-.6	1.017E+00	1.033E+00	1.047E+00	1.058E+00	1.067E+00	1.073E+00	1.078E+00	1.081E+00	1.083E+00	1.088E+00
-.4	1.037E+00	1.075E+00	1.110E+00	1.142E+00	1.168E+00	1.189E+00	1.207E+00	1.220E+00	1.231E+00	1.264E+00
-.2	1.069E+00	1.145E+00	1.223E+00	1.299E+00	1.371E+00	1.436E+00	1.494E+00	1.546E+00	1.592E+00	1.905E+00
0	1.116E+00	1.255E+00	1.413E+00	1.584E+00	1.763E+00	1.949E+00	2.137E+00	2.328E+00	2.520E+00	
.2	1.181E+00	1.420E+00	1.720E+00	2.083E+00	2.514E+00	3.019E+00	3.606E+00	4.284E+00	5.066E+00	
.4	1.268E+00	1.657E+00	2.202E+00	2.946E+00	3.949E+00	5.286E+00	7.062E+00	9.414E+00	1.252E+01	
.6	1.380E+00	1.991E+00	2.950E+00	4.433E+00	6.706E+00	1.017E+01	1.544E+01	2.345E+01	3.559E+01	
.8	1.521E+00	2.453E+00	4.101E+00	6.996E+00	1.206E+01	2.090E+01	3.630E+01	6.315E+01	1.099E+02	
1	1.696E+00	3.087E+00	5.870E+00	1.143E+01	2.257E+01	4.483E+01	8.935E+01	1.784E+02	3.565E+02	
1.2	1.910E+00	3.953E+00	8.590E+00	1.917E+01	4.341E+01	9.898E+01	2.265E+02	5.195E+02	1.192E+03	
1.4	2.170E+00	5.135E+00	1.279E+01	3.276E+01	8.515E+01	2.230E+02	5.862E+02	1.544E+03	4.071E+03	
1.6	2.486E+00	6.751E+00	1.930E+01	5.678E+01	1.695E+02	5.099E+02	1.540E+03	4.658E+03	1.411E+04	
1.8	2.870E+00	8.963E+00	2.945E+01	9.951E+01	3.413E+02	1.179E+03	4.090E+03	1.421E+04	4.945E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = 16

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.998E-01	9.997E-01	9.996E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01
-1.6	9.995E-01	9.991E-01	9.989E-01	9.987E-01	9.987E-01	9.986E-01	9.986E-01	9.986E-01	9.986E-01	9.986E-01
-1.4	9.992E-01	9.985E-01	9.981E-01	9.979E-01	9.977E-01	9.976E-01	9.976E-01	9.976E-01	9.975E-01	9.975E-01
-1.2	9.991E-01	9.984E-01	9.979E-01	9.976E-01	9.974E-01	9.973E-01	9.972E-01	9.972E-01	9.972E-01	9.971E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.003E+00	1.006E+00	1.008E+00	1.010E+00	1.011E+00	1.012E+00	1.012E+00	1.013E+00	1.013E+00	1.013E+00
-.6	1.011E+00	1.021E+00	1.030E+00	1.037E+00	1.043E+00	1.047E+00	1.050E+00	1.052E+00	1.053E+00	1.056E+00
-.4	1.026E+00	1.053E+00	1.078E+00	1.101E+00	1.119E+00	1.135E+00	1.147E+00	1.157E+00	1.164E+00	1.188E+00
-.2	1.054E+00	1.113E+00	1.174E+00	1.233E+00	1.289E+00	1.339E+00	1.385E+00	1.425E+00	1.461E+00	1.705E+00
0	1.097E+00	1.214E+00	1.346E+00	1.489E+00	1.639E+00	1.794E+00	1.952E+00	2.112E+00	2.273E+00	
.2	1.160E+00	1.372E+00	1.637E+00	1.958E+00	2.340E+00	2.787E+00	3.307E+00	3.907E+00	4.599E+00	
.4	1.247E+00	1.606E+00	2.108E+00	2.794E+00	3.717E+00	4.950E+00	6.587E+00	8.754E+00	1.162E+01	
.6	1.360E+00	1.939E+00	2.848E+00	4.254E+00	6.409E+00	9.696E+00	1.469E+01	2.228E+01	3.379E+01	
.8	1.504E+00	2.404E+00	3.997E+00	6.794E+00	1.169E+01	2.023E+01	3.512E+01	6.106E+01	1.062E+02	
1	1.681E+00	3.043E+00	5.766E+00	1.121E+01	2.211E+01	4.389E+01	8.747E+01	1.746E+02	3.489E+02	
1.2	1.890E+00	3.914E+00	8.491E+00	1.894E+01	4.285E+01	9.769E+01	2.236E+02	5.127E+02	1.177E+03	
1.4	2.161E+00	5.104E+00	1.270E+01	3.251E+01	8.450E+01	2.213E+02	5.817E+02	1.532E+03	4.039E+03	
1.6	2.480E+00	6.727E+00	1.923E+01	5.654E+01	1.688E+02	5.077E+02	1.533E+03	4.639E+03	1.405E+04	
1.8	2.866E+00	8.949E+00	2.940E+01	9.934E+01	3.406E+02	1.177E+03	4.083E+03	1.419E+04	4.936E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B1(N,r,mu) for r = 32

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01
-1.6	9.998E-01	9.997E-01	9.996E-01	9.996E-01	9.996E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01
-1.4	9.997E-01	9.994E-01	9.993E-01	9.992E-01	9.991E-01	9.991E-01	9.991E-01	9.991E-01	9.991E-01	9.991E-01
-1.2	9.996E-01	9.993E-01	9.991E-01	9.990E-01	9.989E-01	9.988E-01	9.988E-01	9.988E-01	9.988E-01	9.988E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.002E+00	1.003E+00	1.005E+00	1.006E+00	1.006E+00	1.007E+00	1.007E+00	1.007E+00	1.007E+00	1.008E+00
-.6	1.007E+00	1.014E+00	1.019E+00	1.024E+00	1.028E+00	1.030E+00	1.032E+00	1.033E+00	1.034E+00	1.036E+00
-.4	1.019E+00	1.039E+00	1.057E+00	1.073E+00	1.087E+00	1.098E+00	1.107E+00	1.114E+00	1.119E+00	1.136E+00
-.2	1.043E+00	1.090E+00	1.139E+00	1.186E+00	1.230E+00	1.271E+00	1.307E+00	1.339E+00	1.368E+00	1.563E+00
0	1.083E+00	1.184E+00	1.297E+00	1.420E+00	1.550E+00	1.683E+00	1.819E+00	1.957E+00	2.095E+00	
.2	1.146E+00	1.338E+00	1.579E+00	1.871E+00	2.218E+00	2.625E+00	3.097E+00	3.643E+00	4.272E+00	
.4	1.233E+00	1.572E+00	2.046E+00	2.693E+00	3.565E+00	4.729E+00	6.275E+00	8.320E+00	1.102E+01	
.6	1.348E+00	1.908E+00	2.787E+00	4.147E+00	6.231E+00	9.408E+00	1.424E+01	2.158E+01	3.270E+01	
.8	1.494E+00	2.378E+00	3.940E+00	6.685E+00	1.149E+01	1.986E+01	3.447E+01	5.992E+01	1.042E+02	
1	1.674E+00	3.021E+00	5.716E+00	1.111E+01	2.188E+01	4.344E+01	8.656E+01	1.728E+02	3.453E+02	
1.2	1.893E+00	3.898E+00	8.448E+00	1.883E+01	4.261E+01	9.714E+01	2.223E+02	5.097E+02	1.170E+03	
1.4	2.158E+00	5.091E+00	1.266E+01	3.242E+01	8.425E+01	2.206E+02	5.799E+02	1.527E+03	4.027E+03	
1.6	2.478E+00	6.719E+00	1.920E+01	5.646E+01	1.685E+02	5.070E+02	1.531E+03	4.632E+03	1.403E+04	
1.8	2.865E+00	8.945E+00	2.938E+01	9.928E+01	3.405E+02	1.176E+03	4.081E+03	1.418E+04	4.934E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = 64

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01
-1.4	9.999E-01	9.998E-01	9.997E-01	9.997E-01	9.997E-01	9.997E-01	9.997E-01	9.996E-01	9.996E-01	9.996E-01
-1.2	9.998E-01	9.997E-01	9.996E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01	9.995E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.001E+00	1.002E+00	1.003E+00	1.003E+00	1.004E+00	1.004E+00	1.004E+00	1.004E+00	1.004E+00	1.004E+00
-.6	1.005E+00	1.009E+00	1.013E+00	1.016E+00	1.018E+00	1.020E+00	1.021E+00	1.022E+00	1.022E+00	1.024E+00
-.4	1.014E+00	1.026E+00	1.042E+00	1.054E+00	1.064E+00	1.072E+00	1.078E+00	1.083E+00	1.087E+00	1.100E+00
-.2	1.035E+00	1.073E+00	1.112E+00	1.151E+00	1.187E+00	1.220E+00	1.249E+00	1.275E+00	1.298E+00	1.456E+00
0	1.073E+00	1.161E+00	1.261E+00	1.369E+00	1.482E+00	1.600E+00	1.719E+00	1.840E+00	1.961E+00	
.2	1.135E+00	1.313E+00	1.536E+00	1.808E+00	2.129E+00	2.506E+00	2.944E+00	3.450E+00	4.033E+00	
.4	1.224E+00	1.546E+00	2.003E+00	2.624E+00	3.461E+00	4.576E+00	6.060E+00	8.023E+00	1.042E+01	
.6	1.341E+00	1.889E+00	2.749E+00	4.080E+00	6.119E+00	9.229E+00	1.396E+01	2.114E+01	3.203E+01	
.8	1.489E+00	2.363E+00	3.909E+00	6.624E+00	1.137E+01	1.966E+01	3.411E+01	5.929E+01	1.031E+02	
1	1.670E+00	3.010E+00	5.691E+00	1.105E+01	2.177E+01	4.322E+01	8.611E+01	1.719E+02	3.435E+02	
1.2	1.890E+00	3.891E+00	8.429E+00	1.879E+01	4.251E+01	9.690E+01	2.218E+02	5.085E+02	1.167E+03	
1.4	2.156E+00	5.087E+00	1.265E+01	3.238E+01	8.415E+01	2.204E+02	5.793E+02	1.526E+03	4.022E+03	
1.6	2.477E+00	6.716E+00	1.919E+01	5.644E+01	1.685E+02	5.067E+02	1.530E+03	4.630E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.927E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.933E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = 128

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	9.999E-01	9.999E-01	9.999E-01	9.999E-01
-1.4	1.000E+00	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01
-1.2	9.999E-01	9.999E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01	9.998E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.001E+00	1.001E+00	1.002E+00	1.002E+00	1.002E+00	1.002E+00	1.002E+00	1.002E+00	1.002E+00	1.002E+00
-.6	1.003E+00	1.006E+00	1.008E+00	1.010E+00	1.012E+00	1.013E+00	1.014E+00	1.014E+00	1.015E+00	1.015E+00
-.4	1.010E+00	1.021E+00	1.031E+00	1.040E+00	1.047E+00	1.053E+00	1.058E+00	1.062E+00	1.065E+00	1.074E+00
-.2	1.029E+00	1.060E+00	1.092E+00	1.124E+00	1.153E+00	1.181E+00	1.205E+00	1.226E+00	1.245E+00	1.375E+00
0	1.065E+00	1.144E+00	1.232E+00	1.329E+00	1.430E+00	1.534E+00	1.640E+00	1.748E+00	1.856E+00	
.2	1.127E+00	1.294E+00	1.504E+00	1.759E+00	2.062E+00	2.416E+00	2.827E+00	3.303E+00	3.851E+00	
.4	1.217E+00	1.532E+00	1.973E+00	2.576E+00	3.388E+00	4.471E+00	5.909E+00	7.813E+00	1.033E+01	
.6	1.336E+00	1.877E+00	2.725E+00	4.037E+00	6.048E+00	9.115E+00	1.378E+01	2.086E+01	3.160E+01	
.8	1.486E+00	2.354E+00	3.891E+00	6.589E+00	1.131E+01	1.955E+01	3.391E+01	5.893E+01	1.025E+02	
1	1.668E+00	3.005E+00	5.679E+00	1.103E+01	2.172E+01	4.311E+01	8.589E+01	1.714E+02	3.426E+02	
1.2	1.889E+00	3.887E+00	8.421E+00	1.877E+01	4.246E+01	9.680E+01	2.215E+02	5.079E+02	1.166E+03	
1.4	2.156E+00	5.085E+00	1.265E+01	3.237E+01	8.411E+01	2.203E+02	5.790E+02	1.525E+03	4.021E+03	
1.6	2.477E+00	6.715E+00	1.919E+01	5.643E+01	1.684E+02	5.067E+02	1.530E+03	4.629E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.926E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.932E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = 256

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	9.999E-01	9.999E-01	9.999E-01
-1.2	1.000E+00	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01	9.999E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00
-.6	1.002E+00	1.004E+00	1.005E+00	1.007E+00	1.008E+00	1.008E+00	1.009E+00	1.009E+00	1.010E+00	1.010E+00
-.4	1.008E+00	1.016E+00	1.023E+00	1.029E+00	1.035E+00	1.039E+00	1.043E+00	1.046E+00	1.048E+00	1.055E+00
-.2	1.024E+00	1.050E+00	1.077E+00	1.103E+00	1.127E+00	1.150E+00	1.170E+00	1.188E+00	1.204E+00	1.311E+00
0	1.059E+00	1.130E+00	1.210E+00	1.296E+00	1.388E+00	1.482E+00	1.577E+00	1.674E+00	1.772E+00	
.2	1.121E+00	1.280E+00	1.479E+00	1.722E+00	2.009E+00	2.346E+00	2.737E+00	3.189E+00	3.710E+00	
.4	1.212E+00	1.520E+00	1.952E+00	2.541E+00	3.335E+00	4.394E+00	5.800E+00	7.662E+00	1.012E+01	
.6	1.333E+00	1.869E+00	2.709E+00	4.010E+00	6.003E+00	9.042E+00	1.366E+01	2.068E+01	3.132E+01	
.8	1.484E+00	2.350E+00	3.881E+00	6.570E+00	1.127E+01	1.948E+01	3.380E+01	5.873E+01	1.022E+02	
1	1.668E+00	3.003E+00	5.673E+00	1.101E+01	2.169E+01	4.305E+01	8.578E+01	1.712E+02	3.421E+02	
1.2	1.889E+00	3.886E+00	8.418E+00	1.874E+01	4.244E+01	9.675E+01	2.214E+02	5.077E+02	1.165E+03	
1.4	2.156E+00	5.084E+00	1.264E+01	3.236E+01	8.410E+01	2.202E+02	5.789E+02	1.525E+03	4.020E+03	
1.6	2.477E+00	6.715E+00	1.919E+01	5.642E+01	1.684E+02	5.066E+02	1.530E+03	4.629E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.926E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.932E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = 512

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	1.000E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00	1.001E+00
-.6	1.001E+00	1.003E+00	1.004E+00	1.004E+00	1.005E+00	1.006E+00	1.006E+00	1.006E+00	1.006E+00	1.007E+00
-.4	1.006E+00	1.012E+00	1.017E+00	1.022E+00	1.026E+00	1.030E+00	1.032E+00	1.034E+00	1.036E+00	1.041E+00
-.2	1.020E+00	1.042E+00	1.064E+00	1.086E+00	1.107E+00	1.125E+00	1.142E+00	1.157E+00	1.170E+00	1.261E+00
0	1.054E+00	1.118E+00	1.191E+00	1.270E+00	1.353E+00	1.438E+00	1.526E+00	1.614E+00	1.703E+00	
.2	1.116E+00	1.268E+00	1.460E+00	1.692E+00	1.967E+00	2.290E+00	2.665E+00	3.098E+00	3.598E+00	
.4	1.209E+00	1.512E+00	1.936E+00	2.516E+00	3.296E+00	4.338E+00	5.721E+00	7.553E+00	9.973E+00	
.6	1.331E+00	1.864E+00	2.699E+00	3.992E+00	5.973E+00	8.994E+00	1.359E+01	2.056E+01	3.114E+01	
.8	1.483E+00	2.347E+00	3.875E+00	6.558E+00	1.125E+01	1.945E+01	3.373E+01	5.862E+01	1.020E+02	
1	1.667E+00	3.001E+00	5.670E+00	1.101E+01	2.168E+01	4.303E+01	8.572E+01	1.711E+02	3.419E+02	
1.2	1.889E+00	3.885E+00	8.416E+00	1.874E+01	4.243E+01	9.673E+01	2.214E+02	5.076E+02	1.165E+03	
1.4	2.156E+00	5.084E+00	1.264E+01	3.236E+01	8.410E+01	2.202E+02	5.789E+02	1.525E+03	4.020E+03	
1.6	2.477E+00	6.714E+00	1.919E+01	5.642E+01	1.684E+02	5.066E+02	1.530E+03	4.628E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.926E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.932E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = 1024

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.6	1.001E+00	1.002E+00	1.002E+00	1.003E+00	1.003E+00	1.004E+00	1.004E+00	1.004E+00	1.004E+00	1.004E+00
-.4	1.004E+00	1.009E+00	1.013E+00	1.017E+00	1.020E+00	1.022E+00	1.024E+00	1.026E+00	1.027E+00	1.031E+00
-.2	1.017E+00	1.035E+00	1.054E+00	1.073E+00	1.090E+00	1.106E+00	1.120E+00	1.132E+00	1.144E+00	1.220E+00
0	1.049E+00	1.108E+00	1.175E+00	1.248E+00	1.324E+00	1.402E+00	1.483E+00	1.564E+00	1.645E+00	
.2	1.112E+00	1.259E+00	1.444E+00	1.668E+00	1.934E+00	2.245E+00	2.607E+00	3.025E+00	3.507E+00	
.4	1.206E+00	1.505E+00	1.924E+00	2.497E+00	3.268E+00	4.297E+00	5.663E+00	7.472E+00	9.862E+00	
.6	1.330E+00	1.860E+00	2.693E+00	3.980E+00	5.953E+00	8.963E+00	1.354E+01	2.049E+01	3.102E+01	
.8	1.482E+00	2.345E+00	3.872E+00	6.552E+00	1.124E+01	1.942E+01	3.369E+01	5.855E+01	1.018E+02	
1	1.667E+00	3.001E+00	5.668E+00	1.100E+01	2.167E+01	4.301E+01	8.569E+01	1.711E+02	3.418E+02	
1.2	1.889E+00	3.885E+00	8.416E+00	1.876E+01	4.243E+01	9.672E+01	2.213E+02	5.075E+02	1.165E+03	
1.4	2.156E+00	5.084E+00	1.264E+01	3.236E+01	8.409E+01	2.202E+02	5.789E+02	1.525E+03	4.020E+03	
1.6	2.477E+00	6.714E+00	1.919E+01	5.642E+01	1.684E+02	5.066E+02	1.530E+03	4.628E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.926E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.932E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = 2048

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.6	1.001E+00	1.001E+00	1.002E+00	1.002E+00	1.002E+00	1.002E+00	1.003E+00	1.003E+00	1.003E+00	1.003E+00
-.4	1.003E+00	1.007E+00	1.010E+00	1.012E+00	1.015E+00	1.017E+00	1.018E+00	1.019E+00	1.020E+00	1.023E+00
-.2	1.014E+00	1.030E+00	1.046E+00	1.061E+00	1.076E+00	1.089E+00	1.101E+00	1.112E+00	1.122E+00	1.186E+00
0	1.045E+00	1.100E+00	1.162E+00	1.229E+00	1.299E+00	1.372E+00	1.446E+00	1.521E+00	1.596E+00	
.2	1.108E+00	1.251E+00	1.431E+00	1.648E+00	1.906E+00	2.209E+00	2.560E+00	2.966E+00	3.434E+00	
.4	1.204E+00	1.501E+00	1.916E+00	2.483E+00	3.247E+00	4.266E+00	5.620E+00	7.412E+00	9.780E+00	
.6	1.329E+00	1.858E+00	2.688E+00	3.972E+00	5.941E+00	8.942E+00	1.351E+01	2.044E+01	3.095E+01	
.8	1.482E+00	2.345E+00	3.870E+00	6.548E+00	1.123E+01	1.941E+01	3.367E+01	5.851E+01	1.018E+02	
1	1.667E+00	3.000E+00	5.667E+00	1.100E+01	2.167E+01	4.301E+01	8.568E+01	1.710E+02	3.417E+02	
1.2	1.889E+00	3.885E+00	8.415E+00	1.875E+01	4.243E+01	9.672E+01	2.213E+02	5.075E+02	1.165E+03	
1.4	2.156E+00	5.084E+00	1.264E+01	3.236E+01	8.409E+01	2.202E+02	5.789E+02	1.525E+03	4.020E+03	
1.6	2.477E+00	6.714E+00	1.919E+01	5.642E+01	1.684E+02	5.066E+02	1.530E+03	4.628E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.926E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.932E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

BI(N,r,mu) for r = INFINITY

Mu \ N=	4	8	16	32	64	128	256	512	1024	INF
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
0	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
.2	1.091E+00	1.210E+00	1.360E+00	1.541E+00	1.757E+00	2.009E+00	2.303E+00	2.642E+00	3.033E+00	
.4	1.198E+00	1.487E+00	1.890E+00	2.441E+00	3.184E+00	4.174E+00	5.489E+00	7.231E+00	9.533E+00	
.6	1.327E+00	1.854E+00	2.680E+00	3.958E+00	5.916E+00	8.903E+00	1.344E+01	2.034E+01	3.080E+01	
.8	1.482E+00	2.343E+00	3.867E+00	6.543E+00	1.123E+01	1.940E+01	3.364E+01	5.846E+01	1.017E+02	
1	1.667E+00	3.000E+00	5.667E+00	1.100E+01	2.167E+01	4.300E+01	8.567E+01	1.710E+02	3.417E+02	
1.2	1.889E+00	3.885E+00	8.415E+00	1.875E+01	4.243E+01	9.672E+01	2.213E+02	5.075E+02	1.165E+03	
1.4	2.156E+00	5.084E+00	1.264E+01	3.236E+01	8.409E+01	2.202E+02	5.789E+02	1.525E+03	4.020E+03	
1.6	2.477E+00	6.714E+00	1.919E+01	5.642E+01	1.684E+02	5.066E+02	1.530E+03	4.628E+03	1.402E+04	
1.8	2.865E+00	8.943E+00	2.938E+01	9.926E+01	3.404E+02	1.176E+03	4.080E+03	1.418E+04	4.932E+04	
2	3.333E+00	1.200E+01	4.533E+01	1.760E+02	6.933E+02	2.752E+03	1.097E+04	4.378E+04	1.749E+05	

B2(r, mu)

MU \ r =	.0001	.0003	.001	.003	.01	.03	.1	.3	.5	.7
-2 :	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01
-1.8 :	1.112E-01	1.385E-01	1.762E-01	2.195E-01	2.793E-01	3.479E-01	4.431E-01	5.566E-01	6.264E-01	6.889E-01
-1.6 :	1.874E-02	2.908E-02	4.708E-02	7.306E-02	1.183E-01	1.836E-01	2.979E-01	4.693E-01	5.901E-01	7.013E-01
-1.4 :	3.205E-03	6.196E-03	1.276E-02	2.467E-02	5.081E-02	9.828E-02	2.032E-01	3.999E-01	5.572E-01	7.061E-01
-1.2 :	5.586E-04	1.345E-03	3.525E-03	8.489E-03	2.225E-02	5.362E-02	1.410E-01	3.444E-01	5.273E-01	7.052E-01
-1 :	1.000E-04	3.000E-04	1.000E-03	3.000E-03	1.000E-02	3.000E-02	1.000E-01	3.000E-01	5.000E-01	7.000E-01
-.8 :	1.862E-05	6.956E-05	2.949E-04	1.101E-03	4.662E-03	1.735E-02	7.271E-02	2.642E-01	4.749E-01	6.916E-01
-.6 :	3.687E-06	1.715E-05	9.231E-05	4.280E-04	2.288E-03	1.047E-02	5.438E-02	2.351E-01	4.517E-01	6.808E-01
-.4 :	8.121E-07	4.678E-06	3.174E-05	1.809E-04	1.204E-03	6.664E-03	4.195E-02	2.112E-01	4.303E-01	6.683E-01
-.2 :	2.159E-07	1.510E-06	1.260E-05	8.606E-05	6.921E-04	4.507E-03	3.340E-02	1.915E-01	4.103E-01	6.546E-01
0 :	7.726E-08	6.240E-07	6.065E-06	4.745E-05	4.404E-04	3.250E-03	2.742E-02	1.750E-01	3.915E-01	6.401E-01
.2 :	3.906E-08	3.397E-07	3.594E-06	3.048E-05	3.100E-04	2.494E-03	2.316E-02	1.611E-01	3.740E-01	6.251E-01
.4 :	2.590E-08	2.311E-07	2.530E-06	2.228E-05	2.381E-04	2.020E-03	2.006E-02	1.492E-01	3.574E-01	6.097E-01
.6 :	2.013E-08	1.808E-07	2.001E-06	1.788E-05	1.955E-04	1.708E-03	1.773E-02	1.388E-01	3.416E-01	5.943E-01
.8 :	1.700E-08	1.529E-07	1.697E-06	1.524E-05	1.683E-04	1.493E-03	1.593E-02	1.297E-01	3.267E-01	5.789E-01
1 :	1.500E-08	1.350E-07	1.500E-06	1.349E-05	1.495E-04	1.337E-03	1.450E-02	1.215E-01	3.125E-01	5.635E-01
1.2 :	1.357E-08	1.221E-07	1.356E-06	1.221E-05	1.355E-04	1.216E-03	1.333E-02	1.141E-01	2.989E-01	5.483E-01
1.4 :	1.245E-08	1.120E-07	1.245E-06	1.120E-05	1.244E-04	1.118E-03	1.233E-02	1.074E-01	2.859E-01	5.333E-01
1.6 :	1.152E-08	1.037E-07	1.152E-06	1.037E-05	1.152E-04	1.036E-03	1.147E-02	1.012E-01	2.735E-01	5.186E-01
1.8 :	1.072E-08	9.645E-08	1.072E-06	9.645E-06	1.072E-04	9.642E-04	1.070E-02	9.541E-02	2.615E-01	5.042E-01
2 :	1.000E-08	9.000E-08	1.000E-06	9.000E-06	1.000E-04	9.000E-04	1.000E-02	9.000E-02	2.500E-01	4.900E-01

B2(r, mu)

RU \ r =	1	1.01	1.1	2	4	8	16	32	64	128
-2	1.000E+00	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01
-1.8	1.000E+00	8.614E-01	7.983E-01	7.196E-01	7.062E-01	7.028E-01	7.018E-01	7.015E-01	7.015E-01	7.014E-01
-1.6	1.000E+00	9.429E-01	8.708E-01	7.787E-01	7.561E-01	7.494E-01	7.472E-01	7.465E-01	7.462E-01	7.462E-01
-1.4	1.000E+00	9.776E-01	9.281E-01	8.446E-01	8.192E-01	8.103E-01	8.071E-01	8.058E-01	8.053E-01	8.052E-01
-1.2	1.000E+00	9.929E-01	9.693E-01	9.181E-01	8.990E-01	8.912E-01	8.879E-01	8.865E-01	8.859E-01	8.856E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	1.004E+00	1.024E+00	1.091E+00	1.128E+00	1.148E+00	1.159E+00	1.166E+00	1.170E+00	1.172E+00
-.6	1.000E+00	1.006E+00	1.043E+00	1.192E+00	1.290E+00	1.351E+00	1.392E+00	1.418E+00	1.436E+00	1.447E+00
-.4	1.000E+00	1.007E+00	1.060E+00	1.304E+00	1.494E+00	1.633E+00	1.738E+00	1.817E+00	1.877E+00	1.923E+00
-.2	1.000E+00	1.009E+00	1.075E+00	1.429E+00	1.752E+00	2.027E+00	2.265E+00	2.472E+00	2.652E+00	2.809E+00
0	1.000E+00	1.010E+00	1.089E+00	1.566E+00	2.078E+00	2.581E+00	3.082E+00	3.582E+00	4.082E+00	4.582E+00
.2	1.000E+00	1.011E+00	1.102E+00	1.718E+00	2.489E+00	3.364E+00	4.365E+00	5.514E+00	6.834E+00	8.351E+00
.4	1.000E+00	1.012E+00	1.114E+00	1.886E+00	3.007E+00	4.473E+00	6.404E+00	8.951E+00	1.231E+01	1.674E+01
.6	1.000E+00	1.013E+00	1.126E+00	2.071E+00	3.658E+00	6.051E+00	9.673E+00	1.516E+01	2.348E+01	3.609E+01
.8	1.000E+00	1.014E+00	1.138E+00	2.275E+00	4.475E+00	8.297E+00	1.495E+01	2.653E+01	4.669E+01	8.179E+01
1	1.000E+00	1.015E+00	1.150E+00	2.500E+00	5.500E+00	1.150E+01	2.350E+01	4.750E+01	9.550E+01	1.915E+02
1.2	1.000E+00	1.016E+00	1.162E+00	2.747E+00	6.784E+00	1.607E+01	3.741E+01	8.644E+01	1.991E+02	4.579E+02
1.4	1.000E+00	1.017E+00	1.174E+00	3.018E+00	8.389E+00	2.259E+01	6.008E+01	1.590E+02	4.201E+02	1.109E+03
1.6	1.000E+00	1.018E+00	1.186E+00	3.316E+00	1.039E+01	3.188E+01	9.706E+01	2.947E+02	8.937E+02	2.710E+03
1.8	1.000E+00	1.019E+00	1.198E+00	3.642E+00	1.289E+01	4.513E+01	1.574E+02	5.485E+02	1.910E+03	6.653E+03
2	1.000E+00	1.020E+00	1.210E+00	4.000E+00	1.600E+01	6.400E+01	2.560E+02	1.024E+03	4.096E+03	1.638E+04

RU \ r =	256	512	1024	2048	4096	INF
-2	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01
-1.8	7.014E-01	7.014E-01	7.014E-01	7.014E-01	7.014E-01	7.014E-01
-1.6	7.461E-01	7.461E-01	7.461E-01	7.461E-01	7.461E-01	7.461E-01
-1.4	8.051E-01	8.051E-01	8.051E-01	8.051E-01	8.051E-01	8.051E-01
-1.2	8.855E-01	8.854E-01	8.854E-01	8.854E-01	8.854E-01	8.854E-01
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.173E+00	1.174E+00	1.174E+00	1.174E+00	1.174E+00	1.175E+00
-.6	1.455E+00	1.460E+00	1.463E+00	1.465E+00	1.467E+00	1.470E+00
-.4	1.957E+00	1.983E+00	2.003E+00	2.018E+00	2.029E+00	2.065E+00
-.2	2.945E+00	3.064E+00	3.167E+00	3.257E+00	3.336E+00	3.863E+00
0	5.082E+00	5.582E+00	6.082E+00	6.582E+00	7.082E+00	
.2	1.009E+01	1.209E+01	1.439E+01	1.703E+01	2.006E+01	
.4	2.259E+01	3.031E+01	4.050E+01	5.394E+01	7.167E+01	
.6	5.521E+01	8.418E+01	1.281E+02	1.947E+02	2.955E+02	
.8	1.429E+02	2.493E+02	4.346E+02	7.571E+02	1.319E+03	
1	3.835E+02	7.675E+02	1.536E+03	3.072E+03	6.144E+03	
1.2	1.052E+03	2.418E+03	5.556E+03	1.277E+04	2.933E+04	
1.4	2.928E+03	7.727E+03	2.039E+04	5.382E+04	1.420E+05	
1.6	8.215E+03	2.490E+04	7.549E+04	2.288E+05	6.937E+05	
1.8	2.317E+04	8.067E+04	2.809E+05	9.782E+05	3.406E+06	
2	6.554E+04	2.621E+05	1.049E+06	4.194E+06	1.678E+07	

B3(2,M,r,mu) for r = .01

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2 :	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8 :	1.000E+00	2.214E+00	5.842E+00	1.829E+01	6.502E+01	2.536E+02	1.173E+03	3.308E+03	4.836E+03	6.080E+03	7.145E+03
-1.6 :	1.000E+00	2.419E+00	7.398E+00	2.638E+01	1.018E+02	4.151E+02	1.888E+03	5.075E+03	7.222E+03	8.748E+03	9.882E+03
-1.4 :	1.000E+00	2.617E+00	8.745E+00	3.295E+01	1.306E+02	5.358E+02	2.342E+03	5.964E+03	8.263E+03	9.700E+03	1.063E+04
-1.2 :	1.000E+00	2.811E+00	9.937E+00	3.843E+01	1.534E+02	6.247E+02	2.595E+03	6.253E+03	8.449E+03	9.671E+03	1.036E+04
-1 :	1.000E+00	3.000E+00	1.100E+01	4.300E+01	1.710E+02	6.830E+02	2.674E+03	6.094E+03	8.047E+03	9.024E+03	9.512E+03
-.8 :	1.000E+00	3.180E+00	1.191E+01	4.651E+01	1.823E+02	7.065E+02	2.588E+03	5.578E+03	7.208E+03	7.950E+03	8.280E+03
-.6 :	1.000E+00	3.337E+00	1.238E+01	4.848E+01	1.851E+02	6.877E+02	2.342E+03	4.775E+03	6.045E+03	6.577E+03	6.791E+03
-.4 :	1.000E+00	3.444E+00	1.285E+01	4.812E+01	1.766E+02	6.215E+02	1.957E+03	3.774E+03	4.686E+03	5.041E+03	5.171E+03
-.2 :	1.000E+00	3.463E+00	1.252E+01	4.480E+01	1.557E+02	5.131E+02	1.486E+03	2.712E+03	3.304E+03	3.521E+03	3.594E+03
0 :	1.000E+00	3.368E+00	1.151E+01	3.866E+01	1.253E+02	3.826E+02	1.015E+03	1.753E+03	2.097E+03	2.216E+03	2.254E+03
.2 :	1.000E+00	3.163E+00	9.972E+00	3.090E+01	9.213E+01	2.579E+02	6.253E+02	1.021E+03	1.200E+03	1.259E+03	1.277E+03
.4 :	1.000E+00	2.883E+00	8.251E+00	2.314E+01	6.266E+01	1.595E+02	3.521E+02	5.442E+02	6.280E+02	6.549E+02	6.627E+02
.6 :	1.000E+00	2.574E+00	6.575E+00	1.651E+01	4.018E+01	9.228E+01	1.851E+02	2.709E+02	3.071E+02	3.183E+02	3.215E+02
.8 :	1.000E+00	2.270E+00	5.123E+00	1.141E+01	2.475E+01	5.098E+01	9.280E+01	1.285E+02	1.431E+02	1.476E+02	1.488E+02
1 :	1.000E+00	1.989E+00	3.940E+00	7.728E+00	1.486E+01	2.731E+01	4.505E+01	5.908E+01	6.464E+01	6.629E+01	6.673E+01
1.2 :	1.000E+00	1.737E+00	3.008E+00	5.177E+00	8.775E+00	1.435E+01	2.142E+01	2.660E+01	2.859E+01	2.917E+01	2.933E+01
1.4 :	1.000E+00	1.514E+00	2.289E+00	3.446E+00	5.133E+00	7.440E+00	1.005E+01	1.182E+01	1.248E+01	1.267E+01	1.272E+01
1.6 :	1.000E+00	1.319E+00	1.738E+00	2.285E+00	2.985E+00	3.828E+00	4.678E+00	5.211E+00	5.405E+00	5.459E+00	5.473E+00
1.8 :	1.000E+00	1.149E+00	1.319E+00	1.513E+00	1.729E+00	1.960E+00	2.166E+00	2.286E+00	2.340E+00	2.343E+00	2.343E+00
2 :	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = .03

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2 :	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8 :	1.000E+00	2.216E+00	5.866E+00	1.865E+01	7.138E+01	2.812E+02	4.880E+02	6.534E+02	8.075E+02	9.647E+02	1.143E+03
-1.6 :	1.000E+00	2.422E+00	7.449E+00	2.705E+01	1.119E+02	4.279E+02	7.195E+02	9.170E+02	1.065E+03	1.180E+03	1.274E+03
-1.4 :	1.000E+00	2.623E+00	8.817E+00	3.378E+01	1.411E+02	5.073E+02	8.285E+02	1.022E+03	1.147E+03	1.229E+03	1.283E+03
-1.2 :	1.000E+00	2.816E+00	1.000E+01	3.910E+01	1.607E+02	5.386E+02	8.548E+02	1.027E+03	1.125E+03	1.180E+03	1.212E+03
-1 :	1.000E+00	3.000E+00	1.100E+01	4.300E+01	1.710E+02	5.328E+02	8.221E+02	9.667E+02	1.039E+03	1.075E+03	1.093E+03
-.8 :	1.000E+00	3.164E+00	1.175E+01	4.523E+01	1.719E+02	4.971E+02	7.455E+02	8.606E+02	9.122E+02	9.350E+02	9.451E+02
-.6 :	1.000E+00	3.290E+00	1.217E+01	4.543E+01	1.632E+02	4.378E+02	6.378E+02	7.247E+02	7.600E+02	7.740E+02	7.794E+02
-.4 :	1.000E+00	3.356E+00	1.213E+01	4.331E+01	1.455E+02	3.618E+02	5.126E+02	5.743E+02	5.974E+02	6.057E+02	6.086E+02
-.2 :	1.000E+00	3.339E+00	1.157E+01	3.896E+01	1.214E+02	2.797E+02	3.851E+02	4.262E+02	4.405E+02	4.452E+02	4.468E+02
0 :	1.000E+00	3.229E+00	1.052E+01	3.298E+01	9.462E+01	2.019E+02	2.702E+02	2.957E+02	3.040E+02	3.066E+02	3.074E+02
.2 :	1.000E+00	3.036E+00	9.146E+00	2.635E+01	6.912E+01	1.366E+02	1.776E+02	1.924E+02	1.970E+02	1.983E+02	1.987E+02
.4 :	1.000E+00	2.787E+00	7.640E+00	2.002E+01	4.774E+01	8.737E+01	1.104E+02	1.184E+02	1.209E+02	1.215E+02	1.217E+02
.6 :	1.000E+00	2.510E+00	6.186E+00	1.461E+01	3.153E+01	5.345E+01	6.564E+01	6.975E+01	7.095E+01	7.128E+01	7.136E+01
.8 :	1.000E+00	2.232E+00	4.897E+00	1.035E+01	2.013E+01	3.162E+01	3.772E+01	3.973E+01	4.030E+01	4.045E+01	4.049E+01
1 :	1.000E+00	1.967E+00	3.818E+00	7.179E+00	1.254E+01	1.826E+01	2.116E+01	2.210E+01	2.236E+01	2.242E+01	2.244E+01
1.2 :	1.000E+00	1.725E+00	2.946E+00	4.907E+00	7.688E+00	1.037E+01	1.168E+01	1.209E+01	1.220E+01	1.223E+01	1.224E+01
1.4 :	1.000E+00	1.508E+00	2.259E+00	3.322E+00	4.658E+00	5.829E+00	6.371E+00	6.538E+00	6.583E+00	6.595E+00	6.598E+00
1.6 :	1.000E+00	1.316E+00	1.725E+00	2.234E+00	2.801E+00	3.251E+00	3.450E+00	3.511E+00	3.527E+00	3.531E+00	3.532E+00
1.8 :	1.000E+00	1.148E+00	1.314E+00	1.497E+00	1.676E+00	1.805E+00	1.860E+00	1.876E+00	1.881E+00	1.882E+00	1.882E+00
2 :	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = .1

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	2.229E+00	6.118E+00	2.327E+01	4.419E+01	5.757E+01	6.871E+01	7.831E+01	8.665E+01	9.393E+01	1.003E+02
-1.6	1.000E+00	2.448E+00	7.866E+00	3.138E+01	6.097E+01	7.934E+01	9.267E+01	1.026E+02	1.101E+02	1.158E+02	1.201E+02
-1.4	1.000E+00	2.654E+00	9.270E+00	3.626E+01	6.919E+01	8.832E+01	1.005E+02	1.084E+02	1.136E+02	1.170E+02	1.193E+02
-1.2	1.000E+00	2.842E+00	1.032E+01	3.853E+01	7.133E+01	8.911E+01	9.909E+01	1.048E+02	1.080E+02	1.099E+02	1.109E+02
-1	1.000E+00	3.000E+00	1.100E+01	3.863E+01	6.906E+01	8.453E+01	9.227E+01	9.613E+01	9.807E+01	9.903E+01	9.952E+01
-.8	1.000E+00	3.117E+00	1.127E+01	3.694E+01	6.365E+01	7.645E+01	8.222E+01	8.478E+01	8.591E+01	8.640E+01	8.662E+01
-.6	1.000E+00	3.180E+00	1.114E+01	3.385E+01	5.617E+01	6.632E+01	7.049E+01	7.215E+01	7.280E+01	7.305E+01	7.314E+01
-.4	1.000E+00	3.179E+00	1.061E+01	2.981E+01	4.760E+01	5.532E+01	5.825E+01	5.931E+01	5.968E+01	5.981E+01	5.985E+01
-.2	1.000E+00	3.111E+00	9.757E+00	2.528E+01	3.884E+01	4.447E+01	4.647E+01	4.714E+01	4.735E+01	4.742E+01	4.744E+01
0	1.000E+00	2.980E+00	8.683E+00	2.069E+01	3.059E+01	3.455E+01	3.587E+01	3.629E+01	3.641E+01	3.645E+01	3.646E+01
.2	1.000E+00	2.799E+00	7.502E+00	1.642E+01	2.335E+01	2.689E+01	2.899E+01	2.714E+01	2.721E+01	2.721E+01	2.724E+01
.4	1.000E+00	2.584E+00	6.317E+00	1.269E+01	1.735E+01	1.910E+01	1.964E+01	1.979E+01	1.984E+01	1.985E+01	1.985E+01
.6	1.000E+00	2.352E+00	5.208E+00	9.591E+00	1.262E+01	1.372E+01	1.405E+01	1.414E+01	1.416E+01	1.417E+01	1.417E+01
.8	1.000E+00	2.116E+00	4.221E+00	7.125E+00	9.013E+00	9.683E+00	9.878E+00	9.931E+00	9.944E+00	9.948E+00	9.949E+00
1	1.000E+00	1.888E+00	3.377E+00	5.223E+00	6.354E+00	6.745E+00	6.857E+00	6.886E+00	6.894E+00	6.896E+00	6.896E+00
1.2	1.000E+00	1.674E+00	2.676E+00	3.791E+00	4.435E+00	4.653E+00	4.714E+00	4.730E+00	4.734E+00	4.735E+00	4.735E+00
1.4	1.000E+00	1.477E+00	2.105E+00	2.732E+00	3.074E+00	3.187E+00	3.218E+00	3.226E+00	3.228E+00	3.228E+00	3.229E+00
1.6	1.000E+00	1.299E+00	1.647E+00	1.960E+00	2.120E+00	2.172E+00	2.186E+00	2.189E+00	2.190E+00	2.191E+00	2.191E+00
1.8	1.000E+00	1.141E+00	1.285E+00	1.401E+00	1.457E+00	1.475E+00	1.480E+00	1.481E+00	1.481E+00	1.482E+00	1.482E+00
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = .3

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	2.388E+00	5.582E+00	9.024E+00	1.341E+01	1.980E+01	2.987E+01	4.649E+01	7.464E+01	1.230E+02	2.065E+02
-1.6	1.000E+00	2.669E+00	6.503E+00	9.614E+00	1.237E+01	1.521E+01	1.854E+01	2.285E+01	2.882E+01	3.746E+01	5.023E+01
-1.4	1.000E+00	2.857E+00	6.937E+00	9.770E+00	1.171E+01	1.317E+01	1.440E+01	1.557E+01	1.681E+01	1.825E+01	2.002E+01
-1.2	1.000E+00	2.964E+00	7.013E+00	9.562E+00	1.102E+01	1.188E+01	1.242E+01	1.278E+01	1.305E+01	1.328E+01	1.348E+01
-1	1.000E+00	3.000E+00	6.833E+00	9.083E+00	1.021E+01	1.077E+01	1.105E+01	1.119E+01	1.126E+01	1.130E+01	1.132E+01
-.8	1.000E+00	2.976E+00	6.477E+00	8.422E+00	9.300E+00	9.686E+00	9.850E+00	9.918E+00	9.944E+00	9.952E+00	9.953E+00
-.6	1.000E+00	2.902E+00	6.008E+00	7.655E+00	8.399E+00	8.610E+00	8.713E+00	8.750E+00	8.763E+00	8.766E+00	8.766E+00
-.4	1.000E+00	2.789E+00	5.478E+00	6.844E+00	7.374E+00	7.566E+00	7.633E+00	7.655E+00	7.662E+00	7.664E+00	7.664E+00
-.2	1.000E+00	2.648E+00	4.923E+00	6.036E+00	6.443E+00	6.580E+00	6.624E+00	6.638E+00	6.642E+00	6.643E+00	6.643E+00
0	1.000E+00	2.487E+00	4.374E+00	5.264E+00	5.573E+00	5.671E+00	5.701E+00	5.709E+00	5.712E+00	5.712E+00	5.713E+00
.2	1.000E+00	2.316E+00	3.850E+00	4.548E+00	4.781E+00	4.851E+00	4.871E+00	4.876E+00	4.878E+00	4.878E+00	4.879E+00
.4	1.000E+00	2.140E+00	3.362E+00	3.900E+00	4.073E+00	4.123E+00	4.137E+00	4.140E+00	4.141E+00	4.141E+00	4.142E+00
.6	1.000E+00	1.967E+00	2.919E+00	3.324E+00	3.450E+00	3.486E+00	3.495E+00	3.498E+00	3.498E+00	3.499E+00	3.499E+00
.8	1.000E+00	1.799E+00	2.522E+00	2.820E+00	2.910E+00	2.935E+00	2.942E+00	2.943E+00	2.944E+00	2.944E+00	2.944E+00
1	1.000E+00	1.639E+00	2.171E+00	2.383E+00	2.446E+00	2.463E+00	2.468E+00	2.469E+00	2.469E+00	2.469E+00	2.469E+00
1.2	1.000E+00	1.489E+00	1.864E+00	2.009E+00	2.051E+00	2.062E+00	2.065E+00	2.066E+00	2.066E+00	2.066E+00	2.066E+00
1.4	1.000E+00	1.350E+00	1.597E+00	1.690E+00	1.716E+00	1.723E+00	1.725E+00	1.726E+00	1.726E+00	1.726E+00	1.726E+00
1.6	1.000E+00	1.223E+00	1.367E+00	1.420E+00	1.434E+00	1.438E+00	1.439E+00	1.440E+00	1.440E+00	1.440E+00	1.440E+00
1.8	1.000E+00	1.106E+00	1.169E+00	1.192E+00	1.198E+00	1.200E+00	1.200E+00	1.200E+00	1.200E+00	1.200E+00	1.200E+00
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 1

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-0.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-0.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-0.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-0.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
0	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
1.2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
1.4	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
1.6	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
1.8	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 1.01

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.300E+00	1.771E+00	2.545E+00	3.852E+00	6.095E+00	9.968E+00	1.669E+01	2.836E+01	4.867E+01	8.401E+01
-1.6	1.000E+00	1.105E+00	1.234E+00	1.407E+00	1.650E+00	2.004E+00	2.532E+00	3.324E+00	4.518E+00	6.324E+00	9.057E+00
-1.4	1.000E+00	1.037E+00	1.073E+00	1.110E+00	1.155E+00	1.209E+00	1.278E+00	1.367E+00	1.484E+00	1.637E+00	1.838E+00
-1.2	1.000E+00	1.011E+00	1.019E+00	1.026E+00	1.032E+00	1.039E+00	1.046E+00	1.054E+00	1.063E+00	1.073E+00	1.084E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-0.8	1.000E+00	9.956E-01	9.933E-01	9.921E-01	9.913E-01	9.908E-01	9.905E-01	9.902E-01	9.899E-01	9.897E-01	9.896E-01
-0.6	1.000E+00	9.938E-01	9.910E-01	9.898E-01	9.892E-01	9.889E-01	9.887E-01	9.886E-01	9.885E-01	9.884E-01	9.884E-01
-0.4	1.000E+00	9.932E-01	9.905E-01	9.895E-01	9.891E-01	9.889E-01	9.888E-01	9.888E-01	9.887E-01	9.887E-01	9.887E-01
-0.2	1.000E+00	9.931E-01	9.907E-01	9.899E-01	9.896E-01	9.895E-01	9.895E-01	9.894E-01	9.894E-01	9.894E-01	9.894E-01
0	1.000E+00	9.934E-01	9.913E-01	9.906E-01	9.904E-01	9.903E-01	9.903E-01	9.903E-01	9.903E-01	9.903E-01	9.903E-01
.2	1.000E+00	9.938E-01	9.920E-01	9.914E-01	9.913E-01	9.912E-01	9.912E-01	9.912E-01	9.912E-01	9.912E-01	9.912E-01
.4	1.000E+00	9.944E-01	9.928E-01	9.923E-01	9.922E-01	9.922E-01	9.922E-01	9.922E-01	9.922E-01	9.922E-01	9.922E-01
.6	1.000E+00	9.950E-01	9.936E-01	9.932E-01	9.932E-01	9.931E-01	9.931E-01	9.931E-01	9.931E-01	9.931E-01	9.931E-01
.8	1.000E+00	9.956E-01	9.945E-01	9.942E-01	9.941E-01	9.941E-01	9.941E-01	9.941E-01	9.941E-01	9.941E-01	9.941E-01
1	1.000E+00	9.963E-01	9.954E-01	9.952E-01	9.951E-01	9.951E-01	9.951E-01	9.951E-01	9.951E-01	9.951E-01	9.951E-01
1.2	1.000E+00	9.970E-01	9.963E-01	9.961E-01	9.961E-01	9.961E-01	9.961E-01	9.961E-01	9.961E-01	9.961E-01	9.961E-01
1.4	1.000E+00	9.978E-01	9.972E-01	9.971E-01	9.970E-01	9.970E-01	9.970E-01	9.970E-01	9.970E-01	9.970E-01	9.970E-01
1.6	1.000E+00	9.985E-01	9.981E-01	9.981E-01	9.980E-01	9.980E-01	9.980E-01	9.980E-01	9.980E-01	9.980E-01	9.980E-01
1.8	1.000E+00	9.992E-01	9.991E-01	9.990E-01	9.990E-01	9.990E-01	9.990E-01	9.990E-01	9.990E-01	9.990E-01	9.990E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 1.1

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.494E+00	2.286E+00	3.604E+00	5.845E+00	9.700E+00	1.637E+01	2.795E+01	4.808E+01	8.311E+01	1.441E+02
-1.6	1.000E+00	1.250E+00	1.573E+00	2.017E+00	2.657E+00	3.601E+00	5.011E+00	7.134E+00	1.034E+01	1.519E+01	2.254E+01
-1.4	1.000E+00	1.119E+00	1.243E+00	1.385E+00	1.554E+00	1.772E+00	2.050E+00	2.412E+00	2.887E+00	3.512E+00	4.336E+00
-1.2	1.000E+00	1.044E+00	1.082E+00	1.117E+00	1.153E+00	1.191E+00	1.234E+00	1.282E+00	1.336E+00	1.398E+00	1.470E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	9.740E-01	9.584E-01	9.484E-01	9.414E-01	9.361E-01	9.318E-01	9.282E-01	9.251E-01	9.225E-01	9.202E-01
-.6	1.000E+00	9.590E-01	9.380E-01	9.268E-01	9.202E-01	9.160E-01	9.132E-01	9.111E-01	9.096E-01	9.085E-01	9.077E-01
-.4	1.000E+00	9.510E-01	9.292E-01	9.193E-01	9.145E-01	9.119E-01	9.105E-01	9.096E-01	9.090E-01	9.087E-01	9.085E-01
-.2	1.000E+00	9.476E-01	9.270E-01	9.189E-01	9.156E-01	9.142E-01	9.135E-01	9.131E-01	9.129E-01	9.128E-01	9.128E-01
0	1.000E+00	9.472E-01	9.287E-01	9.223E-01	9.200E-01	9.192E-01	9.189E-01	9.187E-01	9.187E-01	9.186E-01	9.186E-01
.2	1.000E+00	9.490E-01	9.327E-01	9.277E-01	9.261E-01	9.256E-01	9.254E-01	9.254E-01	9.253E-01	9.253E-01	9.253E-01
.4	1.000E+00	9.523E-01	9.382E-01	9.342E-01	9.331E-01	9.327E-01	9.326E-01	9.326E-01	9.326E-01	9.326E-01	9.326E-01
.6	1.000E+00	9.566E-01	9.447E-01	9.415E-01	9.404E-01	9.404E-01	9.403E-01	9.403E-01	9.403E-01	9.403E-01	9.403E-01
.8	1.000E+00	9.617E-01	9.517E-01	9.492E-01	9.485E-01	9.483E-01	9.483E-01	9.483E-01	9.483E-01	9.483E-01	9.483E-01
1	1.000E+00	9.674E-01	9.592E-01	9.572E-01	9.567E-01	9.566E-01	9.565E-01	9.565E-01	9.565E-01	9.565E-01	9.565E-01
1.2	1.000E+00	9.735E-01	9.670E-01	9.655E-01	9.651E-01	9.650E-01	9.649E-01	9.649E-01	9.649E-01	9.649E-01	9.649E-01
1.4	1.000E+00	9.798E-01	9.751E-01	9.739E-01	9.736E-01	9.735E-01	9.735E-01	9.735E-01	9.735E-01	9.735E-01	9.735E-01
1.6	1.000E+00	9.864E-01	9.832E-01	9.825E-01	9.823E-01	9.822E-01	9.822E-01	9.822E-01	9.822E-01	9.822E-01	9.822E-01
1.8	1.000E+00	9.931E-01	9.916E-01	9.912E-01	9.911E-01	9.911E-01	9.911E-01	9.911E-01	9.911E-01	9.911E-01	9.911E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 2

M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.690E+00	2.862E+00	4.877E+00	8.362E+00	1.441E+01	2.491E+01	4.319E+01	7.500E+01	1.304E+02	2.267E+02
-1.6	1.000E+00	1.445E+00	2.089E+00	3.037E+00	4.453E+00	6.583E+00	9.798E+00	1.466E+01	2.203E+01	3.319E+01	5.010E+01
-1.4	1.000E+00	1.255E+00	1.568E+00	1.964E+00	2.473E+00	3.137E+00	4.006E+00	5.150E+00	6.657E+00	8.644E+00	1.126E+01
-1.2	1.000E+00	1.109E+00	1.224E+00	1.347E+00	1.484E+00	1.638E+00	1.813E+00	2.014E+00	2.244E+00	2.507E+00	2.810E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	9.198E-01	8.591E-01	8.113E-01	7.723E-01	7.395E-01	7.115E-01	6.873E-01	6.664E-01	6.483E-01	6.325E-01
-.6	1.000E+00	8.632E-01	7.746E-01	7.151E-01	6.735E-01	6.435E-01	6.213E-01	6.047E-01	5.922E-01	5.828E-01	5.757E-01
-.4	1.000E+00	8.255E-01	7.283E-01	6.727E-01	6.394E-01	6.188E-01	6.057E-01	5.972E-01	5.917E-01	5.881E-01	5.857E-01
-.2	1.000E+00	8.028E-01	7.080E-01	6.614E-01	6.377E-01	6.251E-01	6.183E-01	6.144E-01	6.123E-01	6.110E-01	6.103E-01
0	1.000E+00	7.922E-01	7.052E-01	6.684E-01	6.523E-01	6.451E-01	6.417E-01	6.401E-01	6.393E-01	6.389E-01	6.387E-01
.2	1.000E+00	7.912E-01	7.143E-01	6.860E-01	6.754E-01	6.713E-01	6.697E-01	6.690E-01	6.687E-01	6.686E-01	6.686E-01
.4	1.000E+00	7.979E-01	7.316E-01	7.101E-01	7.031E-01	7.008E-01	7.000E-01	6.997E-01	6.996E-01	6.996E-01	6.996E-01
.6	1.000E+00	8.106E-01	7.547E-01	7.384E-01	7.337E-01	7.323E-01	7.319E-01	7.318E-01	7.318E-01	7.318E-01	7.318E-01
.8	1.000E+00	8.283E-01	7.820E-01	7.696E-01	7.664E-01	7.655E-01	7.653E-01	7.652E-01	7.652E-01	7.652E-01	7.652E-01
1	1.000E+00	8.500E-01	8.125E-01	8.031E-01	8.008E-01	8.002E-01	8.000E-01	8.000E-01	8.000E-01	8.000E-01	8.000E-01
1.2	1.000E+00	8.750E-01	8.456E-01	8.386E-01	8.369E-01	8.365E-01	8.364E-01	8.364E-01	8.363E-01	8.363E-01	8.363E-01
1.4	1.000E+00	9.029E-01	8.811E-01	8.760E-01	8.747E-01	8.745E-01	8.744E-01	8.744E-01	8.744E-01	8.744E-01	8.744E-01
1.6	1.000E+00	9.331E-01	9.186E-01	9.153E-01	9.145E-01	9.143E-01	9.142E-01	9.142E-01	9.142E-01	9.142E-01	9.142E-01
1.8	1.000E+00	9.656E-01	9.582E-01	9.566E-01	9.562E-01	9.561E-01	9.560E-01	9.560E-01	9.560E-01	9.560E-01	9.560E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2, M, r, mu) for r = 4

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.728E+00	2.986E+00	5.170E+00	8.964E+00	1.556E+01	2.705E+01	4.704E+01	8.185E+01	1.425E+02	2.480E+02
-1.6	1.000E+00	1.494E+00	2.232E+00	3.341E+00	5.014E+00	7.544E+00	1.137E+01	1.718E+01	2.597E+01	3.929E+01	5.948E+01
-1.4	1.000E+00	1.296E+00	1.678E+00	2.174E+00	2.823E+00	3.676E+00	4.799E+00	6.279E+00	8.232E+00	1.081E+01	1.420E+01
-1.2	1.000E+00	1.132E+00	1.279E+00	1.444E+00	1.631E+00	1.844E+00	2.088E+00	2.367E+00	2.688E+00	3.057E+00	3.480E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.967E-01	8.123E-01	7.420E-01	6.824E-01	6.313E-01	5.872E-01	5.489E-01	5.157E-01	4.868E-01	4.616E-01
-.6	1.000E+00	8.198E-01	6.934E-01	6.031E-01	5.370E-01	4.880E-01	4.513E-01	4.237E-01	4.028E-01	3.870E-01	3.750E-01
-.4	1.000E+00	7.663E-01	6.256E-01	5.392E-01	4.849E-01	4.502E-01	4.276E-01	4.128E-01	4.032E-01	3.968E-01	3.926E-01
-.2	1.000E+00	7.330E-01	5.946E-01	5.217E-01	4.824E-01	4.607E-01	4.486E-01	4.417E-01	4.378E-01	4.353E-01	4.342E-01
0	1.000E+00	7.167E-01	5.898E-01	5.323E-01	5.057E-01	4.932E-01	4.871E-01	4.841E-01	4.826E-01	4.819E-01	4.815E-01
.2	1.000E+00	7.144E-01	6.030E-01	5.595E-01	5.422E-01	5.351E-01	5.322E-01	5.310E-01	5.305E-01	5.303E-01	5.302E-01
.4	1.000E+00	7.232E-01	6.285E-01	5.962E-01	5.851E-01	5.812E-01	5.798E-01	5.794E-01	5.792E-01	5.791E-01	5.791E-01
.6	1.000E+00	7.409E-01	6.621E-01	6.384E-01	6.312E-01	6.291E-01	6.284E-01	6.282E-01	6.281E-01	6.281E-01	6.281E-01
.8	1.000E+00	7.654E-01	7.013E-01	6.839E-01	6.792E-01	6.779E-01	6.773E-01	6.773E-01	6.773E-01	6.773E-01	6.773E-01
1	1.000E+00	7.955E-01	7.443E-01	7.315E-01	7.283E-01	7.275E-01	7.273E-01	7.273E-01	7.273E-01	7.273E-01	7.273E-01
1.2	1.000E+00	8.298E-01	7.903E-01	7.810E-01	7.788E-01	7.782E-01	7.781E-01	7.781E-01	7.781E-01	7.781E-01	7.781E-01
1.4	1.000E+00	8.679E-01	8.389E-01	8.323E-01	8.307E-01	8.303E-01	8.303E-01	8.302E-01	8.302E-01	8.302E-01	8.302E-01
1.6	1.000E+00	9.090E-01	8.899E-01	8.899E-01	8.856E-01	8.843E-01	8.843E-01	8.843E-01	8.843E-01	8.843E-01	8.843E-01
1.8	1.000E+00	9.531E-01	9.435E-01	9.414E-01	9.408E-01	9.407E-01	9.407E-01	9.407E-01	9.407E-01	9.407E-01	9.407E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2, M, r, mu) for r = 8

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.737E+00	3.019E+00	5.247E+00	9.126E+00	1.588E+01	2.763E+01	4.809E+01	8.371E+01	1.457E+02	2.537E+02
-1.6	1.000E+00	1.509E+00	2.276E+00	3.436E+00	5.191E+00	7.850E+00	1.188E+01	1.798E+01	2.723E+01	4.126E+01	6.251E+01
-1.4	1.000E+00	1.311E+00	1.717E+00	2.251E+00	2.952E+00	3.877E+00	5.096E+00	6.704E+00	8.825E+00	1.162E+01	1.532E+01
-1.2	1.000E+00	1.142E+00	1.302E+00	1.484E+00	1.693E+00	1.932E+00	2.206E+00	2.520E+00	2.882E+00	3.296E+00	3.773E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.851E-01	7.883E-01	7.059E-01	6.351E-01	5.739E-01	5.208E-01	4.747E-01	4.346E-01	3.997E-01	3.693E-01
-.6	1.000E+00	7.965E-01	6.489E-01	5.405E-01	4.600E-01	3.996E-01	3.541E-01	3.198E-01	2.938E-01	2.741E-01	2.592E-01
-.4	1.000E+00	7.331E-01	5.666E-01	4.614E-01	3.939E-01	3.502E-01	3.216E-01	3.028E-01	2.904E-01	2.823E-01	2.769E-01
-.2	1.000E+00	6.931E-01	5.284E-01	4.390E-01	3.896E-01	3.620E-01	3.463E-01	3.374E-01	3.323E-01	3.294E-01	3.277E-01
0	1.000E+00	6.738E-01	5.230E-01	4.526E-01	4.191E-01	4.030E-01	3.952E-01	3.913E-01	3.893E-01	3.884E-01	3.879E-01
.2	1.000E+00	6.719E-01	5.405E-01	4.876E-01	4.661E-01	4.572E-01	4.534E-01	4.518E-01	4.511E-01	4.508E-01	4.507E-01
.4	1.000E+00	6.838E-01	5.732E-01	5.347E-01	5.212E-01	5.164E-01	5.146E-01	5.140E-01	5.137E-01	5.136E-01	5.136E-01
.6	1.000E+00	7.061E-01	6.155E-01	5.878E-01	5.793E-01	5.767E-01	5.759E-01	5.756E-01	5.756E-01	5.755E-01	5.755E-01
.8	1.000E+00	7.361E-01	6.634E-01	6.436E-01	6.382E-01	6.367E-01	6.363E-01	6.362E-01	6.361E-01	6.361E-01	6.361E-01
1	1.000E+00	7.717E-01	7.147E-01	7.004E-01	6.968E-01	6.959E-01	6.957E-01	6.957E-01	6.957E-01	6.957E-01	6.957E-01
1.2	1.000E+00	8.115E-01	7.680E-01	7.578E-01	7.554E-01	7.548E-01	7.546E-01	7.546E-01	7.546E-01	7.546E-01	7.546E-01
1.4	1.000E+00	8.546E-01	8.230E-01	8.159E-01	8.142E-01	8.138E-01	8.137E-01	8.137E-01	8.137E-01	8.137E-01	8.137E-01
1.6	1.000E+00	9.005E-01	8.797E-01	8.751E-01	8.741E-01	8.738E-01	8.738E-01	8.737E-01	8.737E-01	8.737E-01	8.737E-01
1.8	1.000E+00	9.489E-01	9.386E-01	9.363E-01	9.357E-01	9.356E-01	9.356E-01	9.356E-01	9.356E-01	9.356E-01	9.356E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 16

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.740E+00	3.028E+00	5.269E+00	9.171E+00	1.596E+01	2.779E+01	4.838E+01	8.424E+01	1.467E+02	2.553E+02
-1.6	1.000E+00	1.513E+00	2.290E+00	3.467E+00	5.249E+00	7.950E+00	1.204E+01	1.825E+01	2.765E+01	4.191E+01	6.351E+01
-1.4	1.000E+00	1.316E+00	1.732E+00	2.280E+00	3.001E+00	3.953E+00	5.209E+00	6.865E+00	9.051E+00	1.194E+01	1.574E+01
-1.2	1.000E+00	1.146E+00	1.312E+00	1.502E+00	1.720E+00	1.970E+00	2.257E+00	2.587E+00	2.966E+00	3.401E+00	3.901E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.788E-01	7.752E-01	6.860E-01	6.089E-01	5.420E-01	4.839E-01	4.334E-01	3.894E-01	3.511E-01	3.178E-01
-.6	1.000E+00	7.826E-01	6.220E-01	5.026E-01	4.132E-01	3.458E-01	2.949E-01	2.564E-01	2.273E-01	2.052E-01	1.885E-01
-.4	1.000E+00	7.119E-01	5.287E-01	4.112E-01	3.350E-01	2.853E-01	2.527E-01	2.313E-01	2.172E-01	2.078E-01	2.017E-01
-.2	1.000E+00	6.668E-01	4.845E-01	3.838E-01	3.275E-01	2.958E-01	2.777E-01	2.674E-01	2.615E-01	2.581E-01	2.562E-01
0	1.000E+00	6.455E-01	4.785E-01	3.992E-01	3.612E-01	3.426E-01	3.335E-01	3.290E-01	3.267E-01	3.256E-01	3.250E-01
.2	1.000E+00	6.446E-01	5.001E-01	4.410E-01	4.167E-01	4.065E-01	4.021E-01	4.003E-01	3.995E-01	3.991E-01	3.990E-01
.4	1.000E+00	6.597E-01	5.395E-01	4.970E-01	4.819E-01	4.765E-01	4.745E-01	4.738E-01	4.735E-01	4.734E-01	4.734E-01
.6	1.000E+00	6.865E-01	5.891E-01	5.591E-01	5.498E-01	5.469E-01	5.460E-01	5.458E-01	5.457E-01	5.456E-01	5.456E-01
.8	1.000E+00	7.210E-01	6.439E-01	6.227E-01	6.169E-01	6.154E-01	6.149E-01	6.148E-01	6.148E-01	6.148E-01	6.148E-01
1	1.000E+00	7.606E-01	7.008E-01	6.858E-01	6.821E-01	6.812E-01	6.809E-01	6.809E-01	6.809E-01	6.809E-01	6.809E-01
1.2	1.000E+00	8.038E-01	7.586E-01	7.480E-01	7.455E-01	7.449E-01	7.447E-01	7.447E-01	7.447E-01	7.447E-01	7.447E-01
1.4	1.000E+00	8.495E-01	8.169E-01	8.096E-01	8.079E-01	8.075E-01	8.074E-01	8.074E-01	8.074E-01	8.074E-01	8.074E-01
1.6	1.000E+00	8.974E-01	8.762E-01	8.715E-01	8.704E-01	8.701E-01	8.701E-01	8.701E-01	8.701E-01	8.701E-01	8.701E-01
1.8	1.000E+00	9.475E-01	9.370E-01	9.346E-01	9.341E-01	9.340E-01	9.339E-01	9.339E-01	9.339E-01	9.339E-01	9.339E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 32

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.030E+00	5.276E+00	9.184E+00	1.599E+01	2.784E+01	4.847E+01	8.439E+01	1.469E+02	2.558E+02
-1.6	1.000E+00	1.515E+00	2.295E+00	3.477E+00	5.269E+00	7.984E+00	1.210E+01	1.834E+01	2.779E+01	4.212E+01	6.384E+01
-1.4	1.000E+00	1.318E+00	1.738E+00	2.291E+00	3.020E+00	3.982E+00	5.252E+00	6.927E+00	9.137E+00	1.205E+01	1.590E+01
-1.2	1.000E+00	1.147E+00	1.316E+00	1.510E+00	1.732E+00	1.987E+00	2.280E+00	2.616E+00	3.003E+00	3.447E+00	3.957E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.753E-01	7.677E-01	6.747E-01	5.941E-01	5.240E-01	4.630E-01	4.100E-01	3.639E-01	3.237E-01	2.887E-01
-.6	1.000E+00	7.738E-01	6.052E-01	4.789E-01	3.838E-01	3.120E-01	2.577E-01	2.166E-01	1.855E-01	1.619E-01	1.441E-01
-.4	1.000E+00	6.975E-01	5.030E-01	3.770E-01	2.950E-01	2.412E-01	2.059E-01	1.826E-01	1.673E-01	1.572E-01	1.505E-01
-.2	1.000E+00	6.480E-01	4.532E-01	3.445E-01	2.833E-01	2.486E-01	2.288E-01	2.175E-01	2.110E-01	2.073E-01	2.051E-01
0	1.000E+00	6.251E-01	4.465E-01	3.609E-01	3.194E-01	2.991E-01	2.891E-01	2.841E-01	2.816E-01	2.804E-01	2.798E-01
.2	1.000E+00	6.256E-01	4.718E-01	4.085E-01	3.821E-01	3.710E-01	3.663E-01	3.642E-01	3.634E-01	3.630E-01	3.628E-01
.4	1.000E+00	6.440E-01	5.173E-01	4.723E-01	4.561E-01	4.503E-01	4.481E-01	4.473E-01	4.471E-01	4.469E-01	4.469E-01
.6	1.000E+00	6.747E-01	5.733E-01	5.418E-01	5.321E-01	5.290E-01	5.280E-01	5.277E-01	5.276E-01	5.276E-01	5.276E-01
.8	1.000E+00	7.128E-01	6.333E-01	6.114E-01	6.054E-01	6.038E-01	6.033E-01	6.032E-01	6.032E-01	6.031E-01	6.031E-01
1	1.000E+00	7.553E-01	6.941E-01	6.788E-01	6.750E-01	6.740E-01	6.738E-01	6.737E-01	6.737E-01	6.737E-01	6.737E-01
1.2	1.000E+00	8.004E-01	7.545E-01	7.438E-01	7.412E-01	7.406E-01	7.405E-01	7.404E-01	7.404E-01	7.404E-01	7.404E-01
1.4	1.000E+00	8.475E-01	8.146E-01	8.072E-01	8.055E-01	8.051E-01	8.050E-01	8.050E-01	8.049E-01	8.049E-01	8.049E-01
1.6	1.000E+00	8.964E-01	8.750E-01	8.702E-01	8.691E-01	8.689E-01	8.688E-01	8.688E-01	8.688E-01	8.688E-01	8.688E-01
1.8	1.000E+00	9.471E-01	9.365E-01	9.341E-01	9.336E-01	9.334E-01	9.334E-01	9.334E-01	9.334E-01	9.334E-01	9.334E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 64

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.031E+00	5.277E+00	9.188E+00	1.600E+01	2.785E+01	4.849E+01	8.443E+01	1.470E+02	2.559E+02
-1.6	1.000E+00	1.515E+00	2.297E+00	3.481E+00	5.275E+00	7.995E+00	1.212E+01	1.836E+01	2.784E+01	4.219E+01	6.395E+01
-1.4	1.000E+00	1.319E+00	1.740E+00	2.295E+00	3.027E+00	3.993E+00	5.268E+00	6.950E+00	9.170E+00	1.210E+01	1.596E+01
-1.2	1.000E+00	1.148E+00	1.318E+00	1.513E+00	1.737E+00	1.994E+00	2.290E+00	2.629E+00	3.019E+00	3.467E+00	3.981E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.733E-01	7.635E-01	6.683E-01	5.856E-01	5.137E-01	4.512E-01	3.967E-01	3.493E-01	3.081E-01	2.722E-01
-.6	1.000E+00	7.683E-01	5.945E-01	4.637E-01	3.650E-01	2.904E-01	2.340E-01	1.912E-01	1.588E-01	1.343E-01	1.157E-01
-.4	1.000E+00	6.875E-01	4.849E-01	3.531E-01	2.668E-01	2.102E-01	1.730E-01	1.485E-01	1.323E-01	1.216E-01	1.146E-01
-.2	1.000E+00	6.341E-01	4.299E-01	3.152E-01	2.504E-01	2.135E-01	1.924E-01	1.804E-01	1.734E-01	1.695E-01	1.672E-01
0	1.000E+00	6.098E-01	4.225E-01	3.320E-01	2.880E-01	2.663E-01	2.556E-01	2.503E-01	2.476E-01	2.463E-01	2.456E-01
.2	1.000E+00	6.116E-01	4.512E-01	3.846E-01	3.568E-01	3.450E-01	3.400E-01	3.378E-01	3.369E-01	3.363E-01	3.363E-01
.4	1.000E+00	6.332E-01	5.021E-01	4.553E-01	4.384E-01	4.323E-01	4.300E-01	4.292E-01	4.289E-01	4.288E-01	4.287E-01
.6	1.000E+00	6.673E-01	5.634E-01	5.310E-01	5.210E-01	5.178E-01	5.168E-01	5.165E-01	5.164E-01	5.164E-01	5.164E-01
.8	1.000E+00	7.082E-01	6.274E-01	6.051E-01	5.990E-01	5.973E-01	5.969E-01	5.967E-01	5.967E-01	5.967E-01	5.967E-01
1	1.000E+00	7.526E-01	6.908E-01	6.753E-01	6.714E-01	6.705E-01	6.702E-01	6.702E-01	6.702E-01	6.702E-01	6.702E-01
1.2	1.000E+00	7.990E-01	7.527E-01	7.419E-01	7.394E-01	7.388E-01	7.386E-01	7.386E-01	7.386E-01	7.386E-01	7.386E-01
1.4	1.000E+00	8.468E-01	8.137E-01	8.063E-01	8.045E-01	8.041E-01	8.040E-01	8.040E-01	8.040E-01	8.040E-01	8.040E-01
1.6	1.000E+00	8.960E-01	8.745E-01	8.698E-01	8.687E-01	8.684E-01	8.684E-01	8.684E-01	8.684E-01	8.684E-01	8.684E-01
1.8	1.000E+00	9.470E-01	9.363E-01	9.340E-01	9.334E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 128

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.189E+00	1.600E+01	2.786E+01	4.850E+01	8.444E+01	1.470E+02	2.560E+02
-1.6	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.277E+00	7.998E+00	1.212E+01	1.837E+01	2.785E+01	4.221E+01	6.398E+01
-1.4	1.000E+00	1.319E+00	1.741E+00	2.296E+00	3.030E+00	3.997E+00	5.274E+00	6.959E+00	9.182E+00	1.212E+01	1.599E+01
-1.2	1.000E+00	1.148E+00	1.319E+00	1.515E+00	1.739E+00	1.998E+00	2.294E+00	2.635E+00	3.026E+00	3.475E+00	3.992E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.721E-01	7.611E-01	6.647E-01	5.808E-01	5.079E-01	4.444E-01	3.891E-01	3.410E-01	2.992E-01	2.627E-01
-.6	1.000E+00	7.647E-01	5.875E-01	4.539E-01	3.529E-01	2.764E-01	2.186E-01	1.747E-01	1.415E-01	1.164E-01	9.729E-02
-.4	1.000E+00	6.803E-01	4.720E-01	3.359E-01	2.467E-01	1.881E-01	1.494E-01	1.240E-01	1.072E-01	9.614E-02	8.884E-02
-.2	1.000E+00	6.235E-01	4.121E-01	2.928E-01	2.252E-01	1.866E-01	1.646E-01	1.519E-01	1.447E-01	1.405E-01	1.381E-01
0	1.000E+00	5.978E-01	4.036E-01	3.094E-01	2.634E-01	2.407E-01	2.294E-01	2.238E-01	2.210E-01	2.196E-01	2.189E-01
.2	1.000E+00	6.011E-01	4.355E-01	3.665E-01	3.376E-01	3.253E-01	3.200E-01	3.177E-01	3.168E-01	3.163E-01	3.162E-01
.4	1.000E+00	6.256E-01	4.914E-01	4.433E-01	4.259E-01	4.196E-01	4.173E-01	4.164E-01	4.161E-01	4.160E-01	4.159E-01
.6	1.000E+00	6.626E-01	5.571E-01	5.242E-01	5.139E-01	5.107E-01	5.097E-01	5.094E-01	5.093E-01	5.092E-01	5.092E-01
.8	1.000E+00	7.056E-01	6.240E-01	6.016E-01	5.954E-01	5.937E-01	5.932E-01	5.931E-01	5.930E-01	5.930E-01	5.930E-01
1	1.000E+00	7.513E-01	6.891E-01	6.736E-01	6.697E-01	6.687E-01	6.685E-01	6.684E-01	6.684E-01	6.684E-01	6.684E-01
1.2	1.000E+00	7.983E-01	7.520E-01	7.411E-01	7.386E-01	7.380E-01	7.378E-01	7.378E-01	7.378E-01	7.378E-01	7.378E-01
1.4	1.000E+00	8.465E-01	8.133E-01	8.059E-01	8.042E-01	8.038E-01	8.037E-01	8.037E-01	8.037E-01	8.037E-01	8.037E-01
1.6	1.000E+00	8.959E-01	8.744E-01	8.696E-01	8.686E-01	8.683E-01	8.682E-01	8.682E-01	8.682E-01	8.682E-01	8.682E-01
1.8	1.000E+00	9.469E-01	9.363E-01	9.339E-01	9.334E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 256

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.189E+00	1.600E+01	2.786E+01	4.850E+01	8.445E+01	1.470E+02	2.560E+02
-1.6	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.278E+00	7.999E+00	1.212E+01	1.838E+01	2.786E+01	4.222E+01	6.399E+01
-1.4	1.000E+00	1.319E+00	1.741E+00	2.297E+00	3.031E+00	3.999E+00	5.277E+00	6.962E+00	9.187E+00	1.212E+01	1.599E+01
-1.2	1.000E+00	1.149E+00	1.319E+00	1.515E+00	1.740E+00	1.999E+00	2.296E+00	2.637E+00	3.029E+00	3.479E+00	3.996E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.714E-01	7.597E-01	6.626E-01	5.781E-01	5.045E-01	4.405E-01	3.848E-01	3.363E-01	2.941E-01	2.573E-01
-.6	1.000E+00	7.623E-01	5.830E-01	4.475E-01	3.450E-01	2.674E-01	2.086E-01	1.640E-01	1.303E-01	1.047E-01	8.533E-02
-.4	1.000E+00	6.750E-01	4.626E-01	3.235E-01	2.321E-01	1.720E-01	1.323E-01	1.062E-01	8.899E-02	7.763E-02	7.014E-02
-.2	1.000E+00	6.152E-01	3.981E-01	2.753E-01	2.054E-01	1.655E-01	1.427E-01	1.296E-01	1.221E-01	1.178E-01	1.153E-01
0	1.000E+00	5.882E-01	3.885E-01	2.913E-01	2.436E-01	2.201E-01	2.084E-01	2.026E-01	1.997E-01	1.982E-01	1.975E-01
.2	1.000E+00	5.929E-01	4.233E-01	3.525E-01	3.226E-01	3.099E-01	3.045E-01	3.022E-01	3.011E-01	3.007E-01	3.005E-01
.4	1.000E+00	6.202E-01	4.837E-01	4.347E-01	4.170E-01	4.105E-01	4.081E-01	4.072E-01	4.069E-01	4.068E-01	4.067E-01
.6	1.000E+00	6.596E-01	5.530E-01	5.198E-01	5.094E-01	5.061E-01	5.051E-01	5.047E-01	5.046E-01	5.046E-01	5.046E-01
.8	1.000E+00	7.042E-01	6.221E-01	5.995E-01	5.933E-01	5.916E-01	5.911E-01	5.910E-01	5.910E-01	5.910E-01	5.910E-01
1	1.000E+00	7.507E-01	6.883E-01	6.727E-01	6.688E-01	6.679E-01	6.676E-01	6.676E-01	6.675E-01	6.675E-01	6.675E-01
1.2	1.000E+00	7.981E-01	7.516E-01	7.408E-01	7.382E-01	7.376E-01	7.375E-01	7.374E-01	7.374E-01	7.374E-01	7.374E-01
1.4	1.000E+00	8.464E-01	8.132E-01	8.058E-01	8.041E-01	8.036E-01	8.036E-01	8.035E-01	8.035E-01	8.035E-01	8.035E-01
1.6	1.000E+00	8.958E-01	8.743E-01	8.696E-01	8.685E-01	8.682E-01	8.682E-01	8.682E-01	8.682E-01	8.682E-01	8.682E-01
1.8	1.000E+00	9.469E-01	9.362E-01	9.339E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 512

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.190E+00	1.600E+01	2.786E+01	4.850E+01	8.445E+01	1.470E+02	2.560E+02
-1.6	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.278E+00	8.000E+00	1.213E+01	1.838E+01	2.786E+01	4.222E+01	6.400E+01
-1.4	1.000E+00	1.319E+00	1.741E+00	2.297E+00	3.031E+00	4.000E+00	5.277E+00	6.964E+00	9.189E+00	1.212E+01	1.600E+01
-1.2	1.000E+00	1.149E+00	1.319E+00	1.516E+00	1.741E+00	2.000E+00	2.297E+00	2.638E+00	3.030E+00	3.481E+00	3.998E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.711E-01	7.589E-01	6.614E-01	5.765E-01	5.026E-01	4.383E-01	3.823E-01	3.336E-01	2.911E-01	2.542E-01
-.6	1.000E+00	7.608E-01	5.800E-01	4.433E-01	3.398E-01	2.614E-01	2.020E-01	1.570E-01	1.229E-01	9.709E-02	7.751E-02
-.4	1.000E+00	6.712E-01	4.557E-01	3.143E-01	2.213E-01	1.601E-01	1.198E-01	9.316E-02	7.561E-02	6.403E-02	5.639E-02
-.2	1.000E+00	6.085E-01	3.870E-01	2.613E-01	1.897E-01	1.487E-01	1.253E-01	1.118E-01	1.041E-01	9.970E-02	9.716E-02
0	1.000E+00	5.803E-01	3.761E-01	2.764E-01	2.274E-01	2.031E-01	1.911E-01	1.851E-01	1.821E-01	1.806E-01	1.799E-01
.2	1.000E+00	5.863E-01	4.136E-01	3.413E-01	3.108E-01	2.978E-01	2.922E-01	2.898E-01	2.887E-01	2.883E-01	2.881E-01
.4	1.000E+00	6.162E-01	4.781E-01	4.284E-01	4.104E-01	4.039E-01	4.014E-01	4.005E-01	4.002E-01	4.001E-01	4.000E-01
.6	1.000E+00	6.576E-01	5.504E-01	5.169E-01	5.064E-01	5.031E-01	5.021E-01	5.018E-01	5.016E-01	5.016E-01	5.016E-01
.8	1.000E+00	7.033E-01	6.211E-01	5.984E-01	5.921E-01	5.904E-01	5.899E-01	5.898E-01	5.898E-01	5.898E-01	5.898E-01
1	1.000E+00	7.503E-01	6.879E-01	6.723E-01	6.684E-01	6.674E-01	6.672E-01	6.671E-01	6.671E-01	6.671E-01	6.671E-01
1.2	1.000E+00	7.980E-01	7.515E-01	7.406E-01	7.381E-01	7.375E-01	7.373E-01	7.373E-01	7.373E-01	7.373E-01	7.373E-01
1.4	1.000E+00	8.463E-01	8.132E-01	8.057E-01	8.040E-01	8.036E-01	8.035E-01	8.035E-01	8.035E-01	8.035E-01	8.035E-01
1.6	1.000E+00	8.958E-01	8.743E-01	8.696E-01	8.685E-01	8.682E-01	8.682E-01	8.681E-01	8.681E-01	8.681E-01	8.681E-01
1.8	1.000E+00	9.469E-01	9.362E-01	9.339E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 1024

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2 :	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8 :	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.190E+00	1.600E+01	2.786E+01	4.850E+01	8.445E+01	1.470E+02	2.560E+02
-1.6 :	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.278E+00	8.000E+00	1.213E+01	1.838E+01	2.786E+01	4.222E+01	6.400E+01
-1.4 :	1.000E+00	1.319E+00	1.741E+00	2.297E+00	3.031E+00	4.000E+00	5.278E+00	6.964E+00	9.189E+00	1.213E+01	1.600E+01
-1.2 :	1.000E+00	1.149E+00	1.319E+00	1.516E+00	1.741E+00	2.000E+00	2.297E+00	2.639E+00	3.031E+00	3.482E+00	3.999E+00
-1 :	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8 :	1.000E+00	8.708E-01	7.585E-01	6.607E-01	5.756E-01	5.015E-01	4.370E-01	3.809E-01	3.320E-01	2.894E-01	2.524E-01
-.6 :	1.000E+00	7.598E-01	5.781E-01	4.406E-01	3.364E-01	2.575E-01	1.977E-01	1.524E-01	1.181E-01	9.210E-02	7.238E-02
-.4 :	1.000E+00	6.683E-01	4.506E-01	3.076E-01	2.134E-01	1.514E-01	1.105E-01	8.348E-02	6.569E-02	5.396E-02	4.622E-02
-.2 :	1.000E+00	6.031E-01	3.780E-01	2.499E-01	1.769E-01	1.351E-01	1.112E-01	9.744E-02	8.956E-02	8.504E-02	8.244E-02
0 :	1.000E+00	5.737E-01	3.658E-01	2.639E-01	2.138E-01	1.890E-01	1.767E-01	1.705E-01	1.675E-01	1.659E-01	1.652E-01
.2 :	1.000E+00	5.811E-01	4.059E-01	3.323E-01	3.012E-01	2.880E-01	2.823E-01	2.798E-01	2.788E-01	2.783E-01	2.781E-01
.4 :	1.000E+00	6.133E-01	4.740E-01	4.238E-01	4.056E-01	3.990E-01	3.965E-01	3.956E-01	3.953E-01	3.951E-01	3.951E-01
.6 :	1.000E+00	6.564E-01	5.486E-01	5.150E-01	5.045E-01	5.012E-01	5.001E-01	4.998E-01	4.997E-01	4.997E-01	4.996E-01
.8 :	1.000E+00	7.028E-01	6.204E-01	5.977E-01	5.915E-01	5.897E-01	5.893E-01	5.891E-01	5.891E-01	5.891E-01	5.891E-01
1 :	1.000E+00	7.502E-01	6.877E-01	6.721E-01	6.682E-01	6.672E-01	6.670E-01	6.669E-01	6.669E-01	6.669E-01	6.669E-01
1.2 :	1.000E+00	7.979E-01	7.514E-01	7.406E-01	7.380E-01	7.374E-01	7.373E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01
1.4 :	1.000E+00	8.463E-01	8.131E-01	8.057E-01	8.040E-01	8.036E-01	8.035E-01	8.035E-01	8.035E-01	8.035E-01	8.035E-01
1.6 :	1.000E+00	8.958E-01	8.743E-01	8.696E-01	8.685E-01	8.682E-01	8.682E-01	8.681E-01	8.681E-01	8.681E-01	8.681E-01
1.8 :	1.000E+00	9.469E-01	9.362E-01	9.339E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2 :	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 2048

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2 :	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8 :	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.190E+00	1.600E+01	2.786E+01	4.850E+01	8.445E+01	1.470E+02	2.560E+02
-1.6 :	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.278E+00	8.000E+00	1.213E+01	1.838E+01	2.786E+01	4.222E+01	6.400E+01
-1.4 :	1.000E+00	1.320E+00	1.741E+00	2.297E+00	3.031E+00	4.000E+00	5.278E+00	6.964E+00	9.189E+00	1.213E+01	1.600E+01
-1.2 :	1.000E+00	1.149E+00	1.319E+00	1.516E+00	1.741E+00	2.000E+00	2.297E+00	2.639E+00	3.031E+00	3.482E+00	4.000E+00
-1 :	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8 :	1.000E+00	8.707E-01	7.582E-01	6.603E-01	5.751E-01	5.009E-01	4.363E-01	3.800E-01	3.311E-01	2.885E-01	2.514E-01
-.6 :	1.000E+00	7.591E-01	5.768E-01	4.388E-01	3.342E-01	2.549E-01	1.949E-01	1.494E-01	1.149E-01	8.881E-02	6.901E-02
-.4 :	1.000E+00	6.662E-01	4.468E-01	3.025E-01	2.073E-01	1.448E-01	1.035E-01	7.628E-02	5.831E-02	4.646E-02	3.864E-02
-.2 :	1.000E+00	5.987E-01	3.706E-01	2.407E-01	1.664E-01	1.240E-01	9.962E-02	8.565E-02	7.763E-02	7.303E-02	7.038E-02
0 :	1.000E+00	5.681E-01	3.570E-01	2.534E-01	2.023E-01	1.770E-01	1.644E-01	1.582E-01	1.551E-01	1.535E-01	1.527E-01
.2 :	1.000E+00	5.768E-01	3.995E-01	3.250E-01	2.934E-01	2.800E-01	2.742E-01	2.717E-01	2.706E-01	2.701E-01	2.699E-01
.4 :	1.000E+00	6.111E-01	4.709E-01	4.204E-01	4.021E-01	3.953E-01	3.929E-01	3.920E-01	3.916E-01	3.915E-01	3.914E-01
.6 :	1.000E+00	6.555E-01	5.475E-01	5.138E-01	5.032E-01	4.999E-01	4.988E-01	4.985E-01	4.984E-01	4.984E-01	4.984E-01
.8 :	1.000E+00	7.026E-01	6.201E-01	5.973E-01	5.911E-01	5.894E-01	5.889E-01	5.888E-01	5.887E-01	5.887E-01	5.887E-01
1 :	1.000E+00	7.501E-01	6.876E-01	6.720E-01	6.681E-01	6.671E-01	6.669E-01	6.668E-01	6.668E-01	6.668E-01	6.668E-01
1.2 :	1.000E+00	7.979E-01	7.514E-01	7.405E-01	7.380E-01	7.374E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01
1.4 :	1.000E+00	8.463E-01	8.131E-01	8.057E-01	8.040E-01	8.036E-01	8.035E-01	8.035E-01	8.034E-01	8.034E-01	8.034E-01
1.6 :	1.000E+00	8.958E-01	8.743E-01	8.696E-01	8.685E-01	8.682E-01	8.682E-01	8.681E-01	8.681E-01	8.681E-01	8.681E-01
1.8 :	1.000E+00	9.469E-01	9.362E-01	9.339E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2 :	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 4096

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.190E+00	1.600E+01	2.786E+01	4.850E+01	8.445E+01	1.470E+02	2.560E+02
-1.6	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.278E+00	8.000E+00	1.213E+01	1.838E+01	2.786E+01	4.222E+01	6.400E+01
-1.4	1.000E+00	1.320E+00	1.741E+00	2.297E+00	3.031E+00	4.000E+00	5.278E+00	6.964E+00	9.190E+00	1.213E+01	1.600E+01
-1.2	1.000E+00	1.149E+00	1.319E+00	1.516E+00	1.741E+00	2.000E+00	2.297E+00	2.639E+00	3.031E+00	3.482E+00	4.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.706E-01	7.581E-01	6.601E-01	5.748E-01	5.005E-01	4.358E-01	3.796E-01	3.306E-01	2.879E-01	2.508E-01
-.6	1.000E+00	7.587E-01	5.760E-01	4.376E-01	3.327E-01	2.533E-01	1.931E-01	1.474E-01	1.129E-01	8.665E-02	6.679E-02
-.4	1.000E+00	6.646E-01	4.440E-01	2.987E-01	2.030E-01	1.399E-01	9.833E-02	7.089E-02	5.279E-02	4.085E-02	3.297E-02
-.2	1.000E+00	5.950E-01	3.645E-01	2.330E-01	1.578E-01	1.148E-01	9.007E-02	7.590E-02	6.777E-02	6.310E-02	6.042E-02
0	1.000E+00	5.633E-01	3.494E-01	2.443E-01	1.925E-01	1.667E-01	1.539E-01	1.476E-01	1.444E-01	1.428E-01	1.420E-01
.2	1.000E+00	5.733E-01	3.943E-01	3.189E-01	2.870E-01	2.734E-01	2.675E-01	2.650E-01	2.639E-01	2.634E-01	2.632E-01
.4	1.000E+00	6.095E-01	4.686E-01	4.178E-01	3.994E-01	3.926E-01	3.902E-01	3.892E-01	3.889E-01	3.887E-01	3.887E-01
.6	1.000E+00	6.550E-01	5.668E-01	5.130E-01	5.024E-01	4.991E-01	4.980E-01	4.977E-01	4.975E-01	4.975E-01	4.975E-01
.8	1.000E+00	7.024E-01	6.199E-01	5.971E-01	5.909E-01	5.891E-01	5.887E-01	5.885E-01	5.885E-01	5.885E-01	5.885E-01
1	1.000E+00	7.500E-01	6.876E-01	6.719E-01	6.680E-01	6.670E-01	6.668E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01
1.2	1.000E+00	7.979E-01	7.514E-01	7.405E-01	7.380E-01	7.374E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01
1.4	1.000E+00	8.463E-01	8.131E-01	8.057E-01	8.040E-01	8.036E-01	8.035E-01	8.034E-01	8.034E-01	8.034E-01	8.034E-01
1.6	1.000E+00	8.958E-01	8.743E-01	8.696E-01	8.685E-01	8.682E-01	8.682E-01	8.681E-01	8.681E-01	8.681E-01	8.681E-01
1.8	1.000E+00	9.469E-01	9.362E-01	9.339E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

B3(2,M,r,mu) for r = 8192

Mu \ M=	1	2	4	8	16	32	64	128	256	512	1024
-2	1.000E+00	2.000E+00	4.000E+00	8.000E+00	1.600E+01	3.200E+01	6.400E+01	1.280E+02	2.560E+02	5.120E+02	1.024E+03
-1.8	1.000E+00	1.741E+00	3.031E+00	5.278E+00	9.190E+00	1.600E+01	2.786E+01	4.850E+01	8.445E+01	1.470E+02	2.560E+02
-1.6	1.000E+00	1.516E+00	2.297E+00	3.482E+00	5.278E+00	8.000E+00	1.213E+01	1.838E+01	2.786E+01	4.222E+01	6.400E+01
-1.4	1.000E+00	1.320E+00	1.741E+00	2.297E+00	3.031E+00	4.000E+00	5.278E+00	6.964E+00	9.190E+00	1.213E+01	1.600E+01
-1.2	1.000E+00	1.149E+00	1.320E+00	1.516E+00	1.741E+00	2.000E+00	2.297E+00	2.639E+00	3.031E+00	3.482E+00	4.000E+00
-1	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00
-.8	1.000E+00	8.706E-01	7.580E-01	6.599E-01	5.746E-01	5.003E-01	4.356E-01	3.793E-01	3.303E-01	2.876E-01	2.505E-01
-.6	1.000E+00	7.584E-01	5.754E-01	4.368E-01	3.317E-01	2.522E-01	1.918E-01	1.461E-01	1.115E-01	8.523E-02	6.533E-02
-.4	1.000E+00	6.634E-01	4.418E-01	2.959E-01	1.997E-01	1.363E-01	9.444E-02	6.685E-02	4.864E-02	3.663E-02	2.871E-02
-.2	1.000E+00	5.920E-01	3.594E-01	2.266E-01	1.506E-01	1.071E-01	8.211E-02	6.778E-02	5.955E-02	5.483E-02	5.212E-02
0	1.000E+00	5.591E-01	3.429E-01	2.365E-01	1.839E-01	1.578E-01	1.448E-01	1.383E-01	1.351E-01	1.335E-01	1.327E-01
.2	1.000E+00	5.704E-01	3.899E-01	3.139E-01	2.817E-01	2.679E-01	2.620E-01	2.594E-01	2.583E-01	2.578E-01	2.576E-01
.4	1.000E+00	6.083E-01	4.669E-01	4.159E-01	3.974E-01	3.906E-01	3.881E-01	3.872E-01	3.868E-01	3.867E-01	3.867E-01
.6	1.000E+00	6.546E-01	5.463E-01	5.124E-01	5.018E-01	4.985E-01	4.975E-01	4.971E-01	4.970E-01	4.970E-01	4.970E-01
.8	1.000E+00	7.023E-01	6.198E-01	5.970E-01	5.907E-01	5.890E-01	5.885E-01	5.884E-01	5.884E-01	5.884E-01	5.884E-01
1	1.000E+00	7.500E-01	6.875E-01	6.719E-01	6.680E-01	6.670E-01	6.668E-01	6.667E-01	6.667E-01	6.667E-01	6.667E-01
1.2	1.000E+00	7.979E-01	7.514E-01	7.405E-01	7.380E-01	7.374E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01	7.372E-01
1.4	1.000E+00	8.463E-01	8.131E-01	8.057E-01	8.040E-01	8.036E-01	8.035E-01	8.034E-01	8.034E-01	8.034E-01	8.034E-01
1.6	1.000E+00	8.958E-01	8.743E-01	8.696E-01	8.685E-01	8.682E-01	8.682E-01	8.681E-01	8.681E-01	8.681E-01	8.681E-01
1.8	1.000E+00	9.469E-01	9.362E-01	9.339E-01	9.333E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01	9.332E-01
2	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00	1.000E+00

E. APPENDIX - Notes and Errata

The notes listed below are included to bring attention to notation and other problems. The page number provides reference to the location of the problem or comment.

1. page TN-4

AM noise is especially important in the measurement of the (residual) phase noise added by amplifiers and other signal handling components. Often the driving source for the measurements is a frequency synthesizer that has phase and/or amplitude noise that is comparable or larger than the added noise of the component under test. Most measurement systems are configured so that the phase noise of the source cancels out to a large degree (at high Fourier frequencies decorrelation effects limit the cancellation). In such measurement systems, the AM to PM conversion factor and the AM noise of the source may then set the noise floor. Discussion of the effect of AM noise and AM to PM conversion factors on the accuracy and precision of phase noise measurements is found in "Residual Phase Noise Measurements of vhf, uhf, and Microwave Components" by G. K. Montress, T. E. Parker, and M. J. Loboda Proc. *43rd Annual Symposium on Frequency Control*, pp. 349-359 (1989) and in "Accuracy Model for Phase Noise Measurements by F. L. Walls, C. M. Felton, and A. J. D. Clements, *21st Annual Precision Time and Time Interval Meeting* (1989). The notation in these papers as well as that in other parts of the literature differs from that given below. Our notation below is drawn from a modest level of consensus among individuals responsible for setting standards. We expect that it will gradually be adopted within standards committees. The following comments are directed specifically at the specification of noise performance.

The total power spectral density in a signal can be approximated by expanding eq 12-5 of paper B.2 (by Stein) and extending it to include the spectral density of relative amplitude fluctuations, $S_a(f)$. The double-sideband density written in single-sideband form is given by

$$S_V(f) \approx \frac{V_0^2}{2} \left[e^{-I(f)} \delta(f) + e^{-I(f)} S_\phi(f) + S_a(f) \right] \quad 0 < f < \infty,$$

where

$$I(f) \approx \int_f^\infty S_\phi(f) df.$$

$I(f)$ is the integrated phase modulation due to the pedestal and $\delta(f)$ represents the carrier with frequency width $\pm f_c$. The effect of large $S_a(f)$ on power in the carrier has not, to our knowledge, been explored. The power spectral density of relative phase fluctuations, $S_\phi(f)$, is normalized to one rad²/Hz and $S_a(f)$ is normalized to the carrier voltage, but the total power spectral density, $S_V(f)$ is not normalized and has the units of V²/Hz. All of these are single-sided spectral densities. For most measurement purposes, we can disregard the carrier and find that, away from the carrier, the above expression simplifies to

$$S_V(f) = \frac{V_0^2}{2} [S_\phi(f) + S_a(f)] \quad \text{for} \quad 0 \leq f \leq \infty.$$

The more general expression is important only very near the carrier and in certain types of frequency multiplication. The single-sideband, amplitude noise, normalized to the total signal power is given simply by $S_a(f)$ for $0 \leq f \leq \infty$. The measurement of added phase and amplitude noise for amplifiers and other signal handling components should specify the signal level since the AM noise level and the contribution due to AM-to-PM conversion depend on the signal level.

2. page TN-6

For a direct measurement, time accuracy only has meaning when the phase of the time-base oscillator of the frequency counter is known with respect to some time standard. Either it is phase locked or is calibrated with respect to that standard at the time of measurement. The phase of the time-base oscillator can then be measured with respect to the phase of the frequency standard being calibrated (accounting for cable delays, etc.). Except for the cycle ambiguity of the carrier, the phase of the frequency standard being measured can carry time information and have time accuracy. This technique is not common, but is very useful and eliminates divider noise that typically occurs in going from 5 or 10 Mhz to 1 pulse per second. Caution must be exercised to assure that the phase point measured in a sine wave is at a reproducible voltage and impedance so that the cycle ambiguity is an exact integer.

3. page TN-35

A second-order servo loop provides substantially enhanced performance. See, for example, F.L. Walls and S.R. Stein, "Servo techniques in oscillators and measurement systems," *NBS Tech. Note 692* (1976).

4. page TN-35

This error can be identified and corrected using the phase modulation scheme described in paper B.4 on page TN-136.

5. page TN-36

Low-noise DC amplifiers have been substantially improved since publication of this paper.

6. pages TN-37, TN-91, TN-130, TN-174, TN-206 and TN-218

The reader is reminded that the discussions of frequency-domain measurements assume incoherent noise processes. Often the phase noise spectrum of a signal will contain bright spectral features (spurious lines) other than the carrier. Frequency-domain measurements are often useful in identifying such features. But if these spurious lines are narrow compared to the measurement bandwidth, statistical measures such as $S_\phi(f)$ and $\varrho(f)$ are not appropriate. It is better to specify the phase deviation in terms of the rms value of the phase deviations, ϕ_{rms} (rms radians), without reference to bandwidth. This specification in rms radians can be related to the Allan variance (see note # 8 below).

7. page TN-51

Humidity is often an important environmental factor. See, for example, J.E. Gray, H.E. Machlan and D.W. Allan, "The Effect of humidity on commercial cesium beam atomic clocks," *42nd Annual Symp. on Frequency Control*, pp. 514-518 (1988) and F.L. Walls, "The

Influence of Pressure and Humidity on the Medium and Long-Term Frequency Stability of Quartz Oscillators," *42nd Annual Symp. on Frequency Control*, pp. 279-283 (1988).

8. page TN-74

For a bright line (one which is narrow compared with the measurement bandwidth), the solution of eq 12-27 simplifies to

$$\sigma_y(\tau) = \frac{\sqrt{2} \phi_{rms}}{\pi \nu_0 \tau} \sin^2(\pi f \tau),$$

where ϕ_{rms} is the rms value of the phase deviations. The above relationship may be useful where one is trying to determine the effect of a bright line in the time domain. If the bright line is the dominant factor, the plot of $\sigma_y(\tau)$ versus τ has strong $\sin^2(\pi f \tau)$ oscillations and it can be ambiguous. In that case, it is better to provide a specification in terms of ϕ_{rms} without reference to bandwidth. Statistical measures such as $\sigma_y(\tau)$ and $S_\phi(f)$ are not meant to be used to describe coherent signals. For further discussion see paper B.1, section 12.2 (page TN-51).

9. page TN-75

A set of brackets, [], are missing in eq 12-29. The equation should read

$$\text{mod} \sigma_y^2(\tau) = \frac{1}{2\tau^2 N^2 (N-3n+1)} \sum_{j=1}^{N-3n+1} \left[\sum_{i=j}^{n+j-1} (x_{i+2n} - 2x_{i+n} + x_i) \right]^2.$$

10. page TN-119

The reference (Walls and DeMarchi, 1975) is listed as being on pages 310-317. The page numbers should be 210-217.

11. page TN-121

Most of the literature uses the expression $[V_0 + \epsilon(t)]$ instead of V_0 as in eq (2). V_0 is the peak voltage amplitude and $\epsilon(t)$ is the voltage deviation of the amplitude from nominal.

12. page TN-122

In eq (5), most of the literature uses $x_1(t)$ instead of $\epsilon(t)$. $\epsilon(t)$ is usually the voltage deviation as described in note 11 above.

13. page TN-123

There is an error in the caption for figure 7. The last portion of that caption should read: "where f is Fourier frequency ($\omega = 2\pi f$) and $S_y(f) = \omega^2 S_x(f)$."

14. page TN-123

In the right-hand column, last paragraph, 5th line, there is an extraneous minus sign. The quantity $\bar{y}_i^{\tau 0}$ should read \bar{y}_i^0 . Also, note that the use of superscript τ and τ_0 with \bar{y} is not consistent with the new IEEE standard definitions (see paper C.1, page TN-139).

15. page TN-124
In eq 9 the subscript, k - n, should read k + n. That is, the equation should read

$$\bar{y}_k^\tau = \frac{1}{n} \sum_{i=k}^{k+n-1} \bar{y}_i^{\tau_0} = \frac{x_{k+n} - x_k}{\tau}.$$

16. page TN-125
The quantity $\bar{\sigma}_y^2(\tau)$ is now commonly known as $\text{mod}\sigma_y^2(\tau)$. This latter form has been recently adopted by IEEE as the standard terminology.

17. page TN-139
IEEE Std 1139-1988, *IEEE Standard Definitions of Physical Quantities for Fundamental Frequency and Time Metrology*, is an almost exact replica of this paper (D.1). The paper was published during the latter period of the development of the standard. The only substantial difference is that, wherever it occurs, the word "departure" in the paper is replaced in the IEEE standard by "deviation."

18. page TN-146
This widely cited paper provided the *de facto* standards for terminology and oscillator characterization until the recent adoption of the IEEE standard presented in paper C.1 (page TN-139). For terminology, the latest IEEE standard should always take precedence. This paper (C.2) is fairly consistent with the IEEE standard. One exception is that, in this paper, N denotes the number of frequency measurements. The symbol M in the IEEE standard is the same as N in this paper. In the standard, the equation relating M and N is $M = N - 1$.

19. page TN-151
Equation (23) can be substantially simplified as shown, for example, by Stein (page TN-74, eq (12-27)), which is

$$\sigma_y^2(\tau) = \frac{2}{(\pi v_0 \tau)^2} \int_0^{\infty} S_{\phi}(f) \sin^4(\pi f \tau) df.$$

20. page TN-154
In eq 36 the T outside the brackets should be a τ . The equation should read

$$\hat{x}(t_0 + \tau) = x(t_0) + \tau \left[\frac{x(t_0) - x(t_0 - T)}{T} \right].$$

21. page TN-160
The 2 expressions listed as eq (95) are in error. They should read

$$\begin{aligned} -h_{-1} \tau^2 [3 + 2 \ln r - 1/(6r^2)] & \quad r > 1 \\ -h_{-1} T^2 [3 - 2 \ln r] & \quad r < 1. \end{aligned}$$

22. page TN-160
 In eqs (101), (102), (103), (104), and (105), a Greek γ was mistakenly replaced by the number 2. In each of these equations the quantity $[2 + \ln(2\pi f_h \tau)]$ should be replaced by the quantity $[\gamma + \ln(2\pi f_h \tau)]$. γ , Euler's constant, has the value 0.5772156649.....
23. page TN-162
 This paper is included in this collection because it presents the internationally accepted terminology and definitions. There are no substantial inconsistencies with the new IEEE standard (paper C.1, page TN-139), but the latest IEEE standard should be considered to be the most up-to-date authority. A new version of this international report should be issued by the CCIR in 1990.
24. page TN-171
 The definitions for symbols used in this paper are fairly consistent with those adopted in the recent IEEE standard (C.1). One exception is that, in this paper, N denotes the number of frequency measurements. The symbol M in the IEEE standard is the same as N in this paper. Another is that, while this paper uses μ as the exponent of τ in describing the power-law noise processes, the paper adopts the opposite sign convention for μ . $v(t)$ is used where many other papers use $V(t)$ for instantaneous voltage. Some confusion is generated when this small v is typeset in the equations to look almost identical to the Greek ν , a symbol which is used exclusively to represent frequency. For example, In equation (1) the left-hand quantity is voltage, whereas the $\nu(t)$ and ν_0 in equation (4) are clearly frequencies. Finally, the authors of this paper, in equation (2), define $\epsilon(t)$ as the normalized amplitude fluctuations, a very sound choice, but the reader should note that most other papers have not normalized it.
25. page TN-175
 For consistency with figure 12 and the text, $\nu(t)$, the left-hand member of eq (11) should probably be $u(t)$.
26. page TN-177
 Walls, Percival and Irelan (D.4) have recently addressed the more accurate specification of the quantity p in eq (12).
27. page TN-179
 It is important to note that the expressions in Table 2 in this paper are derived assuming use of a single-pole filter. The calculations can also be done using an ideal (infinitely sharp) filter. The solutions in these two limits are useful because they define the range of practical values (using n-pole filters) for the expressions. Table I in this section is an expansion of Table 2 of Lesage and Audoin providing both the single-pole results as well as the results for an infinitely sharp filter. There are discrepancies in several of the coefficients between terms in Table 2 in the paper and those in Table I on the next page.
28. page TN-180
 Barnes and Allan (paper D.8) have recently completed further analysis of the effect of dead time on measurements.

Table I. Asymptotic forms of $\sigma_y^2(\tau)$ for various power-law types and two filter types. Note: $\omega_h/2\pi = f_h$ is the measurement system bandwidth, often called the high-frequency cutoff. $\ln \equiv \log e$.

Name of Noise	α	$S_y(f)$	$\sigma_y^2(\tau)$			
			$\omega_h \tau \gg 1$ Infinite Sharp Filter	$\omega_h \tau >> 1$ Single Pole Filter	$\omega_h \tau < < 1$ Infinite Sharp Filter	$\omega_h \tau < < 1$ Single Pole Filter
White Phase	2	$h_2 f^2$	$\frac{3f_h h_2}{(2\pi)^2 \tau^2}$	$\frac{3f_h h_2}{(2\pi)^2 \tau^2}$	$\frac{2\pi^2 f_h^2 \tau^2 h_2}{5}$	$\frac{f_h^2 h_2}{2\tau}$
Flicker Phase	1	$h_1 f$	$\frac{(1.038 + 3\ln(\omega_h \tau))h_1}{(2\pi)^2 \tau^2}$	$\frac{(3\ln(\omega_h \tau))h_1}{(2\pi)^2 \tau^2}$	$\frac{\pi^2 f_h^4 \tau^2 h_1}{2}$	$2f_h^2 (\ln(2)) h_1$
White Frequency	0	h_0	$\frac{h_0}{2\tau}$	$\frac{h_0}{2\tau}$	$\frac{2\pi^2 f_h^3 \tau^2 h_0}{3}$	$\frac{2\pi^2 f_h^2 \tau h_0}{3}$
Flicker Frequency	-1	$h_{-1} f^{-1}$	$\frac{2(\ln(2))h_{-1}}{2\pi^2 \tau h_2}$	$\frac{2(\ln(2))h_{-1}}{2\tau}$	$\frac{\pi^2 f_h^2 \tau^2 h_{-1}}{3}$	$8\pi^2 f_h^2 \tau^2 h_{-1}$
Random-Walk Frequency	-2	$h_{-2} f^{-2}$	$\frac{2\pi^2 \tau h_2}{3}$	$\frac{2\pi^2 \tau h_2}{3}$	$\frac{2\pi^2 f_h \tau^2 h_2}{3}$	$2\pi^2 f_h \tau^2 h_2$

29. page TN-180
If the ratio of T/τ is constant and greater than 1 (the usual case), the problem described is eliminated. However, in taking data for a plot of $\sigma_y(\tau)$ versus τ , it is difficult to achieve this in the hardware and not possible to do it with software processing alone. For further discussion see paper D.8 (page TN-296).
30. page TN-197
The most recent definitions and concepts for spectral density are given in a new IEEE standard (paper C.1 on page TN-138). This new standard should be consulted as the latest authority on definitions and terminology.
31. page TN-198
The newly accepted definition of $\ell(f)$ is given in paper C.1. This new definition, $\ell(f) \equiv \frac{1}{2}S_{\phi}(f)$, was always valid for Fourier frequencies far from the carrier. It has now been extended to cover all frequencies.
32. page TN-217
Equation (73) should read $20 \log(\text{final frequency/original frequency})$.
33. page TN-239
On page TN-198 the authors refer to a paper by Glaze (1970). The reference, apparently lost in printing, is: Glaze, D.J. (1970). "Improvements in Atomic Beam Frequency Standards at the National Bureau of Standards," *IEEE Trans. Instrum. Meas.* IM-19(3), 156-160.
34. page TN-257
There are two errors in Table 2. Under R(n) the first entry should be 1/n rather than 1. The second item in the same column is not single valued (1), but takes on different values for different measurement bandwidths. The reader is referred to section A.6 (page TN-9) for a discussion of this topic.
35. pages TN-261 and TN-262
Subsequent work on $\text{mod}\sigma_y^2(\tau)$ and R(n) is reported in section A.6 (page TN-9) of this report. There are some differences between the results in A.6 and the ones reported in Tables I and II and Figure 4 in this paper.
36. page TN-264
To be consistent with other papers in the literature, $\phi(t)$ in eq (1) should probably be written as $x_1(t)$.
37. page TN-268
The term $\epsilon(t)$ in eq (3) is normally used to represent the amplitude fluctuations in the output voltage of an oscillator. This term might be better designated $Y_1(t)$.

38. Page TN-30

There are two errors in eq 6.6. The equation for Flicker PM is missing a square root in the exponential. It should read

$$\text{Flicker PM} \quad \text{d.f.} \approx \exp\left(\ln\frac{N-1}{2n} \ln\frac{(2n+1)(N-1)}{4}\right)^{\frac{1}{2}}.$$

In the equation for Flicker FM, the numerator of the upper term should read $2(N-2)^2$ instead of $2(N-2)$. The equation should read

$$\text{Flicker FM} \quad \text{d.f.} \approx \begin{cases} \frac{2(N-2)^2}{2.3N-4.9}, & \text{for } n=1 \\ \frac{5N^2}{4n(N+3n)}, & \text{for } n \geq 2 \end{cases}$$

39. Page TN-85

There are errors in two of the terms of Table 12-4. The small n in the denominator of the first logarithmic term for flicker phase should be an m , and the whole quantity in the square bracket should have an exponent of $\frac{1}{2}$. The equation should read

$$\text{Flicker phase} \quad \exp\left[\ln\left(\frac{N-1}{2m}\right) \ln\left(\frac{(2m+1)(N-1)}{4}\right)\right]^{\frac{1}{2}}.$$

The $(N-2)$ in the flicker frequency term should be replaced by $(N-2)^2$ so that df is

$$\text{Flicker frequency} \quad \frac{2(N-2)^2}{2.3N-4.9} \quad \text{for } m=1.$$

40. Page TN-279

There is a factor N missing in the first expression under eq A6. It should read

$$A = 3[3N(N+1)+2].$$

41. Page TN-307

The fifth identity for B_1 , $B_1(N,1,\mu) = 1$ for $\mu < 0$ and $= [2/N(N-1)] \dots$ for $\mu > 0$, should be deleted. The identity for $\mu < 0$ is incorrect and that for $\mu > 0$ is superfluous considering the identity above it which covers all cases for $\mu \neq 0$.